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SOLA-DF: A Solution Algorithm for Nonequilibrium Two-Phase Flow

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CONTENTS

ABSTRACT	1
I. INTRODUCTION	1
II. EQUATIONS AND CONSTITUTIVE RELATIONS	4
A. Equations of Motion	4
B. Void Fraction	6
C. Equation of State	6
D. Vapor and Liquid Internal Energies	7
E. Relative Velocity	8
F. Phase Transitions	10
G. Pipe Friction	13
III. NUMERICAL SOLUTION METHOD	14
A. Mesh Construction	15
B. Solution Algorithm	15
1. Explicit Updating of Velocities	16
2. Implicit Pressure Calculation	21
3. Updating of Remaining Variables	24
C. Boundary Conditions	29
1. Rigid Free-Slip	29
2. Rigid No-Slip	29
3. Continuative Outflow	30
4. Periodic	30
5. Constant Pressure	31
D. Internal Obstacles	32
E. Variable Time Steps	32
IV. INPUT DATA, COMMON VARIABLES, AND SUBROUTINES AND GRAPHICS OUTPUT	33
A. Input Data	33
B. COMMON Variables	38
C. Subroutines and Graphics Output	42
V. EXAMPLE PROBLEM	42
VI. SUMMARY	43
ACKNOWLEDGMENTS	61
REFERENCES	62
APPENDIX: FORTRAN IV Listing of the SOLA-DF Code	64

SOLA-DF: A SOLUTION ALGORITHM FOR NONEQUILIBRIUM TWO-PHASE FLOW

by

C. W. Hirt, N. C. Romero, M. D. Torrey, and J. R. Travis

ABSTRACT

A numerical solution algorithm, SOLA-DF, is presented for the solution of gas-liquid mixture dynamics in two space dimensions and time. The two-phase system is described by a set of mixture equations plus a relation describing the relative flow of one phase with respect to the other. In addition, the algorithm contains models to represent the interphase exchange rates of mass, momentum, and energy for water-steam mixtures.

I. INTRODUCTION

Fluid dynamic problems involving multiphase mixtures occur in abundance in nearly all branches of engineering and technology. Yet the ability to theoretically study these flows in more than one space dimension and with time-dependent behavior has only recently evolved.^{1,2} For many applications it is unnecessary to solve complete sets of mass, momentum, and energy conservation equations for each phase.³ For example, when the two phases are moving together, a mixture momentum equation can replace the individual momentum equations, or if the phases are in thermodynamic equilibrium, it is unnecessary to keep two energy equations. When small deviations are expected from equilibrium or equal phase velocities, it is possible to introduce correction terms into the mixture equations without having to increase the number of equations. It is the latter approach that is considered in this report.

In the next section a theoretical formulation is described for a two-phase mixture using one variation of the so-called "drift-flux" approximation.⁴

These equations are formulated for two-dimensional planar or axisymmetric coordinate systems. All constitutive relations and equations of state are relatively simple models designed for water-steam mixtures. These relations have been used successfully for many important applications arising in nuclear reactor safety studies.⁵ They should not, however, be used indiscriminately. For other applications, consideration should always be given to their suitability for the temperature-pressure ranges of interest and to the possibility that additional physical processes may be needed in the basic theoretical description. In any case, it is emphasized that the constitutive relations used in this report may be easily changed without altering the basic numerical solution algorithm that forms the heart of the SOLA-DF code.

The solution algorithm used in SOLA-DF has evolved from algorithms used in earlier codes in the SOLA series. The original SOLA code⁶ was designed for problems involving a single, incompressible fluid in a fixed region. The SOLA-SURF code⁶ is an extended form of SOLA that allows for the inclusion of free surfaces. SOLA-ICE,⁷ the third member of the family, was designed to handle single-component, compressible fluids. An implicit solution method is used in SOLA-ICE so that, in addition to shock and rarefaction dominated flows, it can also treat very low speed (incompressible) fluid flows. All of these SOLA codes are available from the National Energy Software Center, 9700 South Cass Avenue, Argonne, Illinois 60439.

The SOLA-DF code is a direct descendent of SOLA-ICE. It utilizes finite-difference approximations with respect to a mesh of equal rectangular cells covering the flow region of interest. A semi-implicit formulation is available as a user option so that large time steps can be used in many circumstances to reduce problem run times.

The solution algorithm also has an option for one-dimensional computations. In this case the mesh is limited to one column of cells, and derivatives of dependent variables normal to this column are automatically suppressed by the setting of a large mesh-interval in the normal direction. Furthermore, for one-dimensional computations of flow in pipes there is an option for a two-phase flow pipe friction model.

In addition, SOLA-DF is formulated with a variable, denoted by A , representing the "thickness" of a mesh cell. That is, in a one-dimensional computation the volume of a cell of length δy is $A\delta y$. In two-dimensional computations the volume of a cell of width δx and height δy is $A\delta x\delta y$. The presence

of the A quantity provides SOLA-DF with many useful features. For example, in one-dimensional applications a variable A can be used to model nozzles or other pipe area changes. If A increases as the square of the distance from the bottom of the mesh column then the equations reduce to a one-dimensional spherical system that could be used to study spherical bubbles or explosive flashing processes.

In two-dimensional applications a variable A may be used to represent flow through a two-dimensional duct of slowly varying thickness. The use of zero values of A in selected mesh cells provides a convenient means of including internal obstacles in the flow region. In particular, a cell with $A = 0$ will allow no flow across its boundaries. Axisymmetric coordinates are generated by having A increase linearly with distance from the axis. In the present version of SOLA-DF the A-quantity must be constant in time. However, it would not be difficult to introduce time-dependent corrections so that flows in flexible pipes and ducts might be modeled.⁸

In constructing SOLA-DF some consideration has been given to providing subroutines for most of the important or controversial constitutive relations and for the equations of state for each phase. Thus, changes in these quantities can be made with relative ease. Although the code was developed on a CDC-7600 computer, ANSI-Standard FORTRAN has been used throughout to facilitate its use on other computers. This does not apply, however, to the output subroutines used in the code for graphic display purposes. A list of these subroutines and their functions is provided in Section IV for those users who have access to graphic display systems and wish to convert them to their system routines. All calls to graphic output routines can be avoided by inputting a plot time larger than the requested problem time.

A sample problem is discussed in Section V to illustrate the type of results that may be obtained with the SOLA-DF code. The results of this test problem are presented in detail so that new users may check that their codes are running correctly.

A listing of the code and its subroutines are provided in the Appendix.

Two SOLA codes,^{5,8} which were referenced above, offer complementary capabilities to SOLA-DF and deserve an expanded explanation. A variation of SOLA-DF is contained in SOLA-LOOP, which is a one-dimensional network code.⁵ The SOLA-LOOP program is designed to handle systems of one-dimensional components

coupled together through junctions. Variable time steps may be used in different components, and provisions are available for defining trips, breaks, valves, and other transient features. SOLA-FLX consists of SOLA-DF coupled to a three-dimensional shell code.⁸ This combination may be used for problems involving coupled fluid-structure interactions in which the structure is a cylindrical shell. Following the basic procedure used in SOLA-FLX, other types of structure models could be coupled to SOLA-DF.

II. EQUATIONS AND CONSTITUTIVE RELATIONS

The drift-flux equations describing the dynamics of two-phase fluid mixtures have been cast in many forms.⁹ In the present case we choose as dependent variables the mixture density ρ , the macroscopic vapor density ρ_v , (vapor mass per unit volume of mixture) the center-of-mass velocity $\underline{u} = (u, v)$, and the mixture specific internal energy I . Important auxiliary variables are the void fraction θ , the relative velocity between phases $\underline{u}_r = (u_r, v_r)$, and the mixture pressure p .

A. Equations of Motion

In terms of the chosen dependent variables, the basic two-dimensional drift-flux equations used in SOLA-DF are

the continuity equations,

$$\frac{\partial \rho}{\partial t} + \frac{1}{A} \left(\frac{\partial A \rho u}{\partial x} + \frac{\partial A \rho v}{\partial y} \right) = 0 \quad (2.1)$$

$$\frac{\partial \rho_v}{\partial t} + \frac{1}{A} \left[\frac{\partial}{\partial x} A \left(\rho_v u + \frac{\rho_v \rho_l}{\rho} u_r \right) + \frac{\partial}{\partial y} A \left(\rho_v v + \frac{\rho_v \rho_l}{\rho} v_r \right) \right] = \Gamma ; \quad (2.2)$$

the momentum equations,

$$\begin{aligned} \frac{\partial p u}{\partial t} + \frac{1}{A} \left[\frac{\partial}{\partial x} A \left(\rho u^2 + \frac{\rho_v \rho_l}{\rho} u_r^2 \right) + \frac{\partial}{\partial y} A \left(\rho u v + \frac{\rho_v \rho_l}{\rho} u_r v_r \right) \right] \\ = - \frac{\partial p}{\partial x} + \rho g_x + f_{visc} \end{aligned}$$

$$\begin{aligned} \frac{\partial \rho v}{\partial t} + \frac{1}{A} \left[\frac{\partial}{\partial x} A \left(\rho u v + \frac{\rho_v \rho_l}{\rho} u_r v_r \right) + \frac{\partial}{\partial y} A \left(\rho v^2 + \frac{\rho_v \rho_l}{\rho} v_r^2 \right) \right] \\ = - \frac{\partial p}{\partial y} + \rho g_y + f_{visy} ; \end{aligned} \quad (2.3)$$

and the internal energy equation,

$$\begin{aligned} \frac{\partial \rho I}{\partial t} + \frac{1}{A} \left\{ \frac{\partial}{\partial x} A \left[\rho I u + \frac{\rho_v \rho_l}{\rho} (I_v - I_l) u_r \right] + \frac{\partial}{\partial y} A \left[\rho I v + \frac{\rho_v \rho_l}{\rho} (I_v - I_l) v_r \right] \right\} \\ = - \frac{p}{A} \left\{ \frac{\partial}{\partial x} A \left[u + \frac{\rho_v \rho_l}{\rho} \left(\frac{1}{\rho_v^0} - \frac{1}{\rho_l^0} \right) u_r \right] + \frac{\partial}{\partial y} A \left[v + \frac{\rho_v \rho_l}{\rho} \left(\frac{1}{\rho_v^0} \right. \right. \right. \\ \left. \left. \left. - \frac{1}{\rho_l^0} \right) v_r \right] \right\} + K(u_r^2 + v_r^2) + \omega_{vis} . \end{aligned} \quad (2.4)$$

In these equations, the independent variables are time t and coordinates x, y . The exchange functions for mass and momentum are Γ and K , respectively. Subscripts v and l refer to properties in the vapor and liquid states while a superscript zero refers to microscopic quantities. The gravitational acceleration is denoted by $\mathbf{g} = (g_x, g_y)$.

The quantity A is the time-independent "thickness" of the flow channel. That is, $A \Delta x \Delta y$ is the volume associated with an area Δx wide and Δy high. In addition to representing variable area ducts, we may use suitably defined A values to represent cylindrical coordinates ($A = x$, the circumferential area per unit azimuthal angle) or obstacle regions ($A = 0$).

To complete these equations a variety of constitutive relations must be defined. It is in the definition of these relations that considerable caution must be exercised. The choices made are governed by the applications for which the code is intended. The best choices are those that can be tested against suitable experimental data. Even with careful testing, however, the prejudices of different researchers often lead to different constitutive relations. In the following we describe one set of simple constitutive models that have been used

in the initial development of the SOLA-DF, SOLA-FLX, and SOLA-LOOP codes. These models should not, therefore, be taken as invariant features of these codes. Instead, the codes are to be thought of as skeletons offering numerical solution algorithms that will work with a variety of constitutive relations.

B. Void Fraction

The volume of vapor per unit volume of mixture, that is, the void fraction θ is defined through the relation

$$\theta = \left(\rho_l^0 - \rho + \rho_v^0 \right) / \rho_l^0 . \quad (2.5)$$

C. Equation of State

In the present formulation of SOLA-DF the equation of state is assumed to be a relation that gives pressure as a function of density and internal energy. Although fits to steam table data could be inserted in the equation-of-state subroutine, for developmental purposes we have chosen to use a much simpler approach. When the void fraction is below a small, predetermined value, $\theta < \theta_c$, (typically $\theta_c \approx 0.005$) the fluid is assumed to be a pure liquid with the equation of state

$$p = p_o + a^2 \left(\rho - \rho_l^0 \right) ,$$

where a is a representative speed of sound in the liquid phase and p_o is chosen (see Eq. 2.6) to ensure pressure continuity between the pure liquid and two-phase states when $\theta = \theta_c$. In the two-phase region, $\theta > \theta_c$, the mixture pressure is equal to that of the vapor and is given by the polytropic gas equation,

$$p = (\gamma - 1) \rho_v^0 I_v .$$

In the code these equations are combined into the single equation

$$p = (\gamma - 1) \rho_v I_v^* / \theta^* + a^2 \rho_l^o (\theta^* - \theta) , \quad (2.6)$$

where $\theta^* = \begin{cases} \theta & \text{if } \theta \geq \theta_c \\ \theta_c & \text{if } \theta < \theta_c \end{cases}$

For saturated conditions we have found that $\gamma = 1.07$ and $a^2 \approx 10^4 \text{ cm}^2/\text{ms}^2$ offer reasonable approximations for many reactor safety problems.

D. Vapor and Liquid Internal Energies

In the equation of state and in the relative velocity convection terms in the internal energy equation, it is necessary to have separate values for vapor and liquid internal energies. Because the basic dependent energy variable is the mixture internal energy, a separate prescription must be given for unfolding the individual phase energies. Two prescriptions have been used. In one the phases are assumed to be at equal temperatures. In the other, the vapor phase is assumed to be saturated. For many applications there is little difference between the two assumptions, because the large heat content of the more massive liquid phase keeps the liquid temperature nearly the same in either case.

To implement either of these assumptions we approximate the temperature dependence of the internal energies as

$$\begin{aligned} I_v &= E_v + C_v (T_v - T_o) \\ I_l &= E_l + C_l (T_l - T_o) , \end{aligned} \quad (2.7)$$

where E_v and E_l are saturated internal energies at temperature T_o and C_v and C_l are constants chosen to fit the steam table versus T curves in the temperature range of interest. For example, in the system of units g, cm, ms, K the

values $E_v = 2.506 \times 10^4$, $E_l = 0.4174 \times 10^4$, $C_v = 6.67$ and $C_l = 44.34$ are good approximations for temperatures up to about $T = 550$ K.

With these definitions, and assuming equal phase temperatures, the mixture temperature can be computed from the mixture internal energy as the solution of the linear equation,

$$\rho I = \rho_v I_v + \rho_l I_l \quad . \quad (2.8)$$

When the vapor is assumed to be saturated its temperature and pressure are related by

$$T_v = 255.2 + 117.8 p^{0.223} \quad , \quad (2.9)$$

where in this expression p must be in bars and T in kelvins. Knowing the vapor temperature it is then an easy task to compute the separate liquid and vapor internal energies and liquid temperature using Eqs. (2.7) and (2.8).

E. Relative Velocity

An equation of motion for the relative velocity can be derived from equations describing a complete two-fluid model,¹⁰ given by

$$\begin{aligned} \frac{\partial u_r}{\partial t} + \frac{1}{2} \frac{\partial}{\partial x} \left\{ u_r \left[2u + \frac{u_r}{\rho} (\rho_l - \rho_v) \right] \right\} &+ v_v \frac{\partial u_v}{\partial y} - v_l \frac{\partial u_l}{\partial y} \\ &= \left(\frac{1}{\rho_l} - \frac{1}{\rho_v} \right) \frac{\partial p}{\partial x} - K \frac{\rho}{\rho_v \rho_l} u_r \\ \frac{\partial v_r}{\partial t} + \frac{1}{2} \frac{\partial}{\partial y} \left\{ v_r \left[2v + \frac{v_r}{\rho} (\rho_l - \rho_v) \right] \right\} &+ u_v \frac{\partial v_v}{\partial x} - u_l \frac{\partial v_l}{\partial x} \\ &= \left(\frac{1}{\rho_l} - \frac{1}{\rho_v} \right) \frac{\partial p}{\partial y} - K \frac{\rho}{\rho_v \rho_l} v_r \end{aligned} \quad (2.10)$$

where $u_r = u_v - u_\ell$, $v_r = v_v - v_\ell$,

$$u = \frac{\rho_v u_v + \rho_\ell u_\ell}{\rho}, \quad v = \frac{\rho_v v_v + \rho_\ell v_\ell}{\rho},$$

$$\rho = \rho_v + \rho_\ell,$$

and K is the interfacial friction coefficient.

The current version of SOLA-DF neglects all but the temporal term on the left side of Eq. 2.10. Assuming the vapor is a dispersed phase of small bubbles when θ is small, or the liquid is a dispersed phase of small droplets when θ is large, we can estimate K from the drag on an individual bubble (or droplet) times the number of bubbles (droplets) per unit volume, N . The result is

$$K = \frac{\rho S}{8\theta_1} \left[C_d |u_r| + \frac{12v}{r_o} \right], \quad (2.11)$$

where

$$\theta_1 = \theta, \quad v = v_\ell (1-\theta)^{-2.5} \quad \text{for } \theta \leq 1/2$$

$$\theta_1 = 1-\theta, \quad v = v_v \theta^{-2.5} \quad \text{for } \theta > 1/2.$$

C_d is a drag coefficient (generally of order unity), S is the cross-sectional area per unit volume of bubbles (droplets) with radius r_o ,

$$S = \begin{cases} \frac{3\theta}{4r_b} \text{ and } r_o = r_b \text{ if } \theta \leq 1/2 \\ \frac{3(1-\theta)}{4r_d} \text{ and } r_o = r_d \text{ if } \theta > 1/2 \end{cases}, \quad (2.12)$$

and ν is the kinematic viscosity. The average radius is related to the number density by the expressions

$$r_b = \left(\frac{3\theta}{4\pi N} \right)^{1/3} \quad \text{for } \theta \leq 1/2$$
$$r_d = \left[\frac{3(1-\theta)}{4\pi N} \right]^{1/3} \quad \text{for } \theta > 1/2 . \quad (2.13)$$

The bubble number N is often assumed to be a constant independent of space and time. This, of course, is an approximation that will not work when preferential nucleating sites are desired. Although N must be estimated for each calculation, as described in the next section, a locally variable N can sometimes be estimated in terms of a critical Weber number.

F. Phase Transitions

The form of the phase change model embodied in Γ is crucial if nonequilibrium effects are to be correctly predicted. The model we describe here is still under development and is not yet sophisticated enough to be used as a predictive tool without some adjustment. Nevertheless, this model has proven useful in numerous applications and its presentation here illustrates the types of considerations necessary in the development of such models.

Defining q as the interfacial heat flux, a simple energy balance shows that

$$\Gamma = \frac{qZ}{\lambda} ,$$

where λ is the latent heat of vaporization and interfacial area Z is related to the bubble radius, r , according to $Z = 3\theta/r$. The heat flux can be further defined as

$$q = k(T_L - T_s)/\lambda ,$$

where T_s is the saturation temperature and k_λ is the thermal conductivity of the liquid whose bulk temperature is T_λ . The length λ characterizes the thickness of the thermal boundary layer over which the liquid temperature changes from its interior, bulk value, T_λ , to the value T_s assumed to exist at the two-phase interface. Thus,

$$\Gamma = \frac{k_\lambda (T_\lambda - T_s)z}{\lambda \lambda} . \quad (2.14)$$

For a single, nontranslating bubble growing in an infinite fluid region, it has been shown¹¹ that $\lambda = \lambda_c$, where

$$\lambda_c = r \left[\frac{6}{\pi} \frac{\rho_l^0 c_l |T_\lambda - T_s|}{\rho_v^0 \lambda} \right]^{-1} .$$

In this expression r is the instantaneous bubble radius, which is defined later.

When the bubbles are translating with respect to the surrounding liquid with speed, U , then $\lambda = \lambda_u$, and Moalem and Sideman¹² give the general expression

$$\lambda_u = r \left(\frac{\pi}{Re_b Pr} \right)^{1/2} ,$$

where $Re_b = 2rU \rho_l^0 / \mu_l$ is the bubble Reynolds number, $Pr = C_l \mu_l / k_l$ is the liquid Prandtl number, and μ_l is the liquid shear viscosity. As the relative speed U increases, the length λ_u rapidly decreases below the value of λ_c , which represents stripping away of the thermal boundary layer by relative flow. In an attempt to smoothly combine both these effects, we have defined λ as the reciprocal average of these limiting characteristic lengths,

$$\frac{1}{\lambda} = \frac{1}{\lambda_c} + \frac{1}{\lambda_u} . \quad (2.15)$$

Equation (2.14) with λ defined by the above equation is a vapor generation rate that includes both finite heat conduction and relative velocity effects. However, the model still requires the definition of r and U .

If we know the number of bubbles per unit volume, then we can calculate the average bubble radius by Eq. (2.13) and use $r=r_o$. Unfortunately, the number of bubbles generally does not remain constant in a dynamic flow environment because bubbles larger than a certain size will break up. The maximum stable bubble radius r_w can be estimated in terms of a critical Weber number W_c ,

$$r_w = \frac{\sigma W_c}{2\rho_o^o U^2} , \quad (2.16)$$

where σ is the interfacial surface tension. The value of W_c is often taken as 4 for turbulent flow conditions.¹³ Thus, we define r as equal to the minimum of r_o and r_w and reserve N as an input parameter that defines the initial number of nucleating sites per unit volume (or more correctly, the minimum number of bubbles).

Finally, the relative speed U could simply be set equal to the magnitude of the average relative speed $|\underline{u}_r|$ between phases, but this would not account for local turbulent fluctuations that have been averaged out in the definition of \underline{u}_r . Fluctuations in \underline{u}_r can locally strip away the individual bubble thermal boundary layers and break up large bubbles. To account for such local effects we define

$$U = |\underline{u}_r| + \beta |\underline{u}| , \quad (2.17)$$

where \underline{u} is the mass averaged mixture velocity and β is a parameter. The β term accounts for turbulent fluctuations. We might expect β to have a magnitude of 0.1 or less because large turbulent velocity fluctuations are often observed to have magnitudes as large as 10% of the mean velocities. In general, the best value of β must be determined by comparisons with experimental data.

Again, it should be stressed that the vapor generation rate described above is preliminary and needs to be critically tested against a wide variety of situations before it can be recommended for general use. Nevertheless, this model does embrace, as special cases, the models used by many other investigators and it has produced good results in several different applications.

G. Pipe Friction

In one-dimensional applications involving flow in pipes it is desirable to have a model for pipe wall friction. SOLA-DF contains a subroutine, PFRIC for this purpose. This subroutine is only called when the one-dimensional option (DIM = 1) is set in the input data list, and only when the pipe radius, RPIPE, is initialized to a nonzero value.

Flow losses arising from wall friction affect the momentum and energy of the flow through the terms f_{vis} and W_{vis} in Eqs. (2.3) and (2.4), respectively. The pipe friction model f_{vis} contained in PFRIC is based on the Armand two-phase flow friction multiplier and uses the Colebrook approximation for the single-phase friction factor

$$f_{vis} = - \frac{f}{R} \frac{\rho_o}{\rho_l} (1-\psi)^2 \phi_{TP} \rho_l^o u^2 . \quad (2.18)$$

Here the friction coefficient f depends on the relative roughness (k/R) and the Reynolds number $Re = 2uR/v_l$,

$$f = a + b Re^{-c} , \quad (2.19)$$

where

$$a = 0.026 (k/2R)^{0.225} + 0.133 (k/2R) ,$$

$$b = 22.0 (k/2R)^{0.44} ,$$

$$c = 1.62 (k/2R)^{0.134} ,$$

and R is the hydraulic radius. The quantity ϕ_{TP} is a two-phase friction multiplier

$$\phi_{TP}^2 = (1-\theta)^{-1.75} ,$$

and ψ accounts for the effects of relative velocity

$$\psi = \rho_v [1 + (\rho - \rho_v) u_r / \rho u]^2 / \rho .$$

The value of W_{vis} is determined from the rate of change of the fluid kinetic energy associated with the f_{vis} flow loss.

III. NUMERICAL SOLUTION METHOD

Numerous schemes can be devised to numerically solve the equations outlined in the previous section. Different schemes will have varying degrees of accuracy, numerical stability, programming simplicity, flexibility for change, and computational efficiency. Unfortunately, these desirable traits are often mutually exclusive. For example, the use of implicit difference equations to achieve unconditional numerical stability can also result in poor accuracy and generally requires more complex programming and more computer memory. Because different applications require different mixtures of the desirable features, the choice of an optimum solution algorithm can rarely be made. Thus, the choice of a numerical solution procedure generally requires a balance to be made primarily between programming simplicity and the flexibility for future evolution versus stability, accuracy, and computational speed. Inevitably, the choice rests on the particular experience and prejudices of the developer.

In the SOLA-DF code an attempt has been made to keep the programming simple and to use a limited implicitness. In all cases, point relaxation methods have been retained in place of direct solvers for coupled sets of equations.

Although point relaxation methods are generally recognized as simple, but inferior to direct methods for linear equation systems, this is not necessarily the case for nonlinear equations where iterative methods must be used anyway. The point relaxation method permits considerable latitude for adding new features, changing boundary conditions, varying time steps, and making other gross changes in the basic code to adapt it to new applications.

Additionally, the code has been written in a modular form consisting of numerous subroutines that isolate individual logical and physical processes. This structure makes the present code particularly easy to modify and extend for new applications.

A. Mesh Construction

The finite-difference mesh used for numerical solution of the above equations consists of rectangular cells of width δx and height δy . The part of the mesh that contains fluid is composed of IBAR cells in the x-direction, labeled with the index i, and JBAR cells in the y-direction, labeled with the index j. This region is surrounded by a single layer of fictitious cells so that the complete mesh is $IMAX = IBAR + 2$ by $JMAX = JBAR + 2$ cells (see Fig. 1). The dependent variables are located within a cell as shown in Fig. 2: x-directed velocities at the middle of the vertical sides, y-directed velocities at the middle of the horizontal sides, and all other variables at the cell center.

Finite-difference subscripting used in the computer code is based on the convention that (i,j) refers to the center of the cell labeled by (i,j) or to the right or top boundary of the cell in the case of velocities, which are located at those cell boundaries.

B. Solution Algorithm

A time cycle of calculation is broken down into four tasks. First, the momentum equations, Eq. (2.3), are advanced explicitly using the previous cycle values for evaluating all contributions. Next an iteration is undertaken to replace the pressure used in the first task with advanced time values. An iteration is needed because the advanced pressures depend on the velocities being calculated. This part of the cycle contains the main implicitness of the numerical scheme. The pressure iteration permits sound waves to propagate more than one mesh cell per cycle. In fact, this scheme is a variant of the ICE technique,¹⁴ which may be used for very low speed (incompressible) flows as well as for high-speed flows containing shock waves and rarefactions. The third

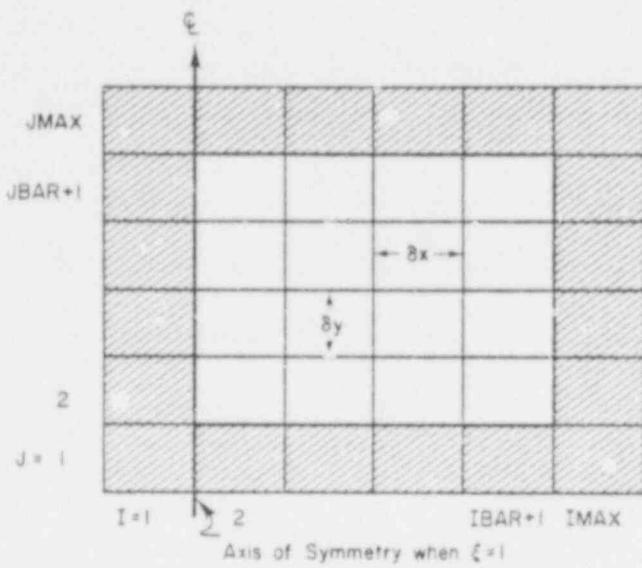


Fig. 1. General mesh arrangement.
Fictitious boundary cells are shaded.

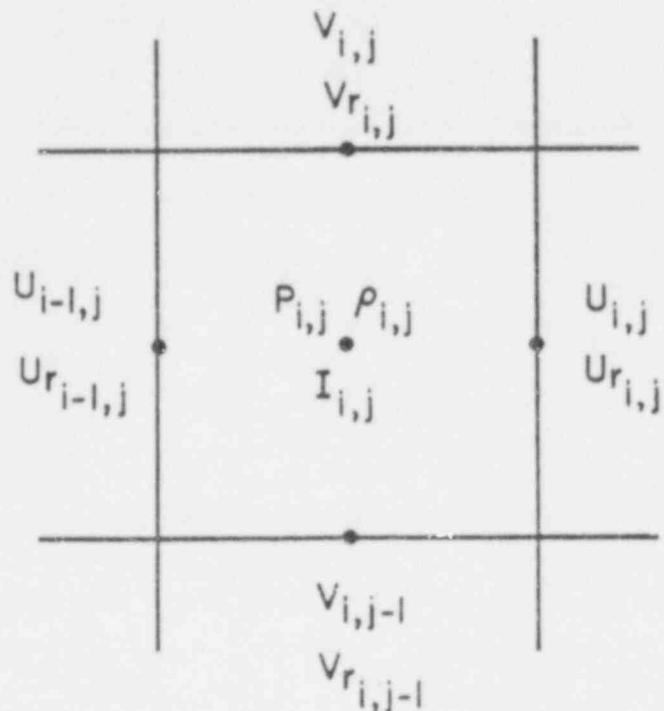


Fig. 2. Arrangement of finite difference variables in a typical mesh cell.

task in a cycle is to update all other dependent variables. Finally, the fourth task consists of data output, time step controls, and bookkeeping operations.

For a purely explicit calculation, the iteration making up the second task may be omitted by setting the input number IMP equal to zero.

1. Explicit Updating of Velocities. Before introducing finite difference approximations for the momentum equation, Eq. (2.3), it is first written in the equivalent differential form,

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + \frac{1}{\rho A} \frac{\partial}{\partial x} \left(\frac{A \rho v^2}{\sigma} - u_r^2 \right) + v \frac{\partial u}{\partial y} + \frac{1}{\rho A} \frac{\partial}{\partial y} \left(\frac{A \rho v^2}{\sigma} - u_r v_r \right) \\ = - \frac{1}{\rho} \frac{\partial p}{\partial x} + g_x + \frac{1}{\rho} f_{visx} \end{aligned}$$

$$\begin{aligned} \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + \frac{1}{\rho A} \frac{\partial}{\partial x} \left(\frac{A \rho v^2}{\rho} u_r v_r \right) + v \frac{\partial v}{\partial y} + \frac{1}{\rho A} \frac{\partial}{\partial y} \left(\frac{A \rho v^2}{\rho} v_r^2 \right) \\ = - \frac{1}{\rho} \frac{\partial p}{\partial y} + g_y + \frac{1}{\rho} f_{visy} \quad . \end{aligned} \quad (3.1)$$

This form is a carryover from previous codes in the SOLA series. Its advantage is that u^{n+1} and v^{n+1} are calculated directly rather than $(\rho u)^{n+1}$ and $(\rho v)^{n+1}$. The disadvantage of this equation is that it is not in conservation form so we do not get rigorous conservation of momentum in the difference approximation. When percentage changes in dependent variables from one cell to the next are not large this nonconservation form of the momentum equations should not cause any problems. In general, it would probably be a good idea to monitor the total momentum of the system in order to have a check on the accuracy of Eq. (3.1), but this has not been included in the present version of the code. The difference equations used to approximate Eq. (3.1) are

$$\begin{aligned} \tilde{u}_{i,j} &= u_{i,j} + \delta t \left[\frac{2(p_{i,j} - p_{i+1,j})}{\delta x(\rho_{i,j} + \rho_{i+1,j})} + g_x \right. \\ &\quad \left. + f_{visx}_{i,j} - FUX - FUY - FDU \right] \\ \tilde{v}_{i,j} &= v_{i,j} + \delta t \left[\frac{2(p_{i,j} - p_{i,j+1})}{\delta y(\rho_{i,j} + \rho_{i,j+1})} + g_y \right. \\ &\quad \left. + f_{visy}_{i,j} - FVX - FVY - FDV \right] , \end{aligned} \quad (3.2)$$

where $\tilde{u}_{i,j}$ and $\tilde{v}_{i,j}$ are the explicit guesses for $u_{i,j}^{n+1}$ and $v_{i,j}^{n+1}$, respectively. In all cases, the boundary values of cell-centered quantities will be defined by linear interpolation, as for ρ used above, unless otherwise noted. The indexing in this equation looks uncentered but recall from Fig. 2 that $u_{i,j}$ and $v_{i,j}$ refer to the velocities at the boundary between cells i and $i+1$ and j and $j+1$, respectively. The convective fluxes are defined as

532 222

$$\begin{aligned}
 F_{UX} &= \frac{1}{2\delta x} [u_{i,j}(u_{i+1,j} - u_{i-1,j}) \\
 &\quad - \alpha|u_{i,j}|(u_{i+1,j} - 2u_{i,j} + u_{i-1,j})] \\
 F_{UY} &= \frac{1}{8\delta y} [(v_{i,j} + v_{i+1,j} + v_{i,j-1} + v_{i+1,j-1})(u_{i,j+1} - u_{i,j-1}) \\
 &\quad - \alpha|v_{i,j} + v_{i+1,j} + v_{i,j-1} + v_{i+1,j-1}|(u_{i,j+1} - 2u_{i,j} + u_{i,j-1})]
 \end{aligned}$$

The parameter α appearing in the convective fluxes is used to give a variable amount of upstream differencing. When α is zero the approximations reduce to the usual centered differenced form; however, this is known to result in an unstable algorithm.¹⁵ When α is unity the approximations are the so-called donor cell or fully upstream (or upwind) difference expressions, which are stable provided fluid does not convect through more than one mesh cell in one time step. In general, numerical stability is expected (see Ref. 15) when α is chosen such that

$$\alpha > \max \left[\frac{u \delta t}{\delta x}, \frac{v \delta t}{\delta y}, \frac{\rho_v \rho_L u_r \delta t}{\rho^2 \delta x}, \frac{\rho_v \rho_L v_r \delta t}{\rho^2 \delta y} \right].$$

The remaining flux terms are

$$\begin{aligned}
 F_{DU} &= \frac{1}{A_R} \left\{ \frac{1}{\delta x} \left[(RUR) \left(u_{r_{i,j}} + u_{r_{i+1,j}} \right) + \alpha |(RUR)| \left(u_{r_{i,j}} - u_{r_{i+1,j}} \right) \right. \right. \\
 &\quad \left. \left. - (RUL) \left(u_{r_{i-1,j}} + u_{r_{i,j}} \right) - \alpha |(RUL)| \left(u_{r_{i-1,j}} - u_{r_{i,j}} \right) \right] \right. \\
 &\quad \left. + \frac{1}{\delta y} \left[(VDT) (RUC - RUT) + \alpha |(VDT)| (RUC - RUT) - (VDB) (RUB + RUC) \right. \right. \\
 &\quad \left. \left. - \alpha |(VDB)| (RUB - RUC) \right] \right\} / (\rho_{i,j} + \rho_{i+1,j}),
 \end{aligned}$$

$$\begin{aligned}
FVX &= \frac{1}{2\delta x} [(u_{i-1,j+1} + u_{i,j+1} + u_{i-1,j} + u_{i,j})(v_{i+1,j} + v_{i-1,j}) \\
&\quad - \alpha |u_{i-1,j+1} + u_{i,j+1} + u_{i-1,j} + u_{i,j}| (v_{i+1,j} - 2v_{i,j} \\
&\quad + v_{i-1,j})] , \\
FVY &= \frac{1}{2\delta y} [v_{i,j}(v_{i,j+1} - v_{i,j-1}) - \alpha |v_{i,j}| (v_{i,j+1} - 2v_{i,j} + v_{i,j-1})] , \\
FDV &= \frac{1}{A_T} \left\{ \frac{1}{\delta x} \left[(UDR)(RVC + RVR) + \alpha |(UDR)| (RVC - RVR) - (UDL)(RVL + RVC) \right. \right. \\
&\quad - \alpha |UDL| (RVL - RVC)] + \frac{1}{\delta y} \left[RVT \left(v_{r_{i,j}} + v_{r_{i,j+1}} \right) + \alpha |(RVT)| \left(v_{r_{i,j}} \right. \right. \\
&\quad - \left. v_{r_{i,j+1}} \right) - (RVB) \left(v_{r_{i,j-1}} + v_{r_{i,j}} \right) - \alpha |(RVB)| \left(v_{r_{i,j-1}} \right. \\
&\quad \left. \left. - v_{r_{i,j}} \right) \right] \left. \right\} / (\rho_{i,j} + \rho_{i,j+1}) ,
\end{aligned}$$

where

$$\begin{aligned}
A_R &= \frac{2 A_{i,j} A_{i+1,j}}{A_{i,j} + A_{i+1,j}} , \\
RUR &= \frac{\rho_{v_{i+1,j}} \rho_{l_{i+1,j}}}{\rho_{i+1,j}} \left(\frac{u_{r_{i,j}} + u_{r_{i+1,j}}}{2} \right) A_{i+1,j} , \\
RUL &= \frac{\rho_{v_{i,j}} \rho_{l_{i,j}}}{\rho_{i,j}} \left(\frac{u_{r_{i-1,j}} + u_{r_{i,j}}}{2} \right) A_{i,j} , \\
VDT &= \left(\frac{v_{r_{i,j}} + v_{r_{i+1,j}}}{2} \right) \left(\frac{A_{i,j} A_{i,j+1}}{A_{i,j} + A_{i,j+1}} + \frac{A_{i+1,j} A_{i+1,j+1}}{A_{i+1,j} + A_{i+1,j+1}} \right) ,
\end{aligned}$$

$$RUC = \frac{u_{r_{i,j}}}{2} \left(\frac{\rho_{v_{i,j}} \rho_{\ell_{i,j}}}{\rho_{i,j}} + \frac{\rho_{v_{i+1,j}} \rho_{\ell_{i+1,j}}}{\rho_{i+1,j}} \right) ,$$

$$RUT = \frac{u_{r_{i,j+1}}}{2} \left(\frac{\rho_{v_{i,j+1}} \rho_{\ell_{i,j+1}}}{\rho_{i,j+1}} + \frac{\rho_{v_{i+1,j+1}} \rho_{\ell_{i+1,j+1}}}{\rho_{i+1,j+1}} \right) ,$$

$$VDB = \left(\frac{v_{r_{i,j-1}} + v_{r_{i+1,j-1}}}{2} \right) \left(\frac{\rho_{v_{i,j-1}} A_{i,j-1}}{\rho_{i,j} + A_{i,j-1}} + \frac{A_{i+1,j-1} A_{i+1,j-1}}{A_{i+1,j} A_{i+1,j-1}} \right) ,$$

$$RUB = \frac{u_{r_{i+1,j-1}}}{2} \left(\frac{\rho_{v_{i+1,j-1}} \rho_{\ell_{i+1,j-1}}}{\rho_{i+1,j-1}} + \frac{\rho_{v_{i+1,j-1}} \rho_{\ell_{i+1,j-1}}}{\rho_{i+1,j-1}} \right) ,$$

$$A_T = \frac{2 A_{i,j} A_{i,j+1}}{A_{i,j} + A_{i,j+1}} ,$$

$$UDR = \left(\frac{u_{r_{i,j}} + u_{r_{i,j+1}}}{2} \right) \left(\frac{A_{i,j} A_{i+1,j}}{A_{i,j} + A_{i+1,j}} + \frac{A_{i,j+1} A_{i+1,j+1}}{A_{i,j+1} + A_{i+1,j+1}} \right) ,$$

$$RVC = \frac{v_{r_{i,j}}}{2} \left(\frac{\rho_{v_{i,j}} \rho_{\ell_{i,j}}}{\rho_{i,j}} + \frac{\rho_{v_{i,j+1}} \rho_{\ell_{i,j+1}}}{\rho_{i,j+1}} \right) ,$$

$$RVR = \frac{v_{r_{i+1,j}}}{2} \left(\frac{\rho_{v_{i+1,j}} \rho_{\ell_{i+1,j}}}{\rho_{i+1,j}} + \frac{\rho_{v_{i+1,j+1}} \rho_{\ell_{i+1,j+1}}}{\rho_{i+1,j+1}} \right) ,$$

$$UDL = \left(\frac{u_{r_{i-1,j}} + u_{r_{i-1,j+1}}}{2} \right) \left(\frac{A_{i,j} A_{i-1,j}}{A_{i,j} + A_{i-1,j}} + \frac{A_{i,j+1} A_{i-1,j+1}}{A_{i,j+1} + A_{i-1,j+1}} \right) ,$$

$$RVL = \frac{v_{r_{i-1,j}}}{2} \left(\frac{\rho_{v_{i-1,j}} \rho_{\ell_{i-1,j}}}{\rho_{i-1,j}} + \frac{\rho_{v_{i-1,j+1}} \rho_{\ell_{i-1,j+1}}}{\rho_{i-1,j+1}} \right) ,$$

$$RVT = \frac{\rho_{v,i,j+1} \rho_{l,i,j+1}}{\rho_{i,j+1}} \left(\frac{v_{r,i,j} + v_{r,i,j+1}}{2} \right) A_{i,j+1}, \text{ and}$$

$$RVB = \frac{\rho_{v,i,j} \rho_{l,i,j}}{\rho_{i,j}} \left(\frac{v_{r,i,j-1} + v_{r,i,j}}{2} \right) A_{i,j}.$$

For the cell boundary areas, we do not use a linearly interpolated value, but instead, prefer a combined geometric and arithmetic average, e.g.,

$$A_{i,j+\frac{1}{2}} = \frac{2 A_{i,j} A_{i,j+1}}{(A_{i,j} + A_{i,j+1})}.$$

This choice has the convenient feature that $A_{i,j+\frac{1}{2}}$ vanishes when either $A_{i,j}$ or $A_{i,j+1}$ is zero.

2. Implicit Pressure Calculation. In this part of the calculational cycle the n level pressures appearing in Eq. (3.2) are to be replaced by approximations for the $n+1$ level pressures. This is done by solving for the pressure in each cell that satisfies the implicit equation

$$F = p - f(\bar{\rho}, \bar{\rho}_v, \bar{I}) = 0, \quad (3.3)$$

where $f(\bar{\rho}, \bar{\rho}_v, \bar{I})$ is the equation of state evaluated using

$$\bar{\rho} = \rho^n / (1 + D),$$

$$\bar{\rho}_v = \rho_v^n / (1 + D),$$

$$\bar{I} = I^n - \frac{p^n}{\rho} D , \quad (3.4)$$

and D is an approximation to cell volume change per unit volume,

$$D = \frac{\delta t}{A} \left[\frac{\partial}{\partial x} (A u) + \frac{\partial}{\partial y} (A v) \right]$$

$$= \frac{\delta t}{A} \left[\frac{1}{\delta x} (A_{i+\frac{1}{2},j} u_{i,j} - A_{i-\frac{1}{2},j} u_{i-1,j}) \right. \\ \left. + \frac{1}{\delta y} (A_{i,j+\frac{1}{2}} v_{i,j} - A_{i,j-\frac{1}{2}} v_{i,j-1}) \right] . \quad (3.5)$$

The $n+1$ level velocities must be used in evaluating D ; that is,

$$u_{i,j}^{n+1} = \tilde{u}_{i,j} - \frac{2\delta t}{(\rho_{i+1,j} + \rho_{i,j})} (p_{i+1,j} - p_{i+1,j}^n - p_{i,j} + p_{i,j}^n) , \quad (3.6)$$

and

$$v_{i,j}^{n+1} = \tilde{v}_{i,j} - \frac{2\delta t}{(\rho_{i,j+1} + \rho_{i,j})} (p_{i,j+1} - p_{i,j+1}^n - p_{i,j} + p_{i,j}^n) . \quad (3.7)$$

Because the $n+1$ level velocities depend on p the implicit nature of Eq. (3.3) is obvious. The pressure that satisfies Eqs. (3.3)-(3.7) is not quite p^{n+1} , because convective fluxes are omitted from the estimates of the new densities $\tilde{\rho}$, $\tilde{\rho}_v$, and energy \bar{I} . It would be equal to p^{n+1} if the remainder of the equations were in Lagrangian form. The difference is not significant, however, because the iteration is always trying to drive p to its equation-of-state value every cycle. In this sense p is a stored variable and is not identically equal to the equation-of-state value unless the iterations are omitted and an explicit

calculation is used. It is important to note that densities and energies are not actually changed during the iteration, because Eq. (3.4) is only used to estimate the new values. To obtain a solution of Eqs. (3.3)-(3.7) a local Newton-Raphson procedure is followed. For this purpose an estimate is needed for $\partial F/\partial p$ in each cell. In SOLA-DF, first guesses for these values are computed by a numerical differentiation. This is done at the beginning of each cycle and the values stored. Once the iteration has started, new values are computed and stored after each iteration. In summary, the following steps are performed for a single cell (i,j):

- a. compute D according to Eq. (3.5) using the most updated values of u and v from Eqs. (3.6) and (3.7),
- b. compute $\bar{\rho}$, $\bar{\rho}_v$, and \bar{I} from Eq. (3.4),
- c. evaluate the equation-of-state function and calculate $\delta p = -F/\partial F/\partial p$,
- d. replace $p_{i,j}$ by $p_{i,j} + \delta p_{i,j}$, and $u_{i,j}$, $u_{i-1,j}$, $v_{i,j}$, and $v_{i,j-1}$ by

$$u_{i,j} = u_{i,j} + \frac{2 \delta t \delta p}{\delta x(p_{i,j} + p_{i+1,j})} ,$$

$$u_{i-1,j} = u_{i-1,j} - \frac{2 \delta t \delta p}{\delta x(p_{i-1,j} + p_{i,j})} ,$$

$$v_{i,j} = v_{i,j} + \frac{2 \delta t \delta p}{\delta y(p_{i,j} + p_{i,j+1})} ,$$

and

$$v_{i,j-1} = v_{i,j-1} - \frac{2 \delta t \delta p}{\delta y(p_{i,j-1} + p_{i,j})} .$$

This iteration process is continued until all cells satisfy the convergence test

$$\left| \frac{p - f(\bar{\rho}, \bar{\rho}_v; \bar{I})}{p + f(\bar{\rho}, \bar{\rho}_v; \bar{I})} \right| < \epsilon ,$$

532 228

where ϵ is typically equal to 0.001.

3. Updating of Remaining Variables. After the implicit portion of the cycle has been completed new time values for the remaining variables are readily computed. The mixture density changes only by convection,

$$\begin{aligned} \rho_{i,j}^{n+1} = & \rho_{i,j} - \frac{\delta t}{A_{i,j}} \left\{ \frac{1}{\delta x} \left[A_{i+\frac{1}{2},j} (\rho u)_{i+\frac{1}{2},j} - A_{i-\frac{1}{2},j} (\rho u)_{i-\frac{1}{2},j} \right] \right. \\ & \left. + \frac{1}{\delta y} \left[A_{i,j+\frac{1}{2}} (\rho v)_{i,j+\frac{1}{2}} - A_{i,j-\frac{1}{2}} (\rho v)_{i,j-\frac{1}{2}} \right] \right\} , \end{aligned} \quad (3.8)$$

where

$$\begin{aligned} (\rho u)_{i+\frac{1}{2},j} = & \frac{1}{2} \left[u_{i,j} (\rho_{i,j} + \rho_{i+1,j}) + \alpha |u_{i,j}| (\rho_{i,j} - \rho_{i+1,j}) \right] , \\ (\rho u)_{i-\frac{1}{2},j} = & \frac{1}{2} \left[u_{i-1,j} (\rho_{i-1,j} + \rho_{i,j}) + \alpha |u_{i-1,j}| (\rho_{i-1,j} - \rho_{i,j}) \right] , \\ (\rho v)_{i,j+\frac{1}{2}} = & \frac{1}{2} \left[v_{i,j} (\rho_{i,j} + \rho_{i,j+1}) + \alpha |v_{i,j}| (\rho_{i,j} - \rho_{i,j+1}) \right] , \end{aligned}$$

and

$$(\rho v)_{i,j-\frac{1}{2}} = \frac{1}{2} \left[v_{i,j-1} (\rho_{i,j} + \rho_{i,j-1}) + \alpha |v_{i,j-1}| (\rho_{i,j} - \rho_{i,j-1}) \right] .$$

Quantities on the right side of Eq. (3.8) are evaluated using n level values for ρ and the available $n+1$ level values of u .

The mixture energy equation is next approximated by

$$\begin{aligned} I_{i,j}^{n+1} = & \frac{\rho_{i,j} I_{i,j}}{\rho_{i,j}^{n+1}} + \frac{\delta t}{\rho_{i,j}^{n+1}} \left\{ - FU - FV - FWK \right. \\ & \left. + \frac{1}{4} K_{i,j} \left[\left(u_{r_{i,j}} + u_{r_{i-1,j}} \right)^2 + \left(v_{r_{i,j}} + v_{r_{i,j-1}} \right)^2 \right] + \left(w_{vis_{i,j}} \right) \right\} , \end{aligned} \quad (3.9)$$

where

$$FU = \frac{1}{A_{i,j} \delta x} \left\{ A_{i+\frac{1}{2},j} \left[(\rho Iu)_{i+\frac{1}{2},j} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) u_r \right)_{i+\frac{1}{2},j} \right] \right.$$

$$\left. - A_{i-\frac{1}{2},j} \left[(\rho Iu)_{i-\frac{1}{2},j} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) u_r \right)_{i-\frac{1}{2},j} \right] \right\} ,$$

$$FV = \frac{1}{A_{i,j} \delta y} \left\{ A_{i,j+\frac{1}{2}} \left[(\rho Iv)_{i,j+\frac{1}{2}} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) v_r \right)_{i,j+\frac{1}{2}} \right] \right.$$

$$\left. - A_{i,j-\frac{1}{2}} \left[(\rho Iv)_{i,j-\frac{1}{2}} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) v_r \right)_{i,j-\frac{1}{2}} \right] \right\} .$$

For this formulation the boundary convective fluxes are approximated by

$$\left[(\rho Iu)_{i+\frac{1}{2},j} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) u_r \right) \right]_{i+\frac{1}{2},j}$$

$$= \frac{1}{2} \left\{ u_{i,j} \left[(\rho I)_{i,j} + (\rho I)_{i+1,j} \right] + \alpha |u_{i,j}| \left[(\rho I)_{i,j} \right. \right.$$

$$\left. - (\rho I)_{i+1,j} \right] + u_r \left[\left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) \right)_{i,j} \right]$$

$$+ \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) \right)_{i+1,j} \right]$$

$$+ \alpha |u_r| \left[\left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) \right)_{i,j} - \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) \right)_{i+1,j} \right] \right\} ,$$

and similarly for $\left[(\rho I u)_{i-\frac{1}{2}, j} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) u_r \right)_{i-\frac{1}{2}, j} \right]$,

$$\left[(\rho I v)_{i, j+\frac{1}{2}} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) v_r \right)_{i, j+\frac{1}{2}} \right], \text{ and } \left[(\rho I v)_{i, j-\frac{1}{2}} + \left(\frac{\rho_v \rho_l}{\rho} (I_v - I_l) v_r \right)_{i, j-\frac{1}{2}} \right].$$

The pressure work term in Eq. (3.9) is approximated as

$$\begin{aligned}
 F_{WK} = & \frac{p_{i,j}}{A_{i,j}} \left\{ \frac{1}{\delta x} \left[A_{i+\frac{1}{2},j} \left(u_{i,j} + \frac{\rho_v p_{i+\frac{1}{2},j} \rho_l p_{i+\frac{1}{2},j}}{\rho_{i+\frac{1}{2},j}} \left(\frac{1}{\rho_v^o p_{i+\frac{1}{2},j}} \right. \right. \right. \right. \right. \\
 & \left. \left. \left. \left. \left. \left. - \frac{1}{\rho_l^o p_{i+\frac{1}{2},j}} \right) u_r \right) - A_{i-\frac{1}{2},j} \left(u_{i-1,j} \right. \right. \right. \right. \\
 & \left. \left. \left. \left. \left. \left. + \frac{\rho_v p_{i-\frac{1}{2},j} \rho_l p_{i-\frac{1}{2},j}}{\rho_{i-\frac{1}{2},j}} \left(\frac{1}{\rho_v^o p_{i-\frac{1}{2},j}} - \frac{1}{\rho_l^o p_{i-\frac{1}{2},j}} \right) u_r \right) \right] \right] \right] \\
 & + \frac{1}{\delta y} \left[A_{i,j+\frac{1}{2}} \left(v_{i,j} + \frac{\rho_v p_{i,j+\frac{1}{2}} \rho_l p_{i,j+\frac{1}{2}}}{\rho_{i,j+\frac{1}{2}}} \left(\frac{1}{\rho_v^o p_{i,j+\frac{1}{2}}} \right. \right. \right. \right. \\
 & \left. \left. \left. \left. \left. - \frac{1}{\rho_l^o p_{i,j+\frac{1}{2}}} \right) v_r \right) - A_{i,j-\frac{1}{2}} \left(v_{i,j-1} \right. \right. \right. \right. \\
 & \left. \left. \left. \left. \left. + \frac{\rho_v p_{i,j-\frac{1}{2}} \rho_l p_{i,j-\frac{1}{2}}}{\rho_{i,j-\frac{1}{2}}} \left(\frac{1}{\rho_v^o p_{i,j-\frac{1}{2}}} - \frac{1}{\rho_l^o p_{i,j-\frac{1}{2}}} \right) v_r \right) \right] \right] \right\}
 \end{aligned}$$

Quantities on the right side of Eq. (3.9) are evaluated using n level values for ρ , ρ_v , and I while $n+1$ level values are used for $u_{i,j}$, $p_{i,j}$, and the overall divisor $\rho_{i,j}$. Finally, the vapor density is updated as

$$(\rho_v)_{i,j}^{n+1} = (\rho_v)_{i,j} + \delta t [- FRU + \Gamma_{i,j}] , \quad (3.10)$$

where

$$\begin{aligned} FRU = & \frac{1}{A_{i,j}} \left\{ \frac{1}{\delta x} \left[A_{i+\frac{1}{2},j} (\rho_v u_v)_{i+\frac{1}{2},j} - A_{i-\frac{1}{2},j} (\rho_v u_v)_{i-\frac{1}{2},j} \right] \right. \\ & \left. + \frac{1}{\delta y} \left[A_{i,j+\frac{1}{2}} (\rho_v v_v)_{i,j+\frac{1}{2}} - A_{i,j-\frac{1}{2}} (\rho_v v_v)_{i,j-\frac{1}{2}} \right] \right\} \end{aligned}$$

and

$$\begin{aligned} (\rho_v u_v)_{i+\frac{1}{2},j} = & \frac{1}{2} \left\{ (u_v)_{i,j} \left[(\rho_v)_{i,j} + (\rho_v)_{i+1,j} \right] \right. \\ & \left. + \alpha |u_v|_{i,j} \left[(\rho_v)_{i,j} - (\rho_v)_{i+1,j} \right] \right\} . \end{aligned}$$

The $i-\frac{1}{2}$ boundary flux is obtained by replacing i with $i-1$ in the above expression. In a similar fashion,

$$\begin{aligned} (\rho_v v_v)_{i,j+\frac{1}{2}} = & \frac{1}{2} \left\{ (v_v)_{i,j} \left[(\rho_v)_{i,j} + (\rho_v)_{i,j+1} \right] \right. \\ & \left. + \alpha |v_v|_{i,j} \left[(\rho_v)_{i,j} - (\rho_v)_{i,j+1} \right] \right\} . \end{aligned}$$

where the $j-\frac{1}{2}$ boundary flux is obtained by replacing j with $j-1$ in the above equation. The vapor velocities used in these expressions are defined as

$$(u_v)_{i,j} = u_{i,j} + (\rho_l)_{i+\frac{1}{2},j} (u_r)_{i,j} / \rho_{i+\frac{1}{2},j} ,$$

with a corresponding expression for the $i-\frac{1}{2}$ boundary obtained by replacing i with $i-1$, and

$$(v_v)_{i,j} = v_{i,j} + (\rho_l)_{i,j+\frac{1}{2}} (v_r)_{i,j} / \rho_{i+\frac{1}{2},j} ,$$

with a corresponding expression for the $j-\frac{1}{2}$ boundary obtained by replacing j with $j-1$.

Some care must be taken with the way the vapor source term Γ , is approximated. When the mixture is not at equilibrium and the relaxation rate is fast, Γ can be large. Under such circumstances Eq. (3.10) may be numerically unstable. To avoid this the densities in Γ should be evaluated at level $n+1$. For general formulations of Γ it is usually necessary to use an iterative technique to solve Eq. (3.10). Such an iteration is provided in subroutine PHCHR for the Γ given by Eq. (2.14). There is, however, another more serious problem that can arise when phase transitions are important. Because the effect of Γ is included at the end of a cycle of calculation, its influence on the pressure, and hence the dynamics, is not accounted for in the implicit pressure iteration. This means that some inaccuracies can be introduced in the propagation of compression and rarefaction waves when significant phase change occurs during a single time step. In addition, a large phase change may also drive the equation-of-state pressure far from the value arrived at in the pressure iteration; so that excessive iterations may be required to solve the implicit equation in the next time cycle. In extreme cases the iteration may not even converge. The above problem can be eliminated by using sufficiently small time increments, δt , but in some cases this may lead to long computing times. A better solution that has been utilized recently in the multidimensional code K-FIX (see Ref. 2) is to incorporate the Γ into the implicit portion of the cycle. Basically, the idea is to include Γ in

Eq. (3.4) for the estimated new time vapor density. Since this more complicated formulation is not in SOLA-DF, the user should check his results for time step dependence (accuracy) by performing a smaller time step calculation when necessary.

Thermodynamic equilibrium calculations can be achieved by using a large phase change rate or by replacing the vapor density equation with a calculation of the saturated vapor density and using an equilibrium equation of state. The latter procedure is often preferable because it effectively puts Γ into the pressure iteration.

C. Boundary Conditions

Various types of boundary conditions may be used at the ends and sides of the computational mesh. Prescribed velocities or prescribed pressures together with densities and temperatures may be used to represent inlet and exit conditions. For example, a guillotine break in a pipe of a reactor system can be represented by assigning the ambient pressure in the containment structure to the end of the mesh.

These boundary conditions are easily imposed by setting appropriate values of the dependent variables in the fictitious cells surrounding the mesh. Five kinds of boundary conditions have been specifically built into the code: (1) rigid free-slip, (2) rigid no-slip, (3) continuative outflow, (4) periodic, and (5) constant pressure. As an example of how to treat these boundary conditions, consider the left boundary.

1. Rigid Free-Slip. The normal velocity component must be zero and the tangential velocity component must have no normal gradient, i.e.,

$$\left. \begin{array}{l} u_{1,j} = 0 \\ v_{1,j} = v_{2,j} \end{array} \right\} \text{for all } j.$$

The variables ρ , p , and I are treated the same as v .

2. Rigid No-Slip. The normal velocity component must be zero and the tangential velocity component at the wall must also be zero, i.e.,

$$\left. \begin{array}{l} u_{1,j} = 0 \\ v_{1,j} = -v_{2,j} \end{array} \right\} \text{for all } j.$$

The variables ρ , p , and I are treated the same as for a free-slip boundary.

For both of the above rigid boundary conditions, the designated conditions are imposed after each pass through the mesh during the pressure iteration.

3. Continuative Outflow. Continuative or outflow boundaries always pose a problem for low-speed calculations, because whatever prescription is chosen it can potentially affect the entire flow field upstream. What is needed is a prescription that permits fluid to flow out of the mesh with a minimum of upstream influence. In this code we have used a continuative boundary condition that involves setting, for the left boundary, for example,

$$\left. \begin{array}{l} u_{1,j} = u_{2,j} \\ v_{1,j} = v_{2,j} \end{array} \right\} \text{for all } j.$$

These conditions, however, are only imposed after applying the complete momentum equations (3.2) and not after each pass through the mesh during the pressure iteration. During the iteration the normal boundary velocities can vary with the changes in pressure, as any interior velocity component. The treatment of ρ , p , and I is the same as (u, v) .

4. Periodic. For periodic boundary conditions in the x -direction, the left and right boundaries must be set to reflect the periodicity. This is easiest when the period length is chosen equal to $(IBAR-1)\delta x$. Then the boundary condition for the fictitious cells on the left are

$$\left. \begin{array}{l} u_{1,j} = u_{IBAR,j} \\ v_{1,j} = v_{IBAR,j} \\ \rho_{1,j} = \rho_{IBAR,j} \\ p_{1,j} = p_{IBAR,j} \\ I_{1,j} = I_{IBAR,j} \end{array} \right\} \text{for all } j.$$

and on the right

$$\left. \begin{array}{l} u_{IBAR+1,j} = u_{2,j} \\ v_{IBAR+2,j} = v_{3,j} \\ \rho_{IBAR+2,j} = \rho_{3,j} \\ p_{IBAR+2,j} = p_{3,j} \\ I_{IBAR+2,j} = I_{3,j} \end{array} \right\} \text{for all } j.$$

In this case these conditions are imposed after applying Eqs. (3.2) and after each pressure iteration.

5. Constant Pressure. When constant pressure is prescribed at a boundary, the fluid must stream freely into or out of the specified pressure region, i.e.,

$$\left. \begin{array}{l} p_{1,j} = p_{2,j} = p_{bc} \\ u_{1,j} = u_{2,j} \\ v_{1,j} = v_{2,j} \end{array} \right\} \text{for all } j.$$

The variables ρ and I are treated the same as v .

Boundary conditions similar to those for the left wall are used at the right, top, and bottom boundaries of the mesh. Of course, the normal and tangential velocities at the top and bottom boundaries are v and u , respectively.

For convenience, the SOLA-DF code has been written so that any of the above boundary conditions can be automatically imposed by setting input numbers. The appropriate input number for the left wall is designated WL, where

$$WL = \begin{cases} 1, & \text{rigid free-slip left wall} \\ 2, & \text{rigid no-slip left wall} \\ 3, & \text{continuative outflow left wall} \\ 4, & \text{periodic in } x \text{ (provided WR} = 4) \\ 5, & \text{constant pressure left wall.} \end{cases}$$

Similar input numbers are used for the right boundary (WR), top boundary (WT), and bottom boundary (WB). Clearly, when periodic conditions are desired in a given direction, both boundaries in that direction must be assigned wall numbers of 4.

D. Internal Obstacles

To increase the usefulness of the basic code, internal obstacles may be inserted within the fluid region. Internal obstacles with rigid boundary conditions can be taken into account by simply setting $A_{i,j} = 0.0$ for the desired obstacle cell. The code is equipped to handle these simple obstacles by means of the INPUT stream; however, if a more complex internal obstacle treatment is desired, the user must provide additional coding at the end of the boundary condition routine BC.

E. Variable Time Steps

SOLA-DF contains provisions that allow the use of different time steps of integration. The time steps are determined by numerical stability requirements and other user-specified conditions. For numerical stability the time step is limited by the flux criterion that $u_r \delta t / \delta x \leq 1/4$, $u_r \delta t / \delta x \leq 1/4$, $v_r \delta t / \delta y \leq 1/4$, and $v_r \delta t / \delta x \leq 1/4$. The minimum time step determined according to this criterion is then increased or decreased by 1%. The direction of this adjustment is deter-

mined by the relative ease of the previous time integration of the system. If fewer than five system iterations were required, the time step is increased, otherwise it is decreased. The time step determined by the above methods is never allowed to be greater than the user-specified maximum time step DELTMX.

IV. INPUT DATA, COMMON VARIABLES, AND SUBROUTINES AND GRAPHICS OUTPUT

The input data, COMMON variables, and subroutines and FORTRAN functions in SOLA-DF are listed and described in this section. While the input quantities are tabulated and defined separately, they also appear in the COMMON variable lists and are identified there simply as an input quantity. These descriptions of the input and COMMON variables are necessarily brief but hopefully will assist in relating the methodology, described in Sec. III, to its implementation in the code. The units the code was written for are g, cm, ms, K with pressure in bars.

A. Input Data

DEFAULT	FORTRAN	ALGEBRAIC	
VALUE	SYMBOL	SYMBOL	DEFINITION
1.0	ALPHA	α	Parameter that determines the amount of upstream differencing in the convective flux terms. Equal to one gives full donor cell differencing.
1.234×10^4	ASQ	a^2	Square of the speed of sound for the liquid phase.
10^4	BUBN	N	Representative bubble (or droplet) number density per cm^3 , used in phase change and interfacial friction model.
0.50	CDG	c_d	Drag coefficient used in the interfacial friction model.

44.34	CHL	C_L	Coefficient of the linear term in the liquid internal energy function.
6.67	CHV	C_V	Similar to CHL but for the vapor energy function.
1.0	CYL	-	Define coordinate system; cylindrical (CYL = 1.0) and Cartesian (CYL = 0.0).
1.0×10^{-4}	DELT	δt	Starting time step for the calculation in ms.
1.0×10^3	DELT MX	δt_{\max}	Maximum time step in ms.
1.0	DELX	δx	Cell length in cm - radial length in cylindrical coordinates.
1.0	DELY	δy	Cell length in cm - axial length in cylindrical coordinates.
0.0	DFVEL	-	Program control parameter that determines whether the relative velocity is to be calculated (DFVEL = 1.) or set to zero (DFVEL = 0.).
2.0	DIM	-	Problem dimension: (DIM = 2.0) is a two-dimensional problem and (DIM = 1.0) is a one-dimensional problem.
4.174×10^3	ECL	E_L	Coniant in the liquid internal energy function.

2.506×10^4	ECV	E_v	Similar to ECL but for the vapor energy function.
1.6×10^{-6}	EDL	-	Ratio of the thermal conductivity to the specific heat for the liquid.
1.6×10^{-7}	EDV	-	Ratio of the thermal conductivity to the specific heat for the vapor.
1.76×10^4	ELHT	λ	Latent heat of vaporization.
0.001	EPSI	ϵ	Pressure iteration convergence test parameter.
1.0	ETEM	-	The liquid and vapor phases are maintained at equal temperatures if ETEM = 1.0 whereas if ETEM = 0.0 the vapor temperature is maintained equal to the saturation temperature at the local pressure.
0.07	GAM1	$(\gamma-1)$	Parameter in the equation of state.
0.0	GX	g_x	Gravitational acceleration in the x- or r-direction.
0.0	GY	g_y	Gravitational acceleration in the y- or z-direction.
10.0	IBAR	-	Number of active computational cells in the x- or r-direction. 532 240

1.0	IMP	-	Program control parameter whose values determines whether an implicit (IMP = 1) or explicit (IMP = 0) solution method is utilized.
10.0	JBAR	-	Number of active computational cells in the y- or z-direction.
NØ	NAME	-	80-character problem name.
1.0	OMG	-	Relaxation parameter for the calculation of δp in the pressure iteration. A value of OMG = 1.7 often improves the rate of convergences for low Mach number flows.
0.0	PBC	-	Constant pressure boundary condition in bars.
1.0	PHCH	-	Parameter whose value determines whether phase change is computed (PHCH = 1) or omitted (PHCH = 0).
1.0	PIN	-	Initial pressure condition in bars.
6 1	PLTDT	-	Time intervals between plots in ms.
0.1	PRTDT	-	Time intervals between prints in ms.
0.0	RADIUS	-	Inner radius of an annulus in cm.
0.0	RG	k	Pipe wall roughness.

0.958	ROL	ρ_l^o	Microscopic density of the liquid in g/cm ³ .
0.0	RPIPE	-	Pipe radius in cm.
.0008	SGWN	-	Product of the liquid surface tension and the critical Weber number.
373.0	TC	T_o	Reference temperature for liquid and vapor internal energy functions in K.
0.001	THC	θ_c	Void fraction below which the fluid is treated as pure liquid.
0.0	THIN	-	Initial void fraction condition.
373.0	TIN	-	Initial temperature condition in K.
10^4	TWFIN	-	Time when the problem is complete in ms.
0.	UI	-	Initial u velocity in cm/ms.
2.0	VELMX	-	Normalizing velocity for velocity plots in cm/ms.
0.	VI	-	Initial v velocity in cm/ms.
3.0×10^{-6}	VISL	v_l	Kinematic viscosity of the liquid.
2×10^{-4}	VISV	v_v	Kinematic viscosity of the vapor.

532 242

1.0	WB	-	Parameter that specifies the boundary condition for the bottom boundary.
1.0	WL	-	Parameter that specifies the boundary condition for the left boundary.
1.0	WR	-	Parameter that specifies the boundary condition for the right boundary.
1.0	WT	-	Parameter that specifies the boundary condition for the top boundary.

Following the NAMELIST data INPUT, the user has the option of defining rigid interior obstacles. If no interior obstacles are going to be defined, one blank card must follow the NAMELIST data INPUT defined above.

First Card After NAMELIST INPUT: NØ (Format I5)

NØ = Number of interior obstacles. If NØ ≠ 0, NØ additional cards are read; i.e., there is one card defining each obstacle.

NO Cards Following the Above Card: IBØB, IEØB, JBØB, JEØB, (Format 4I5)

IBØB = total number of δx cells, including the fictitious boundary cells, to the left side of the obstacle.

IEØB = total number of δx cells, including the fictitious boundary cells, to the right side of the obstacle.

JBØB = total number of δy cells, including the fictitious boundary cells, to the bottom side of the obstacle.

JEØB = total number of δy cells, including the fictitious boundary cells, to the top side of the obstacle.

B. COMMON Variables

FORTRAN	ALGEBRAIC	
SYMBOL	SYMBOL	DEFINITION
ASQ	a^2	Input quantity
CDG	C_d	Input quantity

CHL	C_l	Input quantity
CHV	C_v	Input quantity
CYL	-	Input quantity
DELT	δt	Input quantity
DELTMAX	δt_{max}	Input quantity
DELX	δx	Input quantity
DELY	δy	Input quantity
DFVEL	-	Input quantity
DIM	-	Input quantity
ECL	E_l	Input quantity
ECV	E_v	Input quantity
EDL	-	Input quantity
EDV	-	Input quantity
EIL	-	Initial total energy of the liquid state
EIV	-	Initial total energy of the vapor state
EI2	-	Initial total energy of the two-phase mixture
ELHT	λ	Input quantity
EPSI	ε	Input quantity
ET	-	Total energy
ESTEM	-	Input quantity
ETEM1	-	1.0 - ETEM
FLG	-	Iteration logic control flag
GAM1	($\gamma-1$)	Input quantity
II	-	Indexing variable
IMAX	-	Total number of cells in the x-direction (including fictitious cells)
IM1	-	IMAX - 1
IM2	-	IMAX - 2
IPL	-	Index of the first cell in the x-direction whose centered quantities are computed
IPR	-	Index of the last cell in the x-direction whose centered quantities are computed
ITER	-	Counter for the number of iterations between junctions and components to achieve convergence during one system time step
JJ	-	Indexing variable

JMAX	-	Total number of cells in the y-direction (including fictitious cells)
JM1	-	JMAX - 1
JM2	-	JMAX - 2
JPB	-	Index of the first cell in the y-direction whose centered quantities are computed
JPT	-	Index of the last cell in the y-direction whose centered quantities are computed
ØMG	-	Input quantity
PBC	-	Input quantity
PIN	-	Input quantity
PMAX	-	Temporary pressure variable
PNV	-	$4/3 \pi N$
PT	-	Temporary pressure variable
RDX	$1/\delta x$	Reciprocal δx
RDY	$1/\delta y$	Reciprocal δy
RG	k	Input quantity
RØIL	-	Initial total density of the liquid state
RØIV	-	Initial total density of the vapor state
RØI2	-	Init' . total density of the two-phase mixture
RØL	ρ_l^o	Input quantity
RØT	-	Temporary total density variable
RPIPE	-	Pipe radius in cm
RVIL	-	Initial vapor density in the liquid state
RVIV	-	Initial vapor density in the vapor state
RVI2	-	Initial vapor density in the two-phase mixture
RVT	-	Temporary total vapor density variable
SGWN	-	Input quantity
TC	T_o	Input quantity
THC	θ_c^o	Input quantity
THC1	-	$1.0 - \theta_c^o$
THIN	-	Input quantity
THTM	-	Temporary void fraction variable
TH1	-	Temporary void fraction variable

VISL	v_l	Input quantity
VISV	v_v	Input quantity
WB	-	Input quantity
WL	-	Input quantity
WR	-	Input quantity
WT	-	Input quantity
A(I,J)	-	Thickness of a mesh cell
BETA(I,J)	$(\partial F / \partial p)_{i,j}^{-1}$	Reciprocal derivative of the pressure function
E(I,J)	$I_{i,j}^{n+1}$	Time n+1 specific internal energy
EN(I,J)	$I_{i,j}^n$	Time n specific internal energy
ITITLE(K)	-	Two-word array used by LASL plotting package
NAME(K)	-	Input quantity
P(I,J)	$p_{i,j}$	Pressure
Q(I,J)	-	Two-dimensional array used by LASL plotting package
QL(K)	-	One-dimensional used to set up plotting
R ϕ (I,J)	$\rho_{i,j}^{n+1}$	Time n+1 mixture densities
R ϕ N(I,J)	$\rho_{i,j}^n$	Time n mixture densities
RV(I,J)	$(\rho_v)_{i,j}^{n+1}$	Time n+1 macroscopic vapor density
RVN(I,J)	$(\rho_v)_{i,j}^n$	Time n macroscopic vapor density
U(I,J)	$u_{i,j}^{n+1}$	Time n+1 center-of-mass velocity in the x-direction
UD(I,J)	$(u_r)_{i,j}^{n+1}$	Time n+1 relative or drift velocity in the x-direction
UN(I,J)	$u_{i,j}^n$	Time n center-of-mass velocity in the x-direction
V(I,J)	$u_{i,j}^{n+1}$	Time n+1 center-of-mass velocity in the y-direction
VD(I,J)	$(u_r)_{i,j}^{n+1}$	Time n+1 relative or drift velocity in the y-direction
VN(I,J)	$u_{i,j}^n$	Time n center-of-mass velocity in the y-direction

XC(K)	-	One-dimensional array used by LASL plotting package
YC(K)	-	One-dimensional array used by LASL plotting package

C. Subroutines and Graphics Output

BC	Implements specified boundary conditions.
DRIFT	Calculates the relative velocity between phases.
EQS	Calculates the equation of state.
PFRIC	Calculates the effects of pipe friction.
PHCHR	Calculates the phase change rate
PITER	Determines the n+l pressure and velocities by an iterative procedure.
TSTEP	Calculates the variable time step and the iteration parameter BETA(I,J).

Installations without graphics capabilities should delete lines SOLADF.613 thru SOLADF.698. to eliminate film writing and plotting commands. Installations with graphics should have routines equivalent to those listed below. A brief description of their function is included.

In addition the FORTRAN WRITE(N,XX) has special meaning under LASL operating systems. Its use results in printed data being sent automatically to the system graphics file.

ROUTINE	COMMENTS
ADV	Advance film
CONTRJB	Produce contour plots
CONVRT	Convert problem coordinates to plotting coordinates
DRV	Draw a vector
FRAME	Draw a frame
LINCNT	Position film line counter
SPLLOT	Standard plotting routine

V. EXAMPLE PROBLEM

To verify the proper implementation of the SOLA-DF code on a users computing system, the following example problem is included. A schematic of the problem geometry is shown in Fig. 3. A two-phase equilibrium mixture of steam and water (void fraction = 0.329, temperature = 557.1 K, and pressure = 68.0

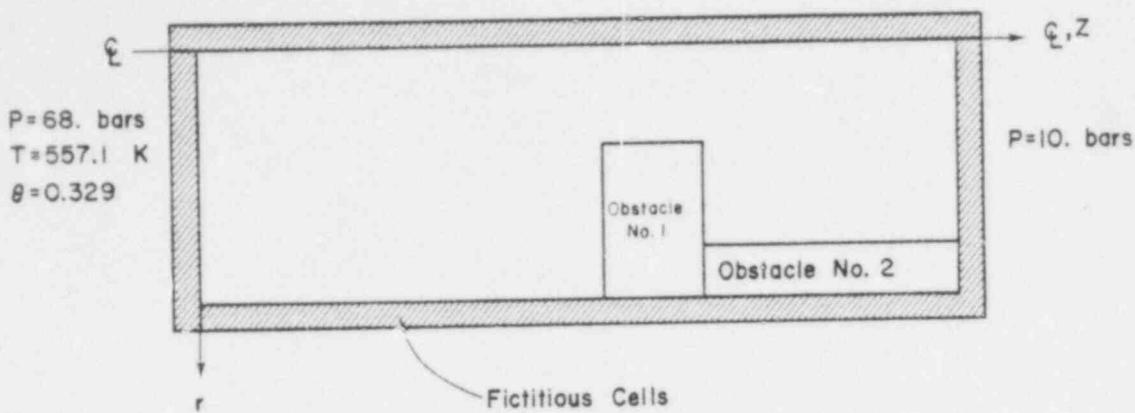


Fig. 3. Schematic of the computing mesh for the example problem.

bars) flows through a nozzle before exiting into a region of constant pressure = 10.0 bars. The flow, which begins from rest, is calculated until a steady state is achieved. The input data that specifies the problem is presented in Table I. The second part of this table shows a listing of the $A(I,J)$ array, which in this case (cylindrical coordinates) merely represents the radial distance from the center line to the center of the mesh cell. Note that obstacle cells set $A(I,J) = 0$. The data that specifies the initial states in the computational mesh is printed in Table II as the CYCLE = 0, TIME = 0 solution. The solution after 1000 time steps or 1s is given in Table III to provide the user with the transient solution check. The steady state solution is 4000 time steps or 4s and is given in Table IV. Velocity vector plots are presented at 1s and 4s in Figs. 4 and 5, respectively.

VI. SUMMARY

A description has been presented of a new computer program, SOLA-DF, for the solution of transient, one- and two-dimensional, two-phase flows. The one-dimensional formulation includes a variable area treatment. The fluid dynamics is described by a nonequilibrium, drift-flux formulation of the fluid conservation laws. An effort has been made to use relatively simple numerical solution procedures and modular programming in order to provide a framework that can be easily modified and adapted to different kinds of flow problems. In addition, a limited amount of implicitness is used to relax excessively restrictive time step limitations encountered in purely explicit integration methods. Even though the SOLA-DF code has a simple structure its flexibility offers capabilities for treating a wide range of two-phase flow problems.

TABLE I
EXAMPLE PROBLEM INPUT DATA

SOLA-OF EXAMPLE PROBLEM 01/20/79

```

ALPHA= 1.00000E+00
ASQ= 1.12470E+04
BURN= 1.00000E+02
CDG= 5.00000E-01
CHL= 4.43400E+01
CHV= 6.67000E+00
CYL= 1.00000E+00
DELT= 4.00000E-03
DELMX= 1.00000E-03
DELX= 2.20027E-01
DELY= 4.81753E-01
DFVEL= 0.
DIM= 2.00000E+00
ECL= 4.17400E+03
ECV= 2.50000E+04
EDL= 1.46600E-06
EDV= 6.00000E-07
ELHT= 1.76000E+04
EPSI= 1.00000E-04
ETEM= 1.00000E+00
GAMI= 7.00000E-02
GX= 0.
GY= 0.
IBAR= 10
IMP= 1.00000E+00
JBAR= 30
OMG= 1.00000E+00
PSI= 1.00000E+01
PHCH= 1.00000E+00
PIN= 6.80000E+01
PLTDT= 1.00000E+00
PRTOT= 1.00000E+00
RADIUS= 0.
RG= 3.66000E-03
ROL= 7.41900E-01
RPIPE= 0.
SGMN= 8.00000E-04
TC= 3.73000E+02
THC= 5.00000E-03
THIN= 3.29000E-01
TIN= 5.57100E+02
TWFIN= 1.00000E+01
UT= 0.
VELMX= 2.00000E+00
VI= 0.
VISL= 2.50000E-06
VISV= 1.23000E-04
WB= 5
WL= 1
WRU= 1
WT= 5

```


POOR ORIGINAL

TABLE II (con't)

TABLE II (con't)

TABLE II (con't)

POOR ORIGINAL

TABLE III (con't)

4	20	7.64582E-02	7.22778E+00	5.13135E+01	2.67776E-01	3.64275E-01	5.39230E+02
5	20	0.	7.08588E+00	5.12915E+01	2.67463E-01	6.04682E-01	5.39214E+02
6	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
7	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
8	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
9	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
10	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	20	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	21	1.42057E-01	8.91252E+00	4.57900E+01	2.22956E-01	7.24602E-01	5.33480E+02
3	21	2.01133E-01	9.07780E+00	4.57169E+01	2.22351E-01	7.25406E-01	5.33397E+02
4	21	2.07763E-01	9.36120E+00	4.65668E+01	2.20960E-01	7.27267E-01	5.33193E+02
5	21	0.	9.80650E+00	4.64211E+01	2.19574E-01	7.29130E-01	5.32981E+02
6	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
7	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
8	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
9	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
10	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	21	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	22	5.27337E-01	1.01670E+01	4.18037E+01	1.81109E-01	7.80124E-01	5.26654E+02
3	22	9.53176E-01	1.02601E+01	4.14104E+01	1.78237E-01	7.83399E-01	5.26124E+02
4	22	1.66567E+00	1.03179E+01	4.06253E+01	1.72665E-01	7.91214E-01	5.25058E+02
5	22	2.95846E+00	1.02084E+01	3.91551E+01	1.62615E-01	8.04339E-01	5.23076E+02
6	22	1.11179E+00	4.36016E+00	3.69096E+01	1.62096E-01	8.03454E-01	5.25687E+02
7	22	4.53415E-01	2.08142E+00	3.69908E+01	1.95101E-01	7.57701E-01	5.34910E+02
8	22	3.19656E-01	1.52628E+00	3.68766E+01	2.08029E-01	7.39707E-01	5.37857E+02
9	22	0.	1.44054E+00	3.63273E+01	2.07909E-01	7.39564E-01	5.38290E+02
10	22	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	22	0.	0.	5.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	23	7.11971E-01	1.08186E+01	3.78987E+01	1.52297E-01	8.17792E-01	5.20736E+02
3	23	1.17115E+00	1.07827E+01	3.76315E+01	1.50424E-01	8.20208E-01	5.20310E+02
4	23	1.70296E+00	1.06688E+01	3.72974E+01	1.48013E-01	8.23334E-01	5.19746E+02
5	23	2.13710E+00	1.03790E+01	3.69362E+01	1.45618E-01	8.26437E-01	5.19178E+02
6	23	1.35084E+00	6.77062E+00	3.70332E+01	1.469615E-01	8.24531E-01	5.19833E+02
7	23	8.81194E-01	4.60567E+00	3.68301E+01	1.55090E-01	8.13149E-01	5.23419E+02
8	23	4.16199E-01	3.76429E+00	3.64112E+01	1.64795E-01	7.99387E-01	5.27430E+02
9	23	0.	3.53939E+00	3.59688E+01	1.67256E-01	7.95704E-01	5.28723E+02
10	23	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	23	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	24	6.31266E-01	1.12441E+01	3.52991E+01	1.34825E-01	8.40328E-01	5.16475E+02
3	24	9.58922E-01	1.111623E+01	3.52629E+01	1.34416E-01	8.40869E-01	5.16333E+02
4	24	1.25206E+00	1.101766E+01	3.53357E+01	1.34711E-01	8.40481E-01	5.16428E+02
5	24	1.45141E+00	1.07821E+01	3.55320E+01	1.35679E-01	8.39267E-01	5.16642E+02
6	24	1.03889E+00	9.39504E+00	3.57554E+01	1.37492E-01	8.36862E-01	5.17266E+02
7	24	5.99631E-01	6.73977E+00	3.58558E+01	1.39373E-01	8.34325E-01	5.18131E+02
8	24	2.43812E-01	6.10357E+00	3.58028E+01	1.40336E-01	8.32971E-01	5.18740E+02
9	24	0.	5.96067E+00	3.58049E+01	1.41010E-01	8.32036E-01	5.19125E+02
10	24	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	24	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	25	4.56556E-01	1.18198E+01	3.34368E+01	1.23177E-01	8.50237E-01	5.13235E+02
3	25	6.67861E-01	1.17478E+01	3.35202E+01	1.23629E-01	8.54654E-01	5.13377E+02
4	25	8.69688E-01	1.16449E+01	3.36895E+01	1.24558E-01	8.53615E-01	5.13571E+02
5	25	1.01738E+00	1.14785E+01	3.39147E+01	1.25758E-01	8.51955E-01	5.13953E+02
6	25	8.21457E-01	9.81713E+00	3.41034E+01	1.26979E-01	8.50368E-01	5.14387E+02
7	25	4.96073E-01	6.56974E+00	3.41966E+01	1.27733E-01	8.49400E-01	5.14725E+02
8	25	2.08905E-01	6.07215E+00	3.42146E+01	1.27848E-01	8.49230E-01	5.14929E+02
9	25	0.	7.95608E+00	3.42225E+01	1.27861E-01	8.49216E-01	5.14987E+02
10	25	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	25	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	26	3.08673E-01	1.27509E+01	3.17520E+01	1.13251E-01	8.67844E-01	5.10201E+02
3	26	4.64919E-01	1.27017E+01	3.18303E+01	1.13729E-01	8.67270E-01	5.10330E+02
4	26	6.10512E-01	1.26300E+01	3.19509E+01	1.14383E-01	8.66412E-01	5.10557E+02
5	26	7.28446E-01	1.25104E+01	3.211120E+01	1.15266E-01	8.65268E-01	5.10839E+02

TABLE III (con't)

6	26	6.29152E-01	1.13137E+01	3.22459E+01	1.16025E-01	8.64329E-01	5.11112E+02
7	26	4.03106E-01	1.03055E+01	3.23216E+01	1.16364E-01	8.63896E-01	5.11323E+02
8	26	1.80217E-01	1.00021E+01	3.23275E+01	1.16144E-01	8.63851E-01	5.11391E+02
9	26	0.	9.90407E+00	3.23513E+01	1.16139E-01	8.63871E-01	5.11442E+02
10	26	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	26	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	27	2.12902E-01	1.41400E+01	2.98623E+01	1.02983E-01	8.80780E-01	5.06762E+02
3	27	3.30814E-01	1.41081E+01	2.99055E+01	1.03219E-01	8.80469E-01	5.06628E+02
4	27	4.37235E-01	1.40573E+01	2.99707E+01	1.03564E-01	8.80051E-01	5.06964E+02
5	27	5.26643E-01	1.39675E+01	3.00676E+01	1.04079E-01	8.79404E-01	5.07145E+02
6	27	4.69615E-01	1.30903E+01	3.01430E+01	1.04440E-01	8.78953E-01	5.07305E+02
7	27	3.08514E-01	1.24057E+01	3.01845E+01	1.04552E-01	8.78822E-01	5.07402E+02
8	27	1.39920E-01	1.21196E+01	3.02052E+01	1.04570E-01	8.78809E-01	5.07456E+02
9	27	0.	1.20430E+01	3.02080E+01	1.04543E-01	8.78849E-01	5.07470E+02
10	27	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	27	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	28	1.55952E-01	1.61958E+01	2.75477E+01	9.11775E-02	8.95514E-01	5.02319E+02
3	28	2.43678E-01	1.61750E+01	2.75513E+01	9.12226E-02	8.95460E-01	5.02327E+02
4	28	3.19094E-01	1.61394E+01	2.75754E+01	9.13310E-02	8.95323E-01	5.02380E+02
5	28	3.90869E-01	1.60739E+01	2.76153E+01	9.15496E-02	8.95054E-01	5.02460E+02
6	28	3.50941E-01	1.54272E+01	2.76407E+01	9.16276E-02	8.94962E-01	5.02517E+02
7	28	2.30672E-01	1.49277E+01	2.76576E+01	9.15850E-02	8.95022E-01	5.02554E+02
8	28	1.04408E-01	1.47194E+01	2.76577E+01	9.15474E-02	8.95081E-01	5.02559E+02
9	28	0.	1.46624E+01	2.76635E+01	9.15199E-02	8.95115E-01	5.02580E+02
10	28	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	28	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	29	1.24545E-01	1.94840E+01	2.44554E+01	7.67776E-02	9.13253E-01	4.96008E+02
3	29	1.91311E-01	1.94697E+01	2.44421E+01	7.67210E-02	9.13322E-01	4.95980E+02
4	29	2.46998E-01	1.94442E+01	2.44381E+01	7.66999E-02	9.13348E-01	4.95970E+02
5	29	3.07963E-01	1.93974E+01	2.44399E+01	7.67298E-02	9.13312E-01	4.95986E+02
6	29	2.75077E-01	1.89250E+01	2.44329E+01	7.66326E-02	9.13434E-01	4.95964E+02
7	29	1.77560E-01	1.85658E+01	2.44193E+01	7.65011E-02	9.13803E-01	4.95936E+02
8	29	7.95629E-02	1.84185E+01	2.44108E+01	7.64162E-02	9.13713E-01	4.95921E+02
9	29	0.	1.83778E+01	2.44067E+01	7.63776E-02	9.13764E-01	4.95915E+02
10	29	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	29	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	30	1.07179E-01	2.59769E+01	1.99430E+01	5.80527E-02	9.35975E-01	4.85683E+02
3	30	1.61900E-01	2.59566E+01	1.99252E+01	5.79816E-02	9.35959E-01	4.85637E+02
4	30	2.09543E-01	2.59316E+01	1.99118E+01	5.79249E-02	9.36027E-01	4.85600E+02
5	30	2.67661E-01	2.58887E+01	1.98638E+01	5.78731E-02	9.36089E-01	4.85569E+02
6	30	2.36996E-01	2.55385E+01	1.98698E+01	5.77173E-02	9.36278E-01	4.85490E+02
7	30	1.50268E-01	2.52773E+01	1.98439E+01	5.75641E-02	9.36467E-01	4.85420E+02
8	30	6.63849E-02	2.51701E+01	1.98265E+01	5.74650E-02	9.36589E-01	4.85376E+02
9	30	0.	2.51395E+01	1.98198E+01	5.74211E-02	9.36644E-01	4.85357E+02
10	30	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	30	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	31	0.	2.59769E+01	1.00000E+01	5.90199E-02	9.34738E-01	4.86228E+02
3	31	0.	2.59566E+01	1.00000E+01	5.89517E-02	9.34819E-01	4.86185E+02
4	31	0.	2.59316E+01	1.00000E+01	5.88760E-02	9.34910E-01	4.86132E+02
5	31	0.	2.58887E+01	1.00000E+01	5.88350E-02	9.34980E-01	4.86108E+02
6	31	0.	2.55385E+01	1.00000E+01	5.87713E-02	9.35040E-01	4.86086E+02
7	31	0.	2.52773E+01	1.00000E+01	5.87138E-02	9.35114E-01	4.86073E+02
8	31	0.	2.51701E+01	1.00000E+01	5.86953E-02	9.35141E-01	4.86051E+02
9	31	0.	2.51395E+01	1.00000E+01	5.86951E-02	9.35174E-01	4.86033E+02
10	31	0.	0.	1.00000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	31	0.	0.	1.00000E+01	5.09973E-01	3.29000E-01	5.57053E+02

POOR ORIGINAL
532 260

TABLE IV
EXAMPLE PROBLEM SOLUTION AT CYCLE 4000

ITER	J	TIME	CYCLE	DELT	F	F1	
		U	V	P	R0	TH	TEM
1	2	0.	8.60350E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	2	0.	8.60322E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
3	2	0.	8.60290E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
4	2	0.	8.60229E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
5	2	0.	8.60174E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
6	2	0.	8.60121E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
7	2	0.	8.60077E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
8	2	0.	8.60043E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
9	2	0.	8.60021E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
10	2	0.	8.60011E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	2	0.	8.60000E-01	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	3	-2.88329E-05	8.60080E-01	6.80329E+01	5.10054E-01	3.28894E-01	5.57053E+02
3	3	-4.21305E-05	8.60050E-01	6.80329E+01	5.10054E-01	3.28893E-01	5.57053E+02
4	3	-5.35231E-05	8.60006E-01	6.80330E+01	5.10054E-01	3.28893E-01	5.57053E+02
5	3	-5.87578E-05	8.59951E-01	6.80332E+01	5.10054E-01	3.28892E-01	5.57053E+02
6	3	-5.72863E-05	8.59892E-01	6.80334E+01	5.10055E-01	3.28892E-01	5.57053E+02
7	3	-5.00591E-05	8.59834E-01	6.80337E+01	5.10056E-01	3.28891E-01	5.57053E+02
8	3	-3.89228E-05	8.59795E-01	6.80340E+01	5.10056E-01	3.28890E-01	5.57053E+02
9	3	-2.59036E-05	8.59746E-01	6.80342E+01	5.10057E-01	3.28889E-01	5.57053E+02
10	3	-1.26071E-05	8.59720E-01	6.80344E+01	5.10057E-01	3.28889E-01	5.57053E+02
11	3	0.	8.59709E-01	6.80345E+01	5.10058E-01	3.28888E-01	5.57053E+02
2	4	-6.41577E-05	8.59648E-01	6.80674E+01	5.10139E-01	3.28782E-01	5.57054E+02
3	4	-9.44277E-05	8.59614E-01	6.80675E+01	5.10139E-01	3.28781E-01	5.57054E+02
4	4	-1.21214E-04	8.59565E-01	6.80676E+01	5.10139E-01	3.28781E-01	5.57054E+02
5	4	-1.34952E-04	8.59503E-01	6.80679E+01	5.10140E-01	3.28780E-01	5.57054E+02
6	4	-1.33896E-04	8.59434E-01	6.80682E+01	5.10141E-01	3.28779E-01	5.57054E+02
7	4	-1.19218E-04	8.59366E-01	6.80686E+01	5.10142E-01	3.28778E-01	5.57054E+02
8	4	-9.44644E-05	8.59306E-01	6.80689E+01	5.10142E-01	3.28777E-01	5.57054E+02
9	4	-6.38978E-05	8.59257E-01	6.80692E+01	5.10143E-01	3.28776E-01	5.57054E+02
10	4	-3.14454E-	8.59225E-01	6.80694E+01	5.10144E-01	3.28775E-01	5.57054E+02
11	4	0.	8.59208E-01	6.80695E+01	5.10144E-01	3.28775E-01	5.57054E+02
2	5	-1.12056E-04	8.59126E-01	6.81038E+01	5.10229E-01	3.29664E-01	5.57054E+02
3	5	-1.65507E-04	8.59082E-01	6.81038E+01	5.10229E-01	3.29664E-01	5.57054E+02
4	5	-2.13868E-04	8.59017E-01	6.81040E+01	5.10229E-01	3.29663E-01	5.57054E+02
5	5	-2.40681E-04	8.58933E-01	6.81041E+01	5.10229E-01	3.29663E-01	5.57054E+02
6	5	-2.42341E-04	8.58839E-01	6.81043E+01	5.10230E-01	3.29652E-01	5.57054E+02
7	5	-2.19841E-04	8.58745E-01	6.81045E+01	5.10230E-01	3.29661E-01	5.57054E+02
8	5	-1.77643E-04	8.58659E-01	6.81047E+01	5.10230E-01	3.29661E-01	5.57054E+02
9	5	-1.22319E-04	8.58569E-01	6.81048E+01	5.10231E-01	3.29660E-01	5.57054E+02
10	5	-6.09572E-05	8.58540E-01	6.81049E+01	5.10231E-01	3.29660E-01	5.57054E+02
11	5	0.	8.58516E-01	6.81050E+01	5.10231E-01	3.29660E-01	5.57054E+02
2	6	-1.89445E-04	8.58731E-01	6.81409E+01	5.10320E-01	3.29544E-01	5.57055E+02
3	6	-2.79094E-04	8.58658E-01	6.81409E+01	5.10320E-01	3.29544E-01	5.57055E+02
4	6	-3.62334E-04	8.58550E-01	6.81408E+01	5.10319E-01	3.29544E-01	5.57055E+02
5	6	-4.10865E-04	8.58493E-01	6.81407E+01	5.10319E-01	3.29544E-01	5.57055E+02
6	6	-4.18219E-04	8.58250E-01	6.81406E+01	5.10319E-01	3.29545E-01	5.57055E+02
7	6	-3.84651E-04	8.58089E-01	6.81404E+01	5.10319E-01	3.29545E-01	5.57055E+02
8	6	-3.15589E-04	8.57941E-01	6.81403E+01	5.10319E-01	3.29546E-01	5.57055E+02
9	6	-2.20412E-04	8.57820E-01	6.81402E+01	5.10319E-01	3.29546E-01	5.57055E+02
10	6	-1.10967E-04	8.57735E-01	6.81401E+01	5.10319E-01	3.29546E-01	5.57055E+02
11	6	0.	8.57692E-01	6.81400E+01	5.10317E-01	3.29546E-01	5.57055E+02
2	7	-3.33659E-04	8.57899E-01	6.81758E+01	5.10408E-01	3.29427E-01	5.57055E+02
3	7	-4.95548E-04	8.57863E-01	6.81766E+01	5.10407E-01	3.29428E-01	5.57055E+02
4	7	-6.45375E-04	8.57835E-01	6.81763E+01	5.10407E-01	3.29429E-01	5.57055E+02
5	7	-7.34327E-04	8.57815E-01	6.81759E+01	5.10406E-01	3.29430E-01	5.57055E+02
6	7	-7.53020E-04	8.57831E-01	6.81753E+01	5.10404E-01	3.29432E-01	5.57055E+02
7	7	-6.93060E-04	8.57505E-01	6.81748E+01	5.10403E-01	3.29434E-01	5.57055E+02
8	7	-5.71254E-04	8.57206E-01	6.81742E+01	5.10402E-01	3.29436E-01	5.57055E+02
9	7	-4.00717E-04	8.56962E-01	6.81737E+01	5.10400E-01	3.29437E-01	5.57055E+02

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TABLE IV (con't)

6	26	1.90987E+00	2.65044E+01	1.27843E+01	3.13681E-02	9.67002E-01	4.63183E+04
7	26	1.34003E+00	2.19408E+01	1.35614E+01	3.38370E-02	9.54215E-01	4.56039E+02
8	26	7.30983E-01	1.69420E+01	1.38845E+01	3.48997E-02	9.63007E-01	4.67253E+02
9	26	0.	1.23919E+01	1.39929E+01	3.53486E-02	9.62468E-01	4.68036E+02
10	26	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	26	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	27	9.75137E-01	3.17932E+01	9.53500E+00	2.20091E-02	9.77387E-01	4.50244E+02
3	27	1.03747E+00	3.12037E+01	9.96543E+00	2.31305E-02	9.76159E-01	4.52019E+02
4	27	8.49159E-01	3.05182E+01	1.06821E+01	2.51114E-02	9.73976E-01	4.55013E+02
5	27	5.82177E-01	2.97156E+01	1.11594E+01	2.77505E-02	9.71053E-01	4.58628E+02
6	27	5.86646E-01	2.70868E+01	1.24002E+01	3.02442E-02	9.68265E-01	4.61819E+02
7	27	6.21527E-01	2.30276E+01	1.28617E+01	3.17052E-02	9.66624E-01	4.63591E+02
8	27	4.25708E-01	1.84796E+01	1.30689E+01	3.23548E-02	9.65888E-01	4.64418E+02
9	27	0.	1.42768E+01	1.31441E+01	3.26224E-02	9.65573E-01	4.64886E+02
10	27	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	27	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	28	-2.28070E-01	3.07341E+01	1.01167E+01	2.33562E-02	9.75919E-01	4.52330E+02
3	28	-4.44996E-01	3.04749E+01	1.05239E+01	2.45147E-02	9.74644E-01	4.54089E+02
4	28	-5.84845E-01	3.02319E+01	1.10467E+01	2.60756E-02	9.72914E-01	4.56359E+02
5	28	-4.47151E-01	2.98198E+01	1.15781E+01	2.77218E-02	9.71082E-01	4.58625E+02
6	28	-6.86925E-02	2.75861E+01	1.19633E+01	2.89624E-02	9.69696E-01	4.60257E+02
7	28	2.17915E-01	2.39495E+01	1.21888E+01	2.96571E-02	9.68920E-01	4.61155E+02
8	28	2.38443E-01	1.97929E+01	1.23052E+01	2.99933E-02	9.68545E-01	4.61580E+02
9	28	0.	1.59187E+01	1.23509E+01	3.01441E-02	9.68363E-01	4.61934E+02
10	28	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	28	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	29	-7.08758E-01	3.04778E+01	1.09105E+01	2.55364E-02	9.73520E-01	4.55650E+02
3	29	-9.32277E-01	3.04827E+01	1.10781E+01	2.60872E-02	9.72910E-01	4.56377E+02
4	29	-9.63059E-01	3.04816E+01	1.12570E+01	2.67091E-02	9.72212E-01	4.57246E+02
5	29	-6.70082E-01	3.02427E+01	1.14056E+01	2.72462E-02	9.71611E-01	4.57975E+02
6	29	-2.34953E-01	2.82432E+01	1.14952E+01	2.75807E-02	9.71236E-01	4.58449E+02
7	29	8.07452E-02	2.49196E+01	1.15509E+01	2.77604E-02	9.71039E-01	4.58730E+02
8	29	1.57163E-01	2.10938E+01	1.15930E+01	2.78633E-02	9.70925E-01	4.58848E+02
9	29	0.	1.74947E+01	1.16144E+01	2.79200E-02	9.70855E-01	4.59034E+02
10	29	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	29	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	30	-6.55800E-01	3.18081E+01	1.11120E+01	2.23802E-02	9.72745E-01	4.56647E+02
3	30	-8.06314E-01	3.17675E+01	1.10893E+01	2.62358E-02	9.72743E-01	4.56612E+02
4	30	-7.74254E-01	3.17273E+01	1.10440E+01	2.61843E-02	9.72792E-01	4.56529E+02
5	30	-4.93386E-01	3.13943E+01	1.09803E+01	2.60696E-02	9.72915E-01	4.56377E+02
6	30	-1.35281E-01	2.94113E+01	1.03222E+01	2.59324E-02	9.73066E-01	4.56176E+02
7	30	1.12993E-01	2.62255E+01	1.08932E+01	2.58424E-02	9.73166E-01	4.56084E+02
8	30	1.63176E-01	2.26174E+01	1.08833E+01	2.38010E-02	9.73213E-01	4.56038E+02
9	30	0.	1.92560E+01	1.08798E+01	2.57971E-02	9.73226E-01	4.56071E+02
10	30	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	30	0.	0.	6.80000E+01	5.09973E-01	3.29000E-01	5.57053E+02
2	31	0.	3.18081E+01	1.00000E+01	2.62316E-02	9.72752E-01	4.56638E+02
3	31	0.	3.17675E+01	1.00000E+01	2.62291E-02	9.72750E-01	4.56608E+02
4	31	0.	3.17273E+01	1.00000E+01	2.61774E-02	9.72800E-01	4.56620E+02
5	31	0.	3.13943E+01	1.00000E+01	2.60620E-02	9.72924E-01	4.56357E+02
6	31	0.	2.94113E+01	1.00000E+01	2.59255E-02	9.73074E-01	4.56165E+02
7	31	0.	2.62255E+01	1.00000E+01	2.58355E-02	9.73174E-01	4.56066E+02
8	31	0.	2.26174E+01	1.00000E+01	2.57920E-02	9.73223E-01	4.56026E+02
9	31	0.	1.92560E+01	1.00000E+01	2.57777E-02	9.73237E-01	4.56052E+02
10	31	0.	0.	1.00000E+01	5.09973E-01	3.29000E-01	5.57053E+02
11	31	0.	0.	1.00000E+01	5.09973E-01	3.29000E-01	5.57053E+02

POOR ORIGINAL

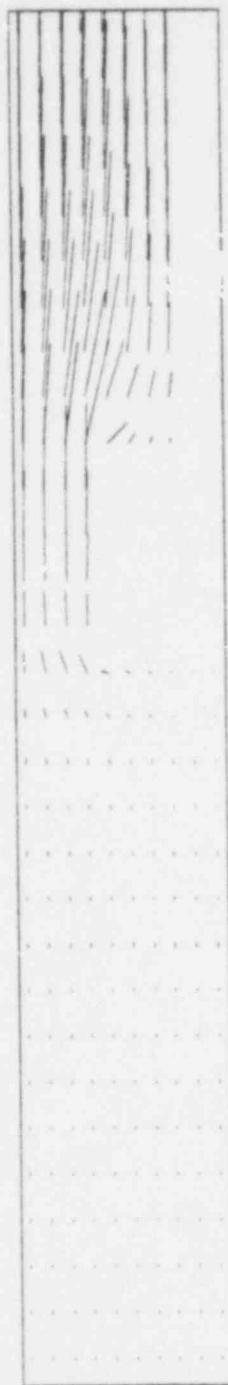


Fig. 4. Velocity vector plot of cycle 1000 for the example problem.

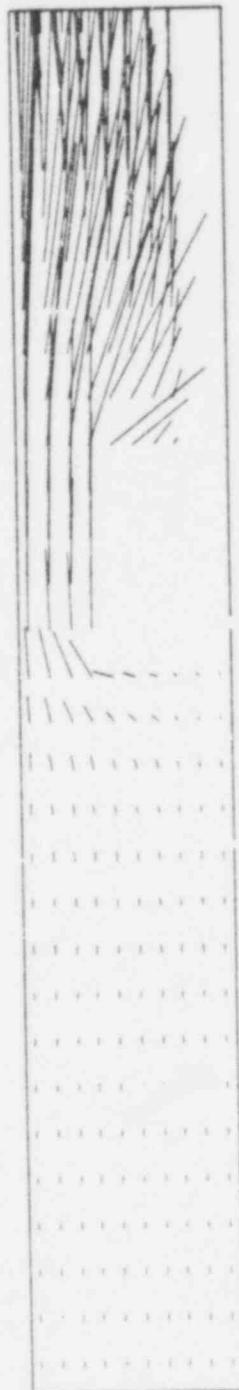


Fig. 5. Velocity vector plot of cycle 4000 for the example problem.

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REFERENCES

1. F. H. Harlow and A. A. Amsden, "Numerical Calculation of Multiphase Fluid Flow," *J. Comput. Phys.* 17, 19 (1975).
2. W. C. Rivard and M. D. Torrey, "K-FIX: A Computer Program for Transient, Two-Dimensional, Two-Fluid Flow," Los Alamos Scientific Laboratory report LA-NUREG-6623 (1977).
3. G. B. Wallis, One Dimensional Two-Phase Flow, McGraw-Hill (1969).
4. N. Zuber, "Flow Excursions and Oscillations in Boiling Two-Phase Flow Systems with Heat Addition," *Sym. Two-Phase Flow Dynamics*, E.U.R. 4288e, 1071 (1967).
5. C. W. Hirt, T. A. Oliphant, W. C. Rivard, N. C. Romero, and M. D. Torrey, "SOLA-LOOP: A Non-Equilibrium, Drift-Flux Code for Two-Phase Flow in Networks," manuscript in preparation.
6. C. W. Hirt, B. D. Nichols, and N. C. Romero, "SOLA - A Numerical Solution Algorithm for Transient Fluid Flows," Los Alamos Scientific Laboratory report LA-5852 (1975) and "SOLA - A Numerical Solution Algorithm for Transient Fluid Flows - Addendum," Los Alamos Scientific Laboratory report LA-5852, Add. (1976).
7. L. D. Cloutman, C. W. Hirt, and N. C. Romero, "SOLA-ICE: A Numerical Solution Algorithm for Transient Compressible Fluid Flows," Los Alamos Scientific Laboratory report LA-6236 (1976).
8. C. W. Hirt, J. K. Dienes, and L. R. Stein, "SOLA-FLX: A Solution Algorithm for Fluid-Structure Interactions of Cylindrical Shells," manuscript in preparation.
9. M. Ishii, Thermo-Fluid Dynamic Theory of Two-Phase Flow, (collection de la Direction des Etudes et Recherches d'Electricite de France, Eyrolles, Paris 1975).
10. J. R. Travis, F. H. Harlow, and A. A. Amsden, "Numerical Calculation of Two-Phase Flows," *Nuc. Sci. and Eng.* 61, 1 (1976).
11. T. G. Theofanous, T. Bohrer, M. Chen, and P. D. Patel, "Universal Solutions for Bubble Growth and the Influence of Microlayers, 15th Nat. Heat Transfer Conf., Paper 15, San Francisco, CA (August 1975).
12. D. Moalem and S. Sideman, "The Effect of Motion on Bubble Collapse," *Int. J. Heat and Mass Transfer* 16, 2321 (1973).

13. G. B. Wallis, H. J. Richter, and J. T. Kuo, "The Separated Flow Model for Two-Phase Flow," Electric Power Research Institute report NP-275, 88 (1976).
 14. F. H. Harlow and A. A. Amsden, "A Numerical Fluid Dynamics Calculation Method for All Flow Speeds," *J. Comput. Phys.* 8, 197 (1971).
 15. C. W. Hirt, "Heuristic Theory for Finite-Difference Equations," *J. Comput. Phys.* 2, 339 (1968).
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APPENDIX
FORTRAN IV Listing of the SOLA-DF Code LASL Code: LP#-0772

```

*COMDECK.COMOK
C
COMMON / SOL1 /
1A (27,82), BETA(27,82), E (27,82), EN (27,82),
2TITLE (2), NAME (10), P (27,82), Q (27,82),
3QL (82), R0 (27,82), RON (27,82), RV (27,82),
4RVN (27,82), U (27,82), UD (27,82), UN (27,82),
5V (27,82), VD (27,82), VN (27,82), XC (27),
6YC (82)
7
C
COMMON / SOL2 /
8ASQ, CDG, CHL, CHV, CYL, DELT, DELTMX, COMOK 1
9DELX, DELY, ECL, ECV, EDL, EDV, EIL, COMOK 12
10EIV, E12, ELHT, EPSI, ET, ETEM, ETEM1, COMOK 13
11EFLG, GAM1, II, IMAX, IM1, IM2, IPL, COMOK 14
12SIPR, ITER, JJ, JMAX, JM1, JM2, JPB, COMOK 15
13SPT, OMO, PSC, PIN, PMAX, PNV, PT, COMOK 16
14TDX, RDY, RG, ROIL, ROI1V, ROI2, ROL, COMOK 17
15BPT, RPIPE, RVIL, RVIV, RV12, RVT, SGN, COMOK 18
16TC, THC, THCI, THIN, THTM, TH1, VISL, COMOK 19
17VISV, WB, WL, WR, WT
18
C
19 INTEGER CYCLE, WB, WL, WR, WT
20 REAL IMP, LONG, NUA, NUC
21
C
22
23
24
25

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POOR ORIGINAL

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*DECK,SOLADF
  PROGRAM SOLA(INPUT,TAPE10=INPUT,OUTPUT,TAPE9=OUTPUT)
*CALL COMDK
C
  DIMENSION ZC(10)
C
  EVCAL(X)=ECV+CHV*(X-TC)
  ELCAL(X)=ECL+CHL*(X-TC)
  SATT(X)=255.2*117.8*X**0.223
  SATP(X)=((X-255.2)/117.8)**4.48
C
  NAMELIST / SOLDA /
  1ALPHA,    ASQ,      BUBN,      CGD,      CHL,      CHV,
  2CYL,      DELT,     DELTMX,    DELX,     DELY,     DFVEL,
  3DIM,      ECL,      ECV,       EDL,      EDV,     ELHT,
  4EPSI,     ETEM,     GAM1,     GX,       GY,      IBAR,
  5IMP,      JBAR,     NAME,     OMO,     PBC,     PHCH,
  6IN,       PLDT,     PRTDT,    RADIUS,   RG,      ROL,
  7PIPE,     SGWN,     TC,       THC,     THIN,    TIN,
  8WIN,     UI,       VELMX,    VI,      VISL,    VISV,
  9WB,      WL,       WR,       WT
C
C  ZERO OUT THE DATA AND CONSTANT STORAGES
  DO 1 I=1,35518
  1 A(I)=0,
  DO 2 I=1,27
  2 XC(I)=0,
  DO 3 I=1,82
  3 YC(I)=0.
C
C  SET DEFAULT VALUES
C
  ALPHA=1.
  ASQ=1.234E+4
  BUBN=1.E+4
  CGD=0.5
  CHL=44.34
  CHV=6.67
  CYL=1.0
  DFVEL = 0.0
  DELT=1.0E-4
  DELTMX=1.0E-3
  DELX=1.0
  DELY=1.0
  DIM=2.0
  ECL=4.174E+3
  ECV=2.506E+4
  EDL=1.6E-6
  EDV=1.6E-7
  ELHT=1.76E+4
  EPSI=0.001
  ETEM=1.0
  GAM1=0.07
  GX = 0.0
  GY = 0.0
  IBAR=10
  IMP=1.0
  JBAR=10
  SOLADF 1
  SOLADF 2
  SOLADF 3
  SOLADF 4
  SOLADF 5
  SOLADF 6
  SOLADF 7
  SOLADF 8
  SOLADF 9
  SOLADF 10
  SOLADF 11
  SOLADF 12
  SOLADF 13
  SOLADF 14
  SOLADF 15
  SOLADF 16
  SOLADF 17
  SOLADF 18
  SOLADF 19
  SOLADF 20
  SOLADF 21
  SOLADF 22
  SOLADF 23
  SOLADF 24
  SOLADF 25
  SOLADF 26
  SOLADF 27
  SOLADF 28
  SOLADF 29
  SOLADF 30
  SOLADF 31
  SOLADF 32
  SOLADF 33
  SOLADF 34
  SOLADF 35
  SOLADF 36
  SOLADF 37
  SOLADF 38
  SOLADF 39
  SOLADF 40
  SOLADF 41
  SOLADF 42
  SOLADF 43
  SOLADF 44
  SOLADF 45
  SOLADF 46
  SOLADF 47
  SOLADF 48
  SOLADF 49
  SOLADF 50
  SOLADF 51
  SOLADF 52
  SOLADF 53
  SOLADF 54
  SOLADF 55
  SOLADF 56
  SOLADF 57
  SOLADF 58

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POOR ORIGINAL

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NAME(1)=10H *NO NAME*
OMG=1.0          SOL ADF    59
PBC=1.0          SOL ADF    60
PHCH=1.0         SOL ADF    61
PIN=1.0          SOL ADF    62
PLTDT=0.1        SOL ADF    63
PRTDT=0.1        SOL ADF    64
RADIUS = 0.0     SOL ADF    65
RG = 0.0          SOL ADF    66
ROL=0.958        SOL ADF    67
RPIPE = 0.0       SOL ADF    68
SGWN=8.0E-4       SOL ADF    69
TC=373.0          SOL ADF    70
THC=0.001        SOL ADF    71
THIN = 0.0        SOL ADF    72
TIN=373.0         SOL ADF    73
TWFIN=1.E+4       SOL ADF    74
UI = 0.0          SOL ADF    75
VELMX=2.0         SOL ADF    76
VI = 0.0          SOL ADF    77
VISL=3.0E-6       SOL ADF    78
VISV=2.0E-4       SOL ADF    79
WB=1              SOL ADF    80
WL=1              SOL ADF    81
WR=1              SOL ADF    82
WT=1              SOL ADF    83
                           SOL ADF    84
C
C READ AND WRITE INITIAL INPUT DATA
C
C READ(10,SOLDA)
C
WRITE(9,901) NAME(1),I=1,10          SOL ADF    85
WRITE(9,985) ALPHA                 SOL ADF    86
WRITE(9,1015) ASQ                  SOL ADF    87
WRITE(9,1050) BUBN                 SOL ADF    88
WRITE(9,1055) CDO                  SOL ADF    89
WRITE(9,1090) CHL                  SOL ADF    90
WRITE(9,1085) CHV                  SOL ADF    91
WRITE(9,930) CYL                  SOL ADF    92
WRITE(9,920) DELT                  SOL ADF    93
WRITE(9,890) DELTMX                SOL ADF    94
WRITE(9,910) DELX                  SOL ADF    95
WRITE(9,915) DELY                  SOL ADF    96
WRITE(9,1030) DFVEL                SOL ADF    97
WRITE(9,880) DIM                  SOL ADF    98
WRITE(9,1100) ECL                  SOL ADF    99
WRITE(9,1095) ECV                  SOL ADF   100
WRITE(9,1115) EDL                  SOL ADF   101
WRITE(9,1110) EDV                  SOL ADF   102
WRITE(9,1120) ELHT                 SOL ADF   103
WRITE(9,935) EPSI                 SOL ADF   104
WRITE(9,1025) ETEM                 SOL ADF   105
WRITE(9,1010) GAM1                 SOL ADF   106
WRITE(9,940) GX                   SOL ADF   107
WRITE(9,945) GY                   SOL ADF   108
WRITE(9,900) IBAR                 SOL ADF   109
WRITE(9,1020) IMP                  SOL ADF   110
WRITE(9,905) JBAR                 SOL ADF   111

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WRITE(9,980) OMG	SOLADF	117
WRITE(9,1028) PBC	SOLADF	118
WRITE(9,1035) PHCH	SOLADF	119
WRITE(9,1130) PIN	SOLADF	120
WRITE(9,975) PLTDT	SOLADF	121
WRITE(9,970) PRTDT	SOLADF	122
WRITE(9,1125) RADIUS	SOLADF	123
WRITE(9,1070) RG	SOLADF	124
WRITE(9,1045) ROL	SOLADF	125
WRITE(9,1065) RPPIPE	SOLADF	126
WRITE(9,1060) SQIN	SOLADF	127
WRITE(9,1105) TC	SOLADF	128
WRITE(9,1040) THC	SOLADF	129
WRITE(9,1140) THIN	SOLADF	130
WRITE(9,1135) TIN	SOLADF	131
WRITE(9,965) TWFIN	SOLADF	132
WRITE(9,950) UI	SOLADF	133
WRITE(9,960) VELMX	SOLADF	134
WRITE(9,955) VI	SOLADF	135
WRITE(9,1080) VISL	SOLADF	136
WRITE(9,1075) VISV	SOLADF	137
WRITE(9,1005) WB	SOLADF	138
WRITE(9,990) WL	SOLADF	139
WRITE(9,995) WR	SOLADF	140
WRITE(9,1000) WT	SOLADF	141
880 FORMAT(10X, 10H DIM= ,1PE12.5)	SOLADF	142
890 FORMAT(10X, 10H DELTMX= ,1PE12.5)	SOLADF	143
900 FORMAT(10X, 10H IBAR= ,15)	SOLADF	144
905 FORMAT(10X, 10H JBAR= ,15)	SOLADF	145
910 FORMAT(10X, 10H DELX= ,1PE12.5)	SOLADF	146
915 FORMAT(10X, 10H DELY= ,1PE12.5)	SOLADF	147
920 FORMAT(10X, 10H DELT= ,1PE12.5)	SOLADF	148
930 FORMAT(10X, 10H CYL= ,1PE12.5)	SOLADF	149
935 FORMAT(10X, 10H EPSI= ,1PE12.5)	SOLADF	150
940 FORMAT(10X, 10H GX= ,1PE12.5)	SOLADF	151
945 FORMAT(10X, 10H OY= ,1PE12.5)	SOLADF	152
950 FORMAT(10X, 10H UI= ,1PE12.5)	SOLADF	153
955 FORMAT(10X, 10H VI= ,1PE12.5)	SOLADF	154
960 FORMAT(10X, 10H VELMX= ,1PE12.5)	SOLADF	155
965 FORMAT(10X, 10H TWFIN= ,1PE12.5)	SOLADF	156
970 FORMAT(10X, 10H PRTDT= ,1PE12.5)	SOLADF	157
975 FORMAT(10X, 10H PLTDT= ,1PE12.5)	SOLADF	158
980 FORMAT(10X, 10H OMG= ,1PE12.5)	SOLADF	159
985 FORMAT(10X, 10H ALPHA= ,1PE12.5)	SOLADF	160
990 FORMAT(10X, 10H WL= ,15)	SOLADF	161
995 FORMAT(10X, 10H WR= ,15)	SOLADF	162
1000 FORMAT(10X, 10H WT= ,15)	SOLADF	163
1005 FORMAT(10X, 10H WB= ,15)	SOLADF	164
1010 FORMAT(10X, 10H GAM1= ,1PE12.5)	SOLADF	165
1015 FORMAT(10X, 10H ASQ= ,1PE12.5)	SOLADF	166
1020 FORMAT(10X, 10H IMP= ,1PE12.5)	SOLADF	167
1025 FORMAT(10X, 10H ETEM= ,1PE12.5)	SOLADF	168
1028 FORMAT(10X, 10H PBC= ,1PE12.5)	SOLADF	169
1030 FORMAT(10X, 10H DFVEL= ,1PE12.5)	SOLADF	170
1035 FORMAT(10X, 10H PHCH= ,1PE12.5)	SOLADF	171
1040 FORMAT(10X, 10H THC= ,1PE12.5)	SOLADF	172
1045 FORMAT(10X, 10H ROL= ,1PE12.5)	SOLADF	173
1050 FORMAT(10X, 10H BUBN= ,1PE12.5)	SOLADF	174

532 272

POOR ORIGINAL

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1055 FORMAT(10X, 10H      CGD= ,1PE12.5)           SOLADF 175
1060 FORMAT(10X, 10H      SGWN= ,1PE12.5)           SOLADF 176
1065 FORMAT(10X, 10H      RPIPE= ,1PE12.5)           SOLADF 177
1070 FORMAT(10X, 10H      RG= ,1PE12.5)             SOLADF 178
1075 FORMAT(10X, 10H      VISY= ,1PE12.5)           SOLADF 179
1080 FORMAT(10X, 10H      VISL= ,1PE12.5)           SOLADF 180
1085 FORMAT(10X, 10H      CHV= ,1PE12.5)           SOLADF 181
1090 FORMAT(10X, 10H      CHL= ,1PE12.5)           SOLADF 182
1095 FORMAT(10X, 10H      ECV= ,1PE12.5)           SOLADF 183
1100 FORMAT(10X, 10H      ECL= ,1PE12.5)           SOLADF 184
1105 FORMAT(10X, 10H      TC= ,1PE12.5)           SOLADF 185
1110 FORMAT(10X, 10H      EDV= ,1PE12.5)           SOLADF 186
1115 FORMAT(10X, 10H      EDL= ,1PE12.5)           SOLADF 187
1120 FORMAT(10X, 10H      ELHT= ,1PE12.5)           SOLADF 188
1125 FORMAT(10X, 10H      RADIUS= ,1PE12.5)         SOLADF 189
1130 FORMAT(10X, 10H      PIN= ,1PE12.5)           SC ADF 190
1135 FORMAT(10X, 10H      TIN= ,1PE12.5)           SOLADF 191
1140 FORMAT(10X, 10H      THIN= ,1PE12.5)
   5 FORMAT(8X,1HJ,7X,1HJ,12X,2HUD,17X,2HVD,18X,2HRV)
   10 FORMAT(5X,12,6X,12,5X,1PE12.5,6X,E12.5,6X,E12.5)
   35 FORMAT(1H1)
901 FORMAT(1H1,10A10,//)
   44 FORMAT(6X,7HCYCLE= ,15,8X,4HTD= ,1PE12.5,8X,4HT2= ,E12.5,
     19X,5HITER= ,15)                                SOLADF 196
   45 FORMAT(10A8)                                     SOLADF 197
   46 FORMAT(1H+,80X,2HT= ,1PE10.3,4X,6HCYCLE= ,14)  S ADF 198
   47 FORMAT(5X,1H1,4X,1HJ,12X,1HU,12X,1HV,17X,1HP,14X,2HRO,14X,2HTH,
     114X,3HTEM)                                     ADF 199
   48 FORMAT(5X,12,3X,12,5X,1PE12.5,3X,E12.5,3X,E12.5,3X,E12.5,
     15X,E12.5)                                     SOLADF 200
   49 FORMAT(2X,5HITER= ,15,7X,6HTIME= ,1PE12.5,3X,7HCYCLE= ,14,
     14X,5HDELT= ,E12.5,4X,3HF= ,E12.5,4X,4HF1= ,E12.5) SOLADF 201
C   COMPUTE CONSTANT TERMS
   IMAX=1BAR+2                                      SOLADF 202
   JMAX=JBAR+2                                      SOLADF 203
   IM1=IMAX-1                                       SOLADF 204
   JM1=JMAX-1                                       SOLADF 205
   IM2=IMAX-2                                       SOLADF 206
   JM2=JMAX-2                                       SOLADF 207
   RDX=1.0/DELX                                      SOLADF 208
   RDX=1.0/DELY                                      SOLADF 209
C   CONTOUR PLOT, SETTING UP VARIABLE VALUES
   XC(1)=0.0                                         SOLADF 210
   DO 15 I=2,IM1                                     SOLADF 211
   XC(I)=DELX*(FLOAT(I)-1.5)                         SOLADF 212
15 CONTINUE
   NNX=-1BAR                                         SOLADF 213
   YC(1)=0.0                                         SOLADF 214
   DO 20 J=2,JM1                                     SOLADF 215
   YC(J)=DELY*(FLOAT(J)-1.5)                         SOLADF 216
20 CONTINUE
   NNY=-JBAR                                         SOLADF 217
   NZX=27                                           SOLADF 218
   NTY=82                                           SOLADF 219
   NC=10                                           SOLADF 220
   ZMN=-1.0                                         SOLADF 221
   ZMX=-1.0                                         SOLADF 222
   DLZ=0.0                                           SOLADF 223
   SOLADF 224
   SOLADF 225
   SOLADF 226
   SOLADF 227
   SOLADF 228
   SOLADF 229
   SOLADF 230
   SOLADF 231
   SOLADF 232

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POOR ORIGINAL

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ZC=1.0                                     SOL ADF    233
DMPX=DELX*(FLOAT(IBAR))                   SOL ADF    234
DMPY=DELY*(FLOAT(JBAR))                   SOL ADF    235
IGRD=0.0                                    SOL ADF    236
NTITLE=10                                   SOL ADF    237
XLABEL=1HX                                 SOL ADF    238
YLABEL=1HY                                 SOL ADF    239
NLBL=1                                      SOL ADF    240
NYLBL=1                                      SOL ADF    241
LONG=IBAR*DELX,                            SOL ADF    242
HIGH=JBAR*DELY                            SOL ADF    243
IYB=916                                     SOL ADF    244
IF(LONG.LE.(1.13556*HIGH))GO TO 30       SOL ADF    245
IXL=0                                       SOL ADF    246
IXR=1022                                    SOL ADF    247
IYT=916+HIGH*1022/LONG                    SOL ADF    248
GO TO 33                                    SOL ADF    249
30 X=LONG*.50/HIGH                         SOL ADF    250
IXL=511-X                                  SOL ADF    251
IXR=511+X                                  SOL ADF    252
IYT=16                                     SOL ADF    253
33 CONTINUE                                 SOL ADF    254
VELMXI=AMIN(DELX,DELY)/VELMX              SOL ADF    255
C   INITIALIZE NUMERICAL CONSTANTS          SOL ADF    256
THPRT=0.                                     SOL ADF    257
THPLT=0.                                     SOL ADF    258
T=0.                                         SOL ADF    259
ITER=0                                       SOL ADF    260
CYCLE=0                                      SOL ADF    261
NEX=0                                         SOL ADF    262
IPL=2                                         SOL ADF    263
IF(WL.EQ.5)IPL=3                           SOL ADF    264
IPR=IM1                                     SOL ADF    265
IF(WR.EQ.5)IPR=IM1-1                       SOL ADF    266
JPB=2                                         SOL ADF    267
IF(WB.EQ.5)JPB=3                           SOL ADF    268
JPT=JM1                                     SOL ADF    269
IF(WT.EQ.5)JPT=JM1-1                       SOL ADF    270
C   INITIALIZE PHYSICAL CONSTANTS          SOL ADF    271
IF(DIM.LT.1.5)DELX=1.0E+10                 SOL ADF    272
ETEM1=1.0-ETEM                            SOL ADF    273
THC1=1.0-THC                             SOL ADF    274
PNV=BUBN*4.1888                           SOL ADF    275
TC2=2.0*TC                                SOL ADF    276
C   INITIALIZE AREAS                      SOL ADF    277
DO 53 J=1,JMAX                           SOL ADF    278
DO 53 I=1,,MAX                           SOL ADF    279
A(I,J)=1.0-CYL*CYL*(RADIUS+DELX*(FLOAT(I)-1.5)) SOL ADF    280
IF(I.EQ.1.AND.WL.LT.3)A(I,J)=0.0           SOL ADF    281
IF(I.EQ.1MAX.AND.WR.LT.3)A(I,J)=0.0         SOL ADF    282
IF(J.EQ.1.AND.WB.LT.3)A(I,J)=0.0           SOL ADF    283
IF(J.EQ.JMAX.AND.WT.LT.3)A(I,J)=0.0         SOL ADF    284
53 CONTINUE                                 SOL ADF    285
C   DEFINE SPECIAL AREAS                  SOL ADF    286
READ(10,210) NO                           SOL ADF    287
210 FORMAT(415)                            SOL ADF    288
IF(NO.LE.0)GO TO 216                      SOL ADF    289
DO 215 K=1,NO                            SOL ADF    290

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POOR ORIGINAL

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READ(10,210) IB0B,IE0B,JB0B,JE0B
DO 211 J=JB0B,JE0B
DO 211 I=IB0B,IE0B
211 A(I,J)=0.0
215 CONTINUE
216 CONTINUE
  WRITE(9,220)
220 FORMAT(1H1,12H AREA ARRAY )
  DO 230 J=1,JMAX
    JU=JMAX+1-J
230 WRITE(9,225)(A(I,JU),I=1,IMAX)
225 FORMAT(1H ,12(1X,1PE9.2))
  WRITE(9,235)
235 FORMAT(1H1)
C   INITIALIZE VARIABLES
  DO 36 J=1,JMAX
  DO 56 I=1,IMAX
    U(I,J)=U1
    V(I,J)=V1
    UD(I,J)=VD(I,J)=0.0
    IF(THIN.LT.THC)GO TO 38
    IF(THIN.LT.THC)GO TO 36
C   VAPOR STATE (P,T)
  TH=THIN
  EVT=EVCAL(TIN)
  RV(1,J)=RV(V=TH*PIN/(GAM)*EVT)
  RO(1,J)=RO(V=RV(1,J)+(1.0-TH)*ROL
  E(1,J)=E(V=(RV(1,J)*EVT*(RO(1,J)-RV(1,J))*ELCAL(TIN))//
  1RO(1,J))
  GO TO 40
36 CONTINUE
  TSAT=SATT(PIN)
  EVSAT=EVCAL(TSAT)
  IF(THIN.LT.THC)GO TO 38
C   SATURATED STATE (P,TH)
  RV12=RV(1,J)=THIN*PIN/(GAM)*EVSAT
  RO12=RO(1,J)=RV(1,J)+(1.0-THIN)*ROL
  E12=E(1,J)=(RV(1,J)*EVSAT+(1.0-THIN)*ROL*ELCAL(TSAT))/RO(1,J)
  GO TO 40
38 CONTINUE
C   LIQUID STATE (P,T)
  EVSAT=EVCAL(TIN)
  PSAT=SATP(TIN)
  RV(1,J)=RV(L=THC*PSAT/(GAM)*EVSAT)
  TH=(PSAT-PIN)/(ASQ*ROL)+THC
  RO(1,J)=RO(L=RV(1,J)+(1.0-TH)*ROL
  E(1,J)=E(L=(RV(1,J)*EVSAT+(1.0-TH)*ROL*ELCAL(TIN))/RO(1,J)
40 CONTINUE
  ET=E(1,J)
  ROT=RO(1,J)
  RVT=RV(1,J)
  II=1
  JJ=J
  CALL EOS
  P(1,J)=PT
56 CONTINUE
  CALL BC
  GO TO 502

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C      START CYCLE
59 CONTINUE
    TER=0
    FLG=IMP
    PMAX=0.0
C      COMPUTE TEMPORARY U AND V
    DO 70 J=2,IM1
    DO 70 I=2,IM1
    U(I,J)=0.0
    AUB=2.0*A(I,J)*A(I+1,J)
    IF(AUB.EQ.0.0)GO TO 60
    AUBR=(A(I,J)+A(I+1,J))/AUB
    FUX=0.5*ROX*(UN(I,J)*UN(I+1,J)-UN(I-1,J))
    1-ALPHA*ABS(UN(I,J))*UN(I+1,J)-2.0*UN(I,J)+UN(I-1,J))
    FUY=ROY/8.0*((VN(I,J)+VN(I+1,J)+VN(I,J-1)+VN(I+1,J-1))
    1*(UN(I,J+1)-UN(I,J-1))
    2-ALPHA*ABS(VN(I,J)+VN(I+1,J)+VN(I,J-1)+VN(I+1,J-1))
    3*(UN(I,J+1)-2.0*UN(I,J)+UN(I,J-1))
    VDT=0.5*(VD(I,J)+VD(I+1,J))*(A(I,J)*A(I,J+1)/(A(I,J)+A(I,J+1))+
    1A(I+1,J)*A(I+1,J+1)/(A(I+1,J)+A(I+1,J+1)))
    VDB=0.5*(VD(I,J-1)+VD(I+1,J-1))*(A(I,J)*A(I,J-1)/(A(I,J)+
    1A(I,J-1)*A(I+1,J)*A(I+1,J-1)/(A(I+1,J)+A(I+1,J-1)))
    RUR=0.5*RVN(I+1,J)*(RON(I+1,J)-RVN(I+1,J))/RO(I+1,J)+
    1*(UD(I,J)+UD(I+1,J))*A(I+1,J)
    RUL=0.5*RVN(I,J)*(RON(I,J)-RVN(I,J))/RON(I,J)*(UD(I-1,J)+UD(I,J))-
    1*A(I,J)
    RUC=0.5*UD(I,J)*(RVN(I,J)*(RON(I,J)-RVN(I,J))/RON(I,J)+RVN(I+1,J)*
    1*(RON(I+1,J)-RVN(I+1,J))/RON(I+1,J))
    RUT=0.5*UD(I,J+1)*(RVN(I,J+1)*(RON(I,J+1)-RVN(I,J+1))/RON(I,J+1)+
    1RVN(I+1,J+1)*(RON(I+1,J+1)-RVN(I+1,J+1))/RON(I+1,J+1))
    RUB=0.5*UD(I,J-1)*(RVN(I,J-1)*(RON(I,J-1)-RVN(I,J-1))/RON(I,J-1)+
    1RVN(I+1,J-1)*(RON(I+1,J-1)-RVN(I+1,J-1))/RON(I+1,J-1))
    FDU=AUBR*(ROX*(RUR*(UD(I,J)+UD(I+1,J))+ALPHA*ABS(RUR)*(UD(I,J)-
    1UD(I+1,J))-RUL*(UD(I-1,J)+UD(I,J))-ALPHA*ABS(RUL)*(UD(I-1,J)-
    2UD(I,J)))+RDY*(VDT*(RUC+RUT)+ALPHA*ABS(VDT)*(RUC-RUT)-
    3VDB*(RUB+RUC)-ALPHA*ABS(VDB)*(RUB-RUC)))/(RO(I,J)+RO(I+1,J))
    VISX=0.0
    U(I,J)=UN(I,J)+DELT*(2.0*(P(I,J)-P(I+1,J))+ROX/(RO(I,J)+RO(I+1,J))-
    1+GX-FUX-FUY+VISX-FDU)
60 CONTINUE
    V(I,J)=0.0
    AVB=2.0*A(I,J)*A(I,J+1)
    IF(AVB.EQ.0.0)GO TO 65
    AVBR=(A(I,J)+A(I,J+1))/AVB
    FVX=ROX/8.0*((UN(I-1,J+1)+UN(I,J+1)+UN(I-1,J)+UN(I,J))*
    1(VN(I+1,J)-VN(I-1,J))-2ALPHA*ABS(UN(I-1,J+1)+UN(I,J+1)+UN(I-1,J)+UN(I,J))*
    3(VN(I+1,J)-2.0*VN(I,J)+VN(I-1,J))-FVY=0.5*ROY*(VN(I,J)*(VN(I,J+1)-VN(I,J-1))-
    1ALPHA*ABS(VN(I,J))*(VN(I,J+1)-2.0*VN(I,J)+VN(I,J-1)))
    UDR=0.5*(UD(I,J)+UD(I+1,J))*(A(I,J)*A(I+1,J)/(A(I,J)+A(I+1,J))+
    1A(I+1,J+1)*A(I+1,J+1)/(A(I+1,J+1)+A(I+1,J+1)))
    UDL=0.5*(UD(I-1,J)+UD(I-1,J+1))*(A(I,J)*A(I-1,J)/(A(I,J)+A(I-1,J))+
    1A(I+1,J+1)*A(I-1,J+1)/(A(I+1,J+1)+A(I-1,J+1)))
    RVT=0.5*RVN(I,J+1)*(RON(I,J+1)-RVN(I,J+1))/RON(I,J+1)*(VD(I,J)*
    1*VD(I,J+1)*A(I,J+1))
    RVB=0.5*RVN(I,J)*(RON(I,J)-RVN(I,J))/RON(I,J)*(VD(I,J-1)+VD(I,J))-
    1*A(I,J)

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532 276

71

POOR ORIGINAL

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RVR=0.5*VD(I+1,J)*(RVN(I+1,J)*(RON(I+1,J)-RVN(I+1,J))/RON(I+1,J)*
1RVN(I+1,J+1)*(RON(I+1,J+1)-RVN(I+1,J+1))/RON(I+1,J+1))
RVC=0.5*VD(I,J)*(RVN(I,J)*(RON(I,J)-RVN(I,J))/RON(I,J)+RVN(I,J+1)*
1*(RON(I,J+1)-RVN(I,J+1))/RON(I,J+1))
RVL=0.5*VD(I-1,J)*(RVN(I-1,J)*(RON(I-1,J)-RVN(I-1,J))/RON(I-1,J)*
1RVN(I-1,J+1)*(RON(I-1,J+1)-RVN(I-1,J+1))/RON(I-1,J+1))
FDV=AVBR*(RDX*(UDR*(RVC+RVR)+ALPHA*ABS(UDR)*(RVC-RVR)-UDL*(RVL+RVC
1)-ALPHA*ABS(UDL)*(RVL-RVC))+RDY*(RVT*(VD(I,J)+VD(I,J+1))+*
2ALPHA*ABS(RVT)*(VD(I,J)-VD(I,J+1))-RVB*(VD(I,J-1)+VD(I,J))-*
3ALPHA*ABS(RVB)*(VD(I,J-1)-VD(I,J)))/(RON(I,J)+RON(I,J+1))
VISY=0.0
V(I,J)=VN(I,J)+DELT*(2.0*(P(I,J)-P(I,J+1))*RDY/(RO(I,J)+RO(I,J+1))+
1+GY-FVX-FVY+VISY-FDV)
65 CONTINUE
C ADD PIPE FRICTION
IF(RPIPE.EQ.0.0.OR.DIM.GT.1.5)GO TO 70
I|=I
JJ=J
CALL PFRIC
70 CONTINUE
CALL BC
250 CONTINUE
C HAS CONVERGENCE BEEN REACHED
IF(FLG.EQ.0.150 TO 400
ITER=ITER+1
IF(ITER.LT.50)GO TO 255
NEX=NEX+1
IF(NEX.LT.1000) GO TO 400
T=1000000000.0
GO TO 502
C COMPUTE UPDATED CELL VELOCITIES U,V
255 CONTINUE
CALL PITER
CALL BC
GO TO 250
400 CONTINUE
C COMPUTE UPDATED QUANTITIES RO, E, RV
D3 4500 J=2,IM1
D0 4500 I=2,IM1
IF(A(I,J).EQ.0.0)GO TO 4500
ABR=2.0*A(I,J)*A(I+1,J)/A(I,J)+A(I+1,J)
ABL=2.0*A(I,J)*A(I-1,J)/A(I,J)+A(I-1,J)
ABT=2.0*A(I,J)*A(I,J+1)/A(I,J)+A(I,J+1)
ABB=2.0*A(I,J)*A(I,J-1)/A(I,J)+A(I,J-1)
C DENSITY EQUATION
ULR=U(I,J)-UD(I,J)*(RVN(I,J)+RVN(I+1,J))/(RON(I,J)+RON(I+1,J))
ULL=U(I+1,J)-UD(I+1,J)*(RVN(I+1,J)+RVN(I,J))/(RON(I+1,J)+RON(I,J))
ULT=V(I,J)-VD(I,J)*(RVN(I,J)+RVN(I,J+1))/(RON(I,J)+RON(I,J+1))
ULB=V(I,J)-VD(I,J-1)*(RVN(I,J-1)+RVN(I,J))/(RON(I,J-1)+RON(I,J))
UVR=U(I,J)*UD(I,J)*(RON(I,J)+RON(I+1,J)-RVN(I,J)-RVN(I+1,J))/
1*(RON(I,J)+RON(I+1,J))
UVL=U(I+1,J)*UD(I+1,J)*(RON(I+1,J)+RON(I,J)-FVN(I+1,J)-RVN(I,J))/
1*(RON(I+1,J)+RON(I,J))
UVT=V(I,J)*VD(I,J)*(RON(I,J)+RON(I,J+1)-RVN(I,J)-RVN(I,J+1))/
1*(RON(I,J)+RON(I,J+1))
UVB=V(I,J-1)*VD(I,J-1)*(RON(I,J-1)+RON(I,J)-RVN(I,J)-RVN(I,J))/
1*(RON(I,J-1)+RON(I,J))
FLR=(ULR*(RON(I,J)+RVN(I,J)+RON(I+1,J)+RVN(I+1,J))+ALPHA*ABS(ULR)*
1+GY-FVX-FVY+VISY-FDV)
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1 (RON(1,J)-RVN(1,J)-RON(1+1,J)+RVN(1+1,J)) * ABR      SOL ADF 467
FLL=(ULL*(RON(1,J)-RVN(1,J)+RON(1,J)-RVN(1,J)) * ALPHA * ABS(ULL) * SOL ADF 468
1 (RON(1,J)-RVN(1,J)-RON(1,J)+RVN(1,J)) * ABL      SOL ADF 469
FLT=(ULT*(RON(1,J)-RVN(1,J)+RON(1,J+1)-RVN(1,J+1)) * ALPHA * ABS(ULT) * SOL ADF 470
1 (RON(1,J)-RVN(1,J)-RON(1,J+1)+RVN(1,J+1)) * ABT      SOL ADF 471
FLB=(ULB*(RON(1,J-1)-RVN(1,J-1)+RON(1,J)-RVN(1,J)) * ALPHA * ABS(ULB) * SOL ADF 472
1 (RON(1,J-1)-RVN(1,J-1)-RON(1,J)+RVN(1,J)) * ABB      SOL ADF 473
FVR=(UVR*(RVN(1,J)+RVN(1+1,J)) * ALPHA * ABS(JVR) * (RVN(1,J)-
IRVN(1+1,J))) * ABR      SOL ADF 474
FVL=(UVL*(RVN(1,J)+RVN(1,J)) * ALPHA * ABS(UVL) * (RVN(1,J)-
IRVN(1,J))) * ABL      SOL ADF 475
FVT=(UVT*(RVN(1,J)+RVN(1,J+1)) * ALPHA * ABS(UVT) * (RVN(1,J)-
IRVN(1,J+1))) * ABT      SOL ADF 476
FVB=(UVB*(RVN(1,J-1)+RVN(1,J)) * ALPHA * ABS(UVB) * (RVN(1,J-1)-
IRVN(1,J))) * ABB      SOL ADF 477
C      UPDATE RO
RO(1,J)=RON(1,J)-0.5*DELT*((FLR+FVR-FLL-FVL)*ROX+(FLT+FVT-FLB-
IFVB)*RDY)/A(1,J)      SOL ADF 478
C      ENERGY EQUATION
ROEC=RON(1,J)*EN(1,J)      SOL ADF 479
ROER=RON(1+1,J)*EN(1+1,J)      SOL ADF 480
ROEL=RON(1-1,J)*EN(1-1,J)      SOL ADF 481
ROET=RON(1,J+1)*EN(1,J+1)      SOL ADF 482
ROEB=RON(1,J-1)*EN(1,J-1)      SOL ADF 483
IF(ETEM.GT.0.5)GO TO 425      SOL ADF 484
TVC=SATT(P(1,J))      SOL ADF 485
TVP=SATT(P(1+1,J))      SOL ADF 486
TVL=SATT(P(1-1,J))      SOL ADF 487
TVT=S_TT(P(1,J+1))      SOL ADF 488
TVB=SA_T(P(1,J-1))      SOL ADF 489
GO TO 410      SOL ADF 490
425 TVC=TC*(ROEC-RON(1,J)*ECL-RVN(1,J)*(ECV-ECL))/(RON(1,J)*CHL+
IRVN(1,J)*(CHV-CHL))      SOL ADF 491
TVR=TC*(ROER-RON(1+1,J)*ECL-RVN(1+1,J)*(ECV-ECL))/(RON(1+1,J)*
CHL+RVN(1+1,J)*(CHV-CHL))      SOL ADF 492
TVL=TC*(ROEL-RON(1-1,J)*ECL-RVN(1-1,J)*(ECV-ECL))/(RON(1-1,J)*
CHL+RVN(1-1,J)*(CHV-CHL))      SOL ADF 493
TVT=TC*(ROET-RON(1,J+1)*ECL-RVN(1,J+1)*(ECV-ECL))/(RON(1,J+1)*
CHL+RVN(1,J+1)*(CHV-CHL))      SOL ADF 494
TVB=TC*(ROEB-RON(1,J-1)*ECL-RVN(1,J-1)*(ECV-ECL))/(RON(1,J-1)*
CHL+RVN(1,J-1)*(CHV-CHL))      SOL ADF 495
430 CONTINUE
REVC=RVN(1,J)*EVCAL(TVC)      SOL ADF 496
REVR=RVN(1+1,J)*EVCAL(TVR)      SOL ADF 497
REVL=RVN(1-1,J)*EVCAL(TVL)      SOL ADF 498
REVt=RVN(1,J+1)*EVCAL(TVT)      SOL ADF 499
REVb=RVN(1,J-1)*EVCAL(TVB)      SOL ADF 500
RELC=ROEC-REVC      SOL ADF 501
RELr=ROER-REVR      SOL ADF 502
RELL=ROEL-REVL      SOL ADF 503
RELT=ROET-REVT      SOL ADF 504
RELb=ROEB-REVb      SOL ADF 505
FRER=0.5*(UVR*(REVC+REVR)+ALPHA*ABS(UVR)*(REVC-REVR) +
ULR*(RELC+RELr)+ALPHA*ABS(ULR)*(RELC-RELr))*ABR      SOL ADF 506
FREL=0.5*(UVL*(REVL+REVC)+ALPHA*ABS(UVL)*(REVL-REVC)+ULL*(
RELL+RELC)+ALPHA*ABS(ULL)*(RELL-RELC))*ABL      SOL ADF 507
FRET=0.5*(UVT*(REVC+REVt)+ALPHA*ABS(UVT)*(REVC-REVT)+ULT*(RELC*(
RELT))+ALPHA*ABS(UU)*(RELC-RELT))*ABT      SOL ADF 508

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POOR ORIGINAL

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FREB=0.5*(UVB*(REV6+REVC)+ALPHA*ABS(UVB)*(REV6-REVC)+ULB*(RELB+
1RELC)+ALPHA*ABS(ULB)*(RELB-RELC))*ABB
FEL=(RDX*(FRER-FR))**2+RDY*(FRET-FREB))/A(I,J)
FPW=P(I,J)*((ABR*U(I,J)-ABL*U(I+1,J))*RDX*(ABT*V(I,J)-
1ABB*V(I,J-1))*RDY)/A(I,J)
TH=(ROL+RVN(I,J)-RON(I,J))/ROL
TH=AMAX1(TH,THC)
TH=AMIN1(TH,THC1)
THI=TH
IF(THI.GT.0.5)THI=1.0-THI
440 RB=(THI/PNV)**0.333
CKN=0.5*CDG*((UD(I,J)+UD(I-1,J))**2+(VD(I,J)+VD(I,J-1))**2)**0.5
1*12.0*V1SV/RB
FDIS=0.25*CKN*(0.375*TH1/RB)*((UD(I,J)+UD(I-1,J))**2+(VD(I,J)+
1VD(I,J-1))**2)*R0(I,J)
FDP=-P(I,J)*((RVN(I,J)+RVN(I+1,J)-RON(I,J)-RON(I+1,J)+
12.0*ROL)/(ROL*2.01-(RVN(I,J)+RVN(I+1,J))/(RON(I,J)+RON(I+1,J)))*
1ABR*
2UD(I,J)-(RVN(I-1,J)+RVN(I,J)-RON(I-1,J)-RON(I,J)+2.0*ROL)/(2.0*
3ROL)-(RVN(I-1,J)+RVN(I,J))/(RON(I-1,J)+RON(I,J)))*UD(I-1,J)
3*ABL)/DELX+
4*((RVN(I,J)+RVN(I,J+1)-RON(I,J)-RON(I,J+1)*2.0*ROL)/(2.0*ROL)
5-(RVN(I,J)+RVN(I,J+1))/(RON(I,J)+RON(I,J+1))*VD(I,J)*ABT
5-((RVN(I,J-1)
5+RVN(I,J)-
6RON(I,J-1)-RON(I,J)+2.0*ROL)/(2.0*ROL)+(RVN(I,J-1)+RVN(I,J))/
7(RON(I,J-1)+RON(I,J+1))*VD(I,J-1)*ABB)/DELY)/A(I,J)
C ENERGY DIFFUSION
TEM=EN(I,J)
TEMR=EN(I+1,J)
TEML=EN(I-1,J)
TEMU=EN(I,J+1)
TEMB=EN(I,J-1)
THR=(ROL+RVN(I+1,J)-RON(I+1,J))/ROL
THR=AMAX1(THR,THC)
THR=AMIN1(THR,THC1)
THL=(ROL+RVN(I-1,J)-RON(I-1,J))/ROL
THL=AMAX1(THL,THC)
THL=AMIN1(THL,THC1)
THT=(ROL+RVN(I,J+1)-RON(I,J+1))/ROL
THT=AMAX1(THT,THC)
THT=AMIN1(THT,THC1)
THB=(ROL+RVN(I,J-1)-RON(I,J-1))/ROL
THB=AMAX1(THB,THC)
THB=AMIN1(THB,THC1)
EKR=0.5*(EDV*(TH+THR)+EDL*(2.0-TH-THR))*ABR
EKT=0.5*(EDV*(TH+THT)+EDL*(2.0-TH-THT))*ABT
EKL=0.5*(EDV*(TH+THL)+EDL*(2.0-TH-THL))*ABL
EKB=0.5*(EDV*(TH+THB)+EDL*(2.0-TH-THB))*ABB
DIFE=(RDX*RDX*(EKR*(TEMR-TEM)+EKL*(TEM-TEML))+
1RDY*RDY*(EKT*(TEMU-TEM)+EKB*(TEM-TEMB)))/A(I,J)
C UPDATE ENERGY
E(I,J)=(ROEC+DELT*(-FEC+FDIS+FDP+DIFE+FPW))/ROL(I,J)
C VAPOR DENSITY EQUATION
RV(I,J)=RVN(I,J)-DELT*0.5*(FVR-FVL)*RDX*(FVT-FVB)*
1RDY/A(I,J)
RV(I,J)=AMIN1(RV(I,J),ROL(I,J))
RV(I,J)=AMAX1(RV(I,J),0.01)

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POOR ORIGINAL

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C   PHASE CHANGE                               SOLADF 581
    IF(PHCH.LT.0.5)GO TO 4500                 SOLADF 582
    II=I                                         SOLADF 583
    JJ=J                                         SOLADF 584
    CALL PHCHR                                    SOLADF 585
4500 CONTINUE                                 SOLADF 586
    CALL BC                                       SOLADF 587
4600 CONTINUE                                 SOLADF 588
C   COMPUTE DRIFT VELOCITY                   SOLADF 589
    IF(DFVEL.LT.0.5)GO TO 502                  SOLADF 590
    CALL DRIFT                                     SOLADF 591
500 CONTINUE                                 SOLADF 592
    CALL BC                                       SOLADF 593
502 CONTINUE                                 SOLADF 594
C   PRINT AND PLOT                            SOLADF 595
    IF(T.GT.0.)GO TO 503                      SOLADF 596
    CALL LINCNT(1)                                SOLADF 597
    WRITE(12,45)                                  SOLADF 598
503 CONTINUE                                 SOLADF 599
C   DEFINE OUTPUT QUANTITY F AND F1          SOLADF 600
    F=0.0                                         SOLADF 601
    INZ=4                                         SOLADF 602
    J=20                                         SOLADF 603
    DO 505 :=2,5                                SOLADF 604
505 F=F+RO(I,J)*V(I,J)*(FLOAT(I)-1.5)
    F=2.0*F/INZ**2                             SOLADF 605
    F1=0.34264*P(5,11)+0.65736*P(5,12)
    WRITE(9,49)ITER,T,CYCLE,DELT,F,F1
    IF(CYCLE.LE.0)GO TO 510
    IF(T.LT.TWPLT)GO TO 560
    TWPLT=TWPLT+PLTDT
510 CONTINUE                                 SOLADF 606
C   FILM LONG PRINT                           SOLADF 607
    CALL ADV(1)                                   SOLADF 608
    CALL LINCNT(3)                                SOLADF 609
    WRITE(12,49)ITER,T,CYCLE,DELT,F,F1
    CALL LINCNT(4)                                SOLADF 610
    WRITE(12,47)                                  SOLADF 611
    DO 525 J=2,IM1                               SOLADF 612
    DO 525 I=2,IM1                               SOLADF 613
    TH=(ROL-RO(I,J)+RV(I,J))/ROL
    D=RDY*(U(I,J)-U(I-1,J))+RDY*(V(I,J)-V(I,J-1))+CYL*(U(I,J)
    +U(I-1,J))/(2.*DELX*(FLOAT(I)-1.5))
    TSAT=SATT(P(I,J))
    EVT=EVCAL(TSAT)*ETEM1+ECV*ETEM
    TEM=TC+(RO(I,J)*(E(I,J)-ECL)-RV(I,J)*(EVT-ECL))/(RV(I,J)*(CHV*
    IETEM-CHL)+RO(I,J)*CHL)
    WRITE(12,48) (I,J,U(I,J),V(I,J),P(I,J),RO(I,J),TH,TEM)
525 CONTINUE                                 SOLADF 614
    CALL ADV(1)                                   SOLADF 615
    CALL LINCNT(3)                                SOLADF 616
    WRITE(12,49)ITER,T,CYCLE,DELT,F,F1
    CALL LINCNT(4)                                SOLADF 617
    WRITE(12,5)                                    SOLADF 618
    DO 528 J=2,IM1                               SOLADF 619
    DO 528 I=2,IM1                               SOLADF 620
    WRITE(12,10)(I,J,UD(I,J),VD(I,J),RV(I,J))
528 CONTINUE                                 SOLADF 621

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POOR ORIGINAL

532 280

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C   VELOCITY VECTOR PLOT                               639
    CALL ADV(1)                                         640
    CALL DGA((XL,IXR,IYT,IYB,0.,LONG,HIGH,0.))        641
    CALL FRAME((XL,IXR,IYT,IYB))                      642
    CALL FRAME((XL,IXR,IYT,IYB))                      643
    CALL LINCNT(61)                                     644
    WRITE(12,45)                                       645
    CALL LINCNT(61)                                     646
    WRITE(12,46,T,CYCLE)                             647
    DO 530 J=2,JM1                                    648
    DO 530 I=2,IM1                                    649
    XCC=DELX*(FLOAT(I)-1.5)                          650
    YCC=DELY*(FLOAT(J)-1.5)                          651
    UVEC=(U(I-1,J)+U(I,J))*0.5*VELMX1+XCC          652
    VVEC=(V(I,J-1)+V(I,J))*0.5*VELMX1+YCC          653
    CALL CONVRT(UVEC,IUVEC,0.,LONG,IXL,IXR)           654
    CALL CONVRT(VVEC,JVVEC,HIGH,0.,IYT,IYB)           655
    CALL CONVRT(XCC,IXCC,0.,LONG,IXL,IXR)             656
    CALL CONVRT(YCC,JYCC,HIGH,0.,IYT,IYB)             657
    CALL DRV((XCC,JYCC,IUVEC,JVVEC))                 658
    530 CONTINUE                                       659
C   CONTOURS PLOT STARTS                           660
    NCO=0                                            661
    532 NCO=NCO+1                                    662
    DO 545 I=1,IM1                                    663
    DO 545 J=1,JM1                                    664
    GO TO(533,534,535,536,537,560),NCO            665
    533 Q(I,J)=RO(I+1,J+1)                          666
    GO TO 540                                         667
    534 Q(I,J)=E(I+1,J+1)                          668
    GO TO 540                                         669
    535 Q(I,J)=P(I+1,J+1)                          670
    GO TO 540                                         671
    536 Q(I,J)=RV(I+1,J+1)                          672
    GO TO 540                                         673
    537 Q(I,J)=(ROL+RV(I+1,J+1)-RO(I+1,J+1))/ROL 674
    540 QL(J)=Q(I,J)                                675
    545 CONTINUE                                       676
    CALL LINCNT(61)                                     677
    WRITE(12,45)                                       678
    CALL LINCNT(61)                                     679
    WRITE(12,46)T,CYCLE)                            680
    GO TO (553,554,555,556,557),NCO                681
    553 ITITLE(1)=10HRC CONTOUR                   682
    GO TO 559                                         683
    554 ITITLE(1)=9HE CONTOUR                     684
    GO TO 559                                         685
    555 ITITLE(1)=9HP CONTOUR                     686
    GO TO 559                                         687
    556 ITITLE(1)=10HRV CONTOUR                   688
    GO TO 559                                         689
    557 ITITLE(1)=9HVOID FRAC                     690
    559 CONTINUE                                       691
    IF((IBAR.GT.1)GO TO 5601
    CALL SPLOT(1,IBAR,YC(2),QL(2),42,1)
    GO TO 532
    5601 CONTINUE                                     692
    CALL CONTRIB(XC(2),NNX,YC(2),NNY,Q,NZX,NZY,NC,ZMN,ZMX,DLZ,ZC, 693
                                                694
                                                695
                                                696

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1DMPX,DMPY,TGRD,ITITLE,NTITLE,XLABEL,NXLBL,YLABEL,NYLBL)      SOLADF 697
GO TO 532
560 CONTINUE
C     LONG PRINT
IF(CYCLE.LE.0) GO TO 565
IF(T.LT.TWPRT)GO TO 580
TWPRT=TWPRT+PRTDT
565 CONTINUE
WRITE(9,35)
WRITE(9,49)ITER,T,CYCLE,DELT,F,F1
WRITE(9,47)
DO 575 J=2,JM1
DO 575 I=2,IM1
TH=(ROL-RO(I,J)+RV(I,J))/ROL
TSAT=SATT(P(I,J))
EVT=EVCAL(TSAT)*ETEM1+ECV*ETEM
TEM=TC +(RO(I,J)*(E(I,J)-ECL)-RV(I,J)*(EVT-ECL))/(RV(I,J)*(CHV*
IETEM-CHL)+RO(I,J)*CHL)
WRITE(9,48)(I,J,U(I,J),V(I,J),P(I,J),RO(I,J),TH,TEM)
575 CONTINUE
C     SET THE ADVANCE TIME QUANTITIES INTO OLD ARRAYS
580 CONTINUE
DO 600 J=1,JMAX
DO 600 I=1,IMAX
UN(I,J)=U(I,J)
VN(I,J)=V(I,J)
RON(I,J)=ROT=RO(I,J)
RVN(I,J)=RVT=RV(I,J)
EN(I,J)=ET=E(I,J)
IF(IMP.GT.0.5)GO TO 600
II=I
JJ=J
CALL EOS
P(I,J)=PT
600 CONTINUE
C     ADJUST TIME STEP
CALL TSTEP
C     ADVANCE TIME T=T+DELT
IF(T.GT.THFIN)GO TO 550
T=T+DELT
CYCLE=CYCLE+1
GO TO 59
650 CALL EXIT()
END

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SUBROUTINE BC                               SOLADF 741
*CALL COMDK                                SOLADF 742
C     SET BOUNDARY CONDITION                 SOLADF 743
EVCAL(X)=ECV+CHV*(X-TC)                   SOLADF 744
ELCAL(X)=ECL+CHL*(X-TC)                   SOLADF 745
SATT(X)=255.2+117.8*X**0.223            SOLADF 746
SATP(X)=((X-255.2)/117.8)**4.48          SOLADF 747
DO 140 J=2,JM1                            SOLADF 749

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POOR ORIGINAL

532 77282

RO(1,J)=RO(2,J)	SOL ADF	749
RV(1,J)=RV(2,J)	SOL ADF	750
P(1,J)=P(2,J)	SOL ADF	751
E(1,J)=E(2,J)	SOL ADF	752
RO(1MAX,J)=RO(1M1,J)	SOL ADF	753
RV(1MAX,J)=RV(1M1,J)	SOL ADF	754
P(1MAX,J)=P(1M1,J)	SOL ADF	755
E(1MAX,J)=E(1M1,J)	SOL ADF	756
GO TO (102,104,106,108,106),WL	SOL ADF	757
102 U(1,J)=0.0	SOL ADF	758
V(1,J)=V(2,J)	SOL ADF	759
GO TO 111	SOL ADF	760
104 U(1,J)=0.0	SOL ADF	761
V(1,J)=-V(2,J)	SOL ADF	762
GO TO 111	SOL ADF	763
106 IF(WL.EQ.5)P(2,J)=PBC	SOL ADF	764
IF(ITER.GT.0.AND.FLG.GT.0)GO TO 111	SOL ADF	765
U(1,J)=U(2,J)	SOL ADF	766
V(1,J)=V(2,J)	SOL ADF	767
GO TO 111	SOL ADF	768
108 U(1,J)=U(1M2,J)	SOL ADF	769
V(1,J)=V(1M2,J)	SOL ADF	770
RO(1,J)=RO(1M2,J)	SOL ADF	771
RV(1,J)=RV(1M2,J)	SOL ADF	772
P(1,J)=P(1M2,J)	SOL ADF	773
E(1,J)=E(1M2,J)	SOL ADF	774
GO TO 111	SOL ADF	775
111 GO TO (122,124,126,128,126),WR	SOL ADF	776
122 U(1M1,J)=0.0	SOL ADF	777
V(1MAX,J)=V(1M1,J)	SOL ADF	778
GO TO 140	SOL ADF	779
124 U(1M1,J)=0.0	SOL ADF	780
V(1MAX,J)=-V(1M1,J)	SOL ADF	781
GO TO 140	SOL ADF	782
126 IF(WR.EQ.5)P(1M1,J)=PBC	SOL ADF	783
IF(ITER.GT.0.AND.FLG.GT.0)GO TO 140	SOL ADF	784
U(1M1,J)=U(1M2,J)*((1M2-1)/(1M1-1)*CYL+(1.0-CYL))	SOL ADF	785
V(1MAX,J)=V(1M1,J)	SOL ADF	786
GO TO 140	SOL ADF	787
128 U(1M1,J)=U(2,J)	SOL ADF	788
V(1MAX,J)=V(3,J)	SOL ADF	789
RO(1MAX,J)=RO(3,J)	SOL ADF	790
RV(1MAX,J)=RV(3,J)	SOL ADF	791
P(1MAX,J)=P(3,J)	SOL ADF	792
E(1MAX,J)=E(3,J)	SOL ADF	793
140 CONTINUE	SOL ADF	794
DO 180 I=2,1M1	SOL ADF	795
RO(I,J)=RO(I,2)	SOL ADF	796
RV(I,J)=RV(I,2)	SOL ADF	797
P(I,J)=P(I,2)	SOL ADF	798
E(I,J)=E(I,2)	SOL ADF	799
RO(I,JMAX)=RO(I,JM1)	SOL ADF	800
RV(I,JMAX)=RV(I,JM1)	SOL ADF	801
P(I,JMAX)=P(I,JM1)	SOL ADF	802
E(I,JMAX)=E(I,JM1)	SOL ADF	803
GO TO (152,154,156,158,156),WT	SOL ADF	804
152 V(1,JM1)=0.0	SOL ADF	805
U(1,JMAX)=U(1,JM1)	SOL ADF	806

POOR ORIGINAL

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      GO TO 161
154 V(I,JM1)=0.0
      U(I,JMAX)=-U(I,JM1)
      GO TO 161
156 IF(WT.EQ.5)P(I,JM1)=PBC
      IF(ITER.GT.0.AND.FLG.GT.0)GO TO 161
      V(I,JM1)=V(I,JM2)
      U(I,JMAX)=U(I,JM1)
      GO TO 161
158 V(I,JM1)=V(I,2)
      U(I,JMAX)=U(I,3)
      RO(I,JMAX)=RO(I,3)
      RV(I,JMAX)=RV(I,3)
      P(I,JMAX)=P(I,3)
      E(I,JMAX)=E(I,3)
      GO TO 161
161 GO TO (172,174,176,178,179),WB
172 V(I,1)*0.0
      U(I,1)=U(I,2)
      GO TO 180
174 V(I,1)=0.0
      U(I,1)=-U(I,2)
      GO TO 180
175 IF(WB.EQ.5)P(I,2)=PBC
      IF(ITER.GT.0.AND.FLG.GT.0)GO TO 180
      V(I,1)=V(I,2)
      U(I,1)=U(I,2)
      GO TO 180
178 V(I,1)=V(I,JM2)
      U(I,1)=U(I,JM2)
      RO(I,1)=RO(I,JM2)
      RV(I,1)=RV(I,JM2)
      P(I,1)=P(I,JM2)
      E(I,1)=E(I,JM2)
180 CONTINUE
182 CONTINUE
C   SPECIAL BOUNDARY CONDITION
DO 188 I=2,IM1
P(I,2)=PIN
U(I,2)=0.0
U(I,JM1)=0.0
IF(THIN.LT.THC) GO TO 38
IF(THIN.LT.THC1) GO TO 36
RV(I,2)=RV1/V
RO(I,2)=RO1/V
E(I,2)=E1/V
GO TO 188
36 RV(I,2)=RV12
RO(I,2)=RO12
E(I,2)=E12
GO TO 188
38 RV(I,2)=RV1L
RO(I,2)=RO1L
E(I,2)=E1L
188 CONTINUE
RETURN
END
      SOL.ADF 807
      SOL.ADF 808
      SOL.ADF 809
      SOL.ADF 810
      SOL.ADF 811
      SOL.ADF 812
      SOL.ADF 813
      SOL.ADF 814
      SOL.DF 815
      SOL.ADF 816
      SOL.ADF 817
      SOL.ADF 818
      SOL.ADF 819
      SOL.ADF 820
      SOL.ADF 821
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      SOL.ADF 849
      SOL.ADF 850
      SOL.ADF 851
      SOL.ADF 852
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      SOL.ADF 856
      SOL.ADF 857
      SOL.ADF 858
      SOL.ADF 859
      SOL.ADF 860
      SOL.ADF 861
      SOL.ADF 862
      SOL.ADF 863

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POOR ORIGINAL

532 79
284

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SUBROUTINE DRIFT
*CALL COMDK                               SOLADF 864
C   CALCULATE RELATIVE OR DRIFT VELOCITY      SOLADF 865
EVCAL(X)=ECV+CHV*(X-TC)                   SOLADF 866
ELCAL(X)=ECL+CHL*(X-TC)                   SOLADF 867
SATT(X)=255.2+117.8*X**0.223             SOLADF 868
SATP(X)=((X-255.2)/117.8)**4.48          SOLADF 869
DO 5000 J=2,IM1                           SOLADF 870
DO 5000 I=2,IM1                           SOLADF 871
IF(A(I,J).EQ.0.0)GO TO 5000              SOLADF 872
TH=(ROL-RO(I,J)+RV(I,J))/ROL            SOLADF 873
IF(TH.GT.THC.AND.TH.LT.THC1)GO TO 4700    SOLADF 874
UD(I,J)=VD(I,J)=0.0                      SOLADF 875
GO TO 5000                                 SOLADF 876
4700 CONTINUE                                SOLADF 877
TH1=TH                                     SOLADF 878
IF(TH1.GT.0.5)TH1=1.0-TH1                  SOLADF 879
4800 RB=(TH1/PNV)**0.333                 SOLADF 880
UDN=UD(I,J)                                SOLADF 881
VDN=VD(I,J)                                SOLADF 882
AREA=3.0*TH1/RB                            SOLADF 883
NUC=VISL                                    SOLADF 884
IF(TH1.GT.0.5)NUC=VISV                    SOLADF 885
NUA=NUC/(1.0-TH1)**2.5                     SOLADF 886
UDO=DELT*(1.0/ROL-2.0*TH/(RV(I,J)+RV(I+1,J)))*RDX*(P(I+1,J)-IP(I,J))
IF(A(I+1,J).EQ.0.0)UDO=0.0                SOLADF 887
VDO=DELT*(1.0/ROL-2.0*TH/(RV(I,J)+RV(I,J+1)))*RDY*(P(I,J+1)-IP(I,J))
IF(A(I,J+1).EQ.0.0)VDO=0.0                SOLADF 888
UDM=(UD(I,J)**2+VD(I,J)**2)**0.5+1.0E-10  SOLADF 889
UKPM=DELT*0.25*AREA*(RO(I,J)+RO(I+1,J))**2/((RV(I,J)+RV(I+1,J))*
  (RO(I,J)+RO(I+1,J))-RV(I,J)-RV(I+1,J)))
VKPM=DELT*0.25*AREA*(RO(I,J)+RO(I,J+1))**2/((RV(I,J)+RV(I,J+1))*
  (RO(I,J)+RO(I,J+1))-RV(I,J)-RV(I,J+1))
4950 UDMT=UDM                               SOLADF 890
UKP=UKPM*(CDG*UDM+12.0*NJA/RB)/TH1       SOLADF 891
VKP=VKPM*(CDG*UDM+12.0*NJA/RB)/TH1       SOLADF 892
UD(I,J)=UD(I,J)-(UD(I,J)+UKP*UD(I,J)-UDO-UDN)/(1.0+UKP+UKPM*CDG)
  *UD(I,J)**2/UDM                           SOLADF 893
VD(I,J)=VD(I,J)-(VD(I,J)+VKP*VD(I,J)-VDO-VDN)/(1.0+VKP+VKPM*CDG)
  *VD(I,J)**2/UDM                           SOLADF 894
UDM=(UDO(I,J)**2+VD(I,J)**2)**0.5+1.0E-10  SOLADF 895
IF(ABS((UDMT-UDM)/(UDMT+UDM)).GT.0.01)GO TO 4950  SOLADF 896
5000 CONTINUE                                SOLADF 897
RETURN                                     SOLADF 898
END                                         SOLADF 899
SOLADF 900
SOLADF 901
SOLADF 902
SOLADF 903
SOLADF 904
SOLADF 905
SOLADF 906
SOLADF 907
SOLADF 908
SOLADF 909
SOLADF 910

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***** / / / / *****

SUBROUTINE EOS
*CALL COMDK                               SOLADF 911
C   EQUATION OF STATE ROUTINE              SOLADF 912
EVCAL(X)=ECV+CHV*(X-TC)                   SOLADF 913
ELCAL(X)=ECL+CHL*(X-TC)                   SOLADF 914
SOLADF 915

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SATT(X)=255.2+117.8*X**0.223          SOL ADF 916
SATP(X)=((X-255.2)/117.8)**4.48        SOL ADF 917
TH=THTM=(ROL+RVT+ROT)/ROL             SOL ADF 918
IF(TH.LT.THC)THTM=THC                 SOL ADF 919
ETEMT=ETEM                           SOL ADF 920
IF(TH.GT.THC)ETEMT=1.0                  SOL ADF 921
TEM=TC+(ROT*(ET-ECL)-RVT*(ECV-ECL))/((RVT*(CHV-CHL)+ROT*CHL)) SOL ADF 922
IF(ETEMT.LT.0.5)TEM=255.2               SOL ADF 923
PT=PTT=GAM1*RVT*(ECV+CHV*(TEM-TC))/THTM+ASQ*ROL*(THTM-TH) SOL ADF 924
IF(ETEMT.GT.0.5)GO TO 6300             SOL ADF 925
PCC=117.8*GAM1*RVT*CHV/THTM           SOL ADF 926
PT=PTT+2.0*(0.223*PCC)**1.287         SOL ADF 927
6260 PTG=PT
    PTA=PCC*PTG**0.223                SOL ADF 928
    PT=PTG+(PTT+PTA-PTG)/(1.0-0.223*PTA/PTG) SOL ADF 929
    IF(ABS((PT-PTG)/(PT+PTG)).GT.0.01)GO TO 6260 SOL ADF 930
    PT=AMAX1(PT,0.0)                   SOL ADF 931
    IF(PT.GT.0.0)GO TO 6300            SOL ADF 932
    WRITE(9,6200)                      SOL ADF 933
6200 FORMAT(5X,1HJ,12X,2HPT,12X,3HROT,17X,3HRVT,17X,2HET)
    WRITE(9,6263)I,J,PT,ROT,RVT,ET      SOL ADF 934
6263 FORMAT(5X,213.5X,1PE12.5,6X,E12.5,6X,E12.5,6X,E12.5) SOL ADF 935
6300 CONTINUE
    RETURN
    END
                                         SOL ADF 936
                                         SOL ADF 937
                                         SOL ADF 938
                                         SOL ADF 939
                                         SOL ADF 940

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SUBROUTINE PFRIC
*CALL COMDK
C PIPE FRICTION ROUTINE
EVCAL(X)=ECV+CHV*(X-TC)
ELCAL(X)=ECL+CHL*(X-TC)
SATT(X)=255.2+117.8*X**0.223
SATP(X)=((X-255.2)/117.8)**4.48
I=1
J=JU
VA1=ABS(V(I,J))
IF(VA1.LT.EPS1)GO TO 60
TH=TH=(ROL+RV(I,J)-RO(I,J))/ROL
IF(TH.LT.THC)TH=THC
IF(TH.GT.THC)TH=THC
CH1=RV(I,J)/RO(I,J)*(1.0+(RO(I,J)-RV(I,J))/RO(I,J))*VD(I,J)/V(I,J)
REN=2.0*RP1PE*ABS(V(I,J))/VISL
RGR=0.5*RG/RP1PE
FLA=0.026*RGR**0.225+0.133*RGR
FLB=22.0*RGR**0.44
FLC=1.62*RGR**0.134
FLC2=(FLA+FLB/REN**FLC)
PHIS=1.0/(1.0-TH)**1.75

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532 286

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FLC=FLC2*(R0(I,J)/ROL)*(1.0-CHI)**2.0/RPIPE*PHIS          SOLADF 963
FLCT=2.0*DELT*FLC                                         SOLADF 964
V(I,J)=SIGN(1.0,V(I,J))*(-1.0+(1.0+2.0*FLCT*VAL)**0.5)/FLCT  SOLADF 965
EE=0.25*(VA)**2-V(I,J)**2                                  SOLADF 966
E(I,J)=E(I,J)+EE                                         SOLADF 967
E(I,J+1)=E(I,J+1)+EE                                     SOLADF 968
60 CONTINUE
RETURN
END

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***** = = = = = / / / / / = = = = = *****

SUBROUTINE PHCHR                                         SOLADF 972
*CALL COMDK
C   PHASE CHANGE ROUTINE
EVCAL(X)=ECV+CHV*(X-TC)                                 SOLADF 973
ELCAL(X)=ECL+CHL*(X-TC)                                 SOLADF 974
SATT(X)=255.2+117.8*X**0.223                            SOLADF 975
SATP(X)=((X-255.2)/117.8)**4.48                         SOLADF 976
I=11
J=JJ
RVT=RVO=RV(I,J)
ROT=RO(I,J)
ET=E(I,J)
TH1=TH=(ROL-ROT+RVO)/ROL                               SOLADF 977
IF(TH.LT.THC.OR.TH.GT.THC1)GO TO 5000                  SOLADF 978
IF(TH.GT.0.5)TH1=1.0-TH                                SOLADF 979
VAVE=0.5*((U(I,J)+U(I-1,J))**2+
1(V(I,J)+V(I,J-1))**2)**0.5                           SOLADF 980
VDAVE=0.5*((UD(I,J)+UD(I-1,J))**2+(VD(I,J)+VD(I,J-1))**2)**0.5
VTB=VDAVE+0.1*VAVE                                      SOLADF 981
4410 RB=(TH1/PNV)**0.333                                SOLADF 982
RBW=SWIN/(ROT*VTB**2+1.0E-10)                          SOLADF 983
RB=AMINI(RB,RBW)                                       SOLADF 984
AREA=3.0*TH1/RB                                         SOLADF 985
TVAP=SATT(P(I,J))
EVT=EVCAL(TVAP)*ETEM1+ECV*ETEM
TLQ=TC+(ROT*ET-RVO*EVT-(ROT-RVO)*ECL)/(RVC*CHV*ETEM+(ROT-RVO)*CHL)
TLQ=AMAX1(TLQ,256.0)                                    SOLADF 986
PSAT=SATP(TLQ)
EVSAT=EVCAL(TLQ)
RSAT=TH*PSAT/(GAM1*EVSAT)                             SOLADF 987
IVSL=0
RBETA=1.0E-5*ROL
IRVT=0
RTOT=(RSAT+RVO)/2.0
RET=ROT*(ET-ECL)
RCH=ROT*CHL
CEC=CHV*ETEM-CHL
C   START ITERATION
4412 CONTINUE
RVT=RV(I,J)
II=I
JJ=J
CALL EOS

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TVAP=SATT(PT)
EVTM=EVCAL(TVAP)*ETEM1+ECV*ETEM-ECL
TLQ=TC+(RET-RVT+EVTM)/(RCH+PVT*CEC)
ELTBL=(1.0+1.91*ROL*TH/RTOT*CHL/ELHT*ABS(TLQ-TVAP)
1+(0.637*VTB*R8*ROL/EDL)**0.5)/R8
RATE=DELT*(ELTBL*AREA/ELHT*EDL*CHL*(TLQ-TVAP))*100.0
RVEQ=RVT-RVO-RATE
IF(ABS(RVEQ/RTOT).LT.0.001)GO TO 5000
IRVT=IRVT+
IF(IPVT.GT.25)GO TO 5000
IF(IRV1.EQ.1)SRH=SIGN(1.0,RVEQ)
IF(IVSL.EQ.1)GO TO 4412
IF(RVEQ.GE.1.0)GO TO 4415
C      SEEK BOUNDS
FMN=RVEQ
RVMN=RV(1,J)
IF(SRH.GE.0.0)IVSL=1
RV(1,J)=AMIN1(RVT+RBETA,ROT)
IF(IVSL.EQ.1)RV(1,J)=0.5*(RVMN+RVMX)
RBETA=2.0*RBETA
GO TO 4412
4415 FMX=RVEQ
RVMX=RV(1,J)
IF(SRH.LT.0.0)IVSL=1
RV(1,J)=AMAX1(RVT-RBETA,0.0)
IF(IVSL.EQ.1)RV(1,J)=0.5*(RVMN+RVMX)
RBTCA=2.0*RBETA
TO 4412
C      CONTINUE
C      INVERSE BOUNDS
IF(RVEQ.LT.0.0)GO TO 4422
FVTP=RVT-RVEQ*(RVMX-RVT)/(FMX-RVEQ)
FMD=RVEQ
RVMX=RVT
RV(1,J)=RVTP
IF(RVTP.LT.RVMN)RV(1,J)=0.5*(RVMN+RVMX)
GO TO 4412
4422 RVTP=RVT-RVEQ*(RVT-RVMN)/(RVEQ-FMN)
FMN=RVEQ
RVMN=RVT
RV(1,J)=RVTP
IF(RVTP.GT.RVMX)RV(1,J)=0.5*(RVMN+RVMX)
GO TO 4412
5000 CONTINUE
RETURN
END

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SOL ADF	1015
SOL ADF	1016
SOL ADF	1017
SOL ADF	1018
SOL ADF	1019
SOL ADF	1020
SOL ADF	1021
SOL ADF	1022
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SOL ADF	1060

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SUBROUTINE PITER
C      PRESSURE ITERATION ROUTINE
*CALL COMDK
EVCAL(X)=ECV+CHV*(X-TC)
ELCAL(X)=ECL+CHL*(X-TC)
SATT(X)=255.2+117.8*X**0.223

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SOL ADF	1061
SOL ADF	1062
SOL ADF	1063
SOL ADF	1064
SOL ADF	1065
SOL ADF	1066

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532 288

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SATP(X)=(X-255.2)/117.8)**4.48
FLG=0.0
DO 300 J=JPB,JPT
DO 300 I=IPL,IPR
IF(A(I,J).EQ.0.0)GO TO 300
ABR=2.0*A(I,J)*A(I+1,J)/(A(I,J)*A(I+1,J))
ABL=2.0*A(I,J)*A(I-1,J)/(A(I,J)*A(I-1,J))
ABT=2.0*A(I,J)*A(I,J+1)/(A(I,J)*A(I,J+1))
ABR=2.0*A(I,J)*A(I,J-1)/(A(I,J)*A(I,J-1))
PMAX=AMAX1(PMAX,P(I,J))
SM=-1.0E+10
ICT=0
XMX=1.0E+10
XMN=0.0
PBB=0.0
260 PB=P(I,J)
D=(RDX*(ABR*U(I,J)-ABL*U(I-1,J))+ROY*(ABT*V(I,J)+IABB*V(I,J-1)))*DELT/A(I,J)
DMAX=EN(I,J)*RO(I,J)/PB
D=AMIN1(D,DMAX)
D=AMAX1(D,-0.99)
ROT=RON(I,J)/(1.0+D)
RVT=RV(I,J)/(1.0+D)
ET=EN(I,J)-P(I,J)*D/RO(I,J)
II=I
JJ=J
CALL EOS
S=PB-PT
IF((ICT.NE.0)BETA(I,J)=(PB-PBB)/(S-SM)
P(I,J)=PB-BETA(I,J)*S
IF(S.GE.0.0)GO TO 262
XMN=PB
IF(P(I,J).GE.XMX)P(I,J)=0.5*(XMN+XMX)
GO TO 266
262 XMX=PB
IF(P(I,J).LE.XMN)P(I,J)=0.5*(XMN+XMX)
266 CONTINUE
DELP=P(I,J)-PB
IF(ABS(DELP).LE.EPS1*PMAX)ICT=.00
IF(ABR.EQ.0.0)GO TO 270
U(I,J)=U(I,J)+2.0*DELT*RDX*DELP/(RO(I,J)+RO(I+1,J))
270 IF(ABL.EQ.0.0)GO TO 272
U(I-1,J)=U(I-1,J)-2.0*DELT*RDX*DELP/(RO(I-1,J)+RO(I,J))
272 IF(ABB.EQ.0.0)GO TO 274
V(I,J-1)=V(I,J-1)-2.0*DELT*ROY*DELP/(RO(I,J-1)+RO(I,J))
274 IF(ABT.EQ.0.0)GO TO 276
V(I,J)=V(I,J)+2.0*DELT*ROY*DELP/(RO(I,J)+RO(I,J+1))
276 CONTINUE
SM=S
PBB=PB
ICT=ICT+1
IF((ICT.GT.10)GO TO 300
FLG=1.0
GO TO 260
300 CONTINUE
RETURN
END

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SOL ADF 1067
 SOL ADF 1068
 SOL ADF 1069
 SOL ADF 1070
 SOL ADF 1071
 SOL ADF 1072
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 SOL ADF 1108
 SOL ADF 1109
 SOL ADF 1110
 SOL ADF 1111
 SOL ADF 1112
 SOL ADF 1113
 SOL ADF 1114
 SOL ADF 1115
 SOL ADF 1116
 SOL ADF 1117
 SOL ADF 1118
 SOL ADF 1119
 SOL ADF 1120
 SOL ADF 1121
 SOL ADF 1122
 SOL ADF 1123

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***** = = = = = / / / / / = = = = = *****

SUBROUTINE TSTEP          SOL ADF 1124
*CALL COMDK             SOL ADF 1125
C VARIABLE TIME STEP AND ITERATION PARAMETER ROUTINE   SOL ADF 1126
EVCAL(X)=ECV+CHV*(X-TC)   SOL ADF 1127
ELCAL(X)=ECL+CHL*(X-TC)   SOL ADF 1128
SATT(X)=255.2+117.8*X**0.223   SOL ADF 1129
SATP(X)=((X-255.2)/117.8)**4.48   SOL ADF 1130
DUMX=DVMX=1.0E-10   SOL ADF 1131
DO 615 I=2,IM1   SOL ADF 1132
DO 615 J=2,JM1   SOL ADF 1133
UDM=ABS(UD(I,J))   SOL ADF 1134
VDM=ABS(VD(I,J))   SOL ADF 1135
UMM=ABS(U(I,J))   SOL ADF 1136
VMM=ABS(V(I,J))   SOL ADF 1137
DUMX=AMAX1(DUMX,UDM,UMM)   SOL ADF 1138
615 DVMX=AMAX1(DVMX,VDM,VMM)   SOL ADF 1139
DUMX=4.0*DUMX   SOL ADF 1140
DVMX=4.0*DVMX   SOL ADF 1141
DTMY=1.01   SOL ADF 1142
IF(ITER.GT.5)DTMY=0.99   SOL ADF 1143
DELT0=DELT*DTMY   SOL ADF 1144
DELT=AM(NI)DELT0,DELX/DUMX,DELY/DVMX)   SOL ADF 1145
DELT=AM(NI)(DELT,DELMX)   SOL ADF 1146
C RELAXATION FACTOR   SOL ADF 1147
DO 620 I=PL,IPR   SOL ADF 1148
DO 620 J=JPB,JPT   SOL ADF 1149
IF(A(I,J).EQ.0.0)GO TO 620   SOL ADF 1150
ET=E(I,J)   SOL ADF 1151
ROT=RO(I,J)   SOL ADF 1152
RVT=RV(I,J)   SOL ADF 1153
II=I   SOL ADF 1154
JJ=J   SOL ADF 1155
CALL EOS   SOL ADF 1156
PT0=PT   SOL ADF 1157
DELP=-EPS1*PT0*THTM*RO(I,J)   SOL ADF 1158
UR=4.0*DELT*ROX*DELP/(RO(I,J)+RO(I+1,J))*A(I,J)*A(I+1,J)/
I(A(I,J)+A(I+1,J))   SOL ADF 1159
IF(A(I+1,J).EQ.0.0)UR=0.0   SOL ADF 1160
UL=-4.0*DELT*ROX*DELP/(RO(I-1,J)+RO(I,J))*A(I,J)*A(I-1,J)/
I(A(I,J)+A(I-1,J))   SOL ADF 1161
IF(A(I-1,J).EQ.0.0)UL=0.0   SOL ADF 1162
VT=4.0*DELT*ROY*DELP/(RO(I,J)+RO(I,J+1))*A(I,J)*A(I,J+1)/
I(A(I,J)+A(I,J+1))   SOL ADF 1163
IF(A(I,J+1).EQ.0.0)VT=0.0   SOL ADF 1164
VB=-4.0*DELT*ROY*DELP/(RO(I,J-1)+RO(I,J))*A(I,J)*A(I,J-1)/
I(A(I,J)+A(I,J-1))   SOL ADF 1165
IF(A(I,J-1).EQ.0.0)VB=0.0   SOL ADF 1166
DT=DELT*(ROX*(UR-UL)+ROY*(VT-VB))/A(I,J)   SOL ADF 1167
ROT=RO(I,J)/(1.0+DT)   SOL ADF 1168
RVT=RV(I,J)/(1.0+DT)   SOL ADF 1169
ET=E(I,J)-PT0/RO(I,J)*DT   SOL ADF 1170
II=I   SOL ADF 1171
JJ=J   SOL ADF 1172

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532 290

85

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