## NRC Research ind Tach. Accession no. Assistance Report

$\qquad$

Contract Program or Project Title: Thermal Hydraulic LMFBR Safety Experiments

Subject of this Document: "Heat Removal Characteristics of Volume-Heated EAling Pools
With Inclined Boundaries"

Type of Document: Informal Report

Author (s): G. A. Greene, O. C. Jones, Jr., C. E. Schwarz, and N. Abuaf

Date of Document: June 1979

Responsible NRC T dividual and NRC Office or Division:
M. Silberberg

Division of Reactor Safety Research USs. Nuclear Regulatory Commisefou Washington, D.C. 2055:

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Ntiftof, NY 11973
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Prepared for
U.S. Nuclear regulatory Commission Washington, D. C. 20555
Under Interagency Agreement EY-76-C-02-0016
NRC FIN No. A- 3024
INTERIM REPORT 494 035

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# HEAT rEMOVAL CHARACTER*STICS OF VOLUME HEATED BOILING POOLS WITH INCLINED BOUNDARIES 

## By

G. A. Greene, O. C. Jones, Jr., C. E. Schwarz and N. Abuaf

Brookhaven National Laboratory
Thermal Hydraulic Development Division
Department of Nuclear Energy
Upton, New York 11973

MAY 1979

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Prepared for the U.S. Nuclear Regulatory Commission
    Office of Nuclear Regulatory Research
            Contract No. EY-76-C-02-0016
                FIN NO. A-3024
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#### Abstract

The state-of-the-art of heat transfer from boiling liquids having internal heat generation is reviewed. Considerable scatter is found in the existing data. Attempts to correlate these data have relied on both natural and forced convection concepts. This report describes a new series of experiments wherein the data scacter appears to have been improved by a factor of four to six from previous experiments finen compared on the basis of standard deviation in correlation coefficients.

Local heat transfer data to both vertical and inclined surfaces (up to $30^{\circ}$ from vert\&cal) are reported having aaximum to minimum heat transfer ratios of up to $5: 1$. It is shown that with surface vapor fluxes up to twice the free bubble rise velocities given by Harmathy ${ }^{22}$ there are two distinct flow regimes: bubbly and churn-turbulent.

In bubbly flows, the pool is generally quiescent and surface temperature fluctuations negligible. Two heat transfer regimes were identified: laminar-where $\overline{\mathrm{Nu}}=1.54 \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}, \theta)^{0.25}$ for $\mathrm{Ra}^{*} \leq 1.865 \times 10^{11}$, and turbu-lent-where $\overline{\mathrm{Nu}}=0.0314 \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}, \theta)^{0.40}$ for $\mathrm{Ra}^{*}>1.865 \times 10^{11}$. Standard deviations in the correlation coefficients were 0.08 and 0,0016 respectively.

In churn-turbulent flows, the pool is generally chaotic and three dimensional. The surface temperatures showed large fluctuations up tc de maximum pool-to-wall difference indicating intermittent destruction and renewal of boundary layer. Heat transfer coefficients were more uniform, and the maximum was observed to be in the range $.25-.30 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}^{\circ} \mathrm{C}$.


The data reported herein are in general agreement with the data reported by Gabor, et al. ${ }^{11}$, but with significantly less scatter. On the other hand, the more recent data of Gustavson, et al. ${ }^{12}$ are lower than those reported herein by approximately a factor of two.

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## NOMENCLATURE


V Velocity
Vj Drift velocity
x Coordinate (along wall)
Y Defined in Eq. 8
$Z$ Defined in Eq. 9
$\alpha$ Void fraction
$\bar{\alpha} \quad$ Average void fraction
B$\varepsilon \quad$ Fractional uncertainty
$\mu \quad$ Dynamic viscosity
$\checkmark \quad$ Kinematic viscosicy
$\rho$ Density
$\sigma \quad$ Standard deviation
$\theta \quad$ Wall angle from vertical
Tv Volumetric vapor source
૬ Normalized coordinate
Subscripts
ave Average
B Boiling
BN Boron nitride
C Forced convective
eff Effective
f Film 494 ..... 045
g ..... Gas

NOMENCLATURE (Cont'd)

| Q | Liquid |
| :--- | :--- |
| L | Laminar |
| N | Natural convective |
| o | Initial |
| T | Turbulent |
| TC | Thermal convective |
| v | Vapor |
| w | Wall |
| $\infty$ | Infinity |
| * | Modified |

## 1. INTRODUCTION

The heat transfer characteristics of volume-heated boiling pools are of importance in the safety analysis of hypothetical core disruptive accidents (HCDA) in liquid metal fast breeder reactors (LMFBR). In general, these pools would be composed of molten fuel and steel and would generate heat as a result of fission product decay. The fluid dynamic characteristics, as well as the containability of such boiling systems, would depend intimately on the heat loads applied to the surrounding boundaries. In addition, the thermodynamic and hydrodynamic states of the boiling mixture might determine the initial or boundary conditions for separate but related phenomena, such as nuclear recriticality, structural integrity, flow and freezing of multiphase fluids, etc. Confidence in the conceptualization, as well as computation of such hypothetical events depends to a great deal upon the ability to predict the vapor generation rate, void fraction, and local boundary heat flux from such volume-boiling pools. It is the purpose of this report to present new experimental data for local boundary heat transfer coefficients and average void fraction in volume-boiling pools and compare these results to previous experimental data, as well as to existing empirical models.

## 2. HISTORICAL REVIEW

Numerous studies exist in the literature concerning heat transfer from liquid pools with an internal heat source. A brief review of this literature is indicated in Table $1 .{ }^{1-7,13}$ However, investigations into the heat

TABLE 1

REVIEW OF NATURAL CONVECTION HEAT TRANSFER

transfer and hydrodynamic behavior of volume-heated boiling pools have been few and none are known to exist prior to this decade. ${ }^{8-12,15,18}$

The earliest known attempt to consider the heat transfer from volumeheated boiling pools is the work performed at Argonne National Laboratory by Stein, et al. ${ }^{9}$ In this work, a solution of NaCl and water was boiled in an open container by joule heating. The average downward and horizontal heat fluxes were measured by thermocouples soldered in small dead-end holes in the plates making up the electrodes and base, and in the coolant system.

A model was presented which separated the boundary heat transfer into a natural convection and forced convection regime. The natural convection regime was shown to agree with the correlation below,

$$
\begin{equation*}
\overline{\mathrm{Nu}}_{\mathrm{N}}=.677[\mathrm{Pr} /(.952+\mathrm{Pr})]^{1 / 4} \mathrm{Ra}^{1 / 4} \tag{1}
\end{equation*}
$$

where $\overline{\mathrm{Nu}}_{\mathrm{N}}=\dot{Q}_{\mathrm{N}} \mathrm{L} /(\mathrm{k} \Delta \mathrm{T}), \mathrm{Ra}=\mathrm{PrGr}$, and $\mathrm{Gr}=\mathrm{g} \beta \Delta \mathrm{TL}{ }^{3} / \mathrm{v}^{2}$. The forced convection regime was shown to be correlated by the relation,

$$
\begin{equation*}
\overline{\mathrm{Nu}}_{\mathrm{C}}=.644 \mathrm{Pr}^{1 / 3} \mathrm{Re}^{1 / 2} \tag{2}
\end{equation*}
$$

where $\overline{N u}_{C}=\dot{Q}_{C} L /(k \Delta T)$ and $R e=V_{B} L / v$. Both relations were valid only for laminar flow conditions. For convenience, a thermal convection reference velocity, $V_{T C}$, was defined as

$$
\begin{equation*}
\mathrm{V}_{\mathrm{TC}}=(\mathrm{gB} \mathrm{\Delta TL})^{1 / 2}=\left[\mathrm{g}\left(\rho_{\mathrm{w}}-\rho_{\infty}\right) \mathrm{L} / \rho_{\mathrm{f}}\right]^{1 / 2} \tag{3}
\end{equation*}
$$

and an equivalent free stream velocity, $V_{B}$, was defined as below;

$$
\begin{equation*}
V_{B}=40 \dot{Q}_{B}^{0.72} \tag{4}
\end{equation*}
$$

it was reported that for $V_{B} / V_{T C} \leq 0.2$, forced correction heat transfer was negligible and Eq. 1 was applied. For $V_{B} / V_{T C}>3.0$, therme? convection was negligible and Eq. 2 applied. For values of $V_{B} / V_{T C}$ intermediate to these values, mixed convection existed. The results of this investigation indicate that downward heat fluxes were found to be significantly larger than predictions from conduction theory would indicate; in addition, at the higher boiling beat fluxes, horizontal heat transfer was found to be significanty larger than values calculated by thermal convection alone, and could be correlated empirically by the laminar forced convection model.

The next attempt to experimentally characterize boundary heat transfer from volume-boiling pools was the work of Gabor, et al. ${ }^{11}$ from Argonne National Laboratory. In their work, they used simulant solutions of $\mathrm{ZnSO}_{4}$ in water. Base plates of two lengths (191 and 381 mm ) and three electrode heights ( $64,114,230 \mathrm{~mm}$ ) were used. The volumetric boiling power was supplied by joule heating as in the previous work. The electrodes and base plate were used as the heat transfer surfaces; thermocouples were buried halfway into the copper plates for temperature measurements, seven into the base plate, and two in each of the electrodes. Boundary heat losses were measured by calculating the enthalpy increase of the water coolant flowing in coils of copper tubing brazed to the backs of the electrodes. In these tests, the heat transfer rate to the vertical electrode was measured in two segments; for the 114 mm pool depth, the electrode was split into separately cooled segments of 25 mm at the top and 89 mm at the bottom. For the 230 mm pool depth, the electrode was split into a 25 mm upper segment and a 205 mm
lower segment. The opposite electrode was unsegmented and of the same overall length.

The ratios of the boundary heat fluxes, $\dot{Q}_{\text {upper }} / \dot{Q}_{\text {lower }}$, were investigated as a function of the boiling heat flux, $\dot{Q}_{B}$. It was found that for low boiling heat flux $\left(\dot{Q}_{B}\right.$ less than $3.5 \mathrm{cal} / \mathrm{cm}^{2}$ s for the 230 mm pool; $\dot{Q}_{B}$ less than $6 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}$ for the 114 mm pool), the ratio $\dot{Q}_{\text {upper }} \dot{Q}_{\text {lower }}$ was in the range of 1.5 to 2 , in agreement with the prediction from thermal convection theory. For high boiling heat flux $\dot{Q}_{B}$ greater than $4.5 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}$ for the 230 mm pool; $\dot{Q}_{B}$ greater than $9 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}$ for the 114 mm pool), the heat flux ratio was more nearly equal unity and equal to the heat flux to the unsplit electrode. The data was correlated in terms of a Nusselt number and Reynolds number based on the superficial vapor velocity. The Prandtl number was assigned separate exponential weight of $0,1 / 3$, and $1 / 2$ powers. As a result, a new model was presented for horizontal heat flux based on bubble-induced laminar forced convection of the form

$$
\begin{equation*}
\overline{\mathrm{Nu}}=\mathrm{CR} \mathrm{Re}^{1 / 2} \tag{5}
\end{equation*}
$$

where the constant $C$ included the effect of the Prandtl number and the superficial vapor velocity was defined as in Eq. 4.

While the studies reported so far have contributed to the understanding of some hydrodynamic and heat transfer processes occurring in internally heated boiling pools, they do not provide a mechanistic model for predicting local boundary heat transfer or void fraction in such pools. Recognizing this shortcoming, Gustavson, et al. ${ }^{12}$ undertook an investigation into the local distribution of boundary heat transfer and void fraction in internally
heated boiling pools. In their work, they also considered a rectangular pool of $\mathrm{ZnSO}_{4}$ and water, joule heated by passing a-c current through the pool between two electrodes. Instead of using the electrode as the heat transfer surface, an instrumented test plate was installed, designed to allow measurement of local heat transfer to thermally isolated segments. Each segment was cooled by flowing water through. separate cooling channels, and each flow rate was separately controlled to insure an isothermal poolside surface temperature. The heat flux to each segment was measured by measuring the temperature rise and the flow rate of the coolant for each segvent. The surface cemperature of each segment was determined by extrapolacing the interior thermocouple reading at the segment centerline to the rest wall surface across 0.38 mm of aluminum and 0.76 mm Teflon st a 5 , which was cemented to the aluminum test wall surface for electrical insulation from the pool. A constant level weir was connected to an inlet at the pool bottom, which fed a steady flow of fluid to the pool to identically replace the losses due to vaporization. In this fashion, the net power for vaporization could be determined. The accuracy of the measured heat transfer coefficient in these tests was reported to be $\pm 40$ percent.

The authors proposed that boundary heat transfer from volume-heated boiling pools was a mixed convection-type heat transfer phenomenon in which the effects were superimposable. They proposed, for laminar flow, that the thermal convective component be modeled as

$$
\begin{equation*}
\mathrm{Nu}_{\mathrm{N}}(\mathrm{x})=0.42[\mathrm{Gr}(\mathrm{x}) \cdot \mathrm{Pr}]^{0.25} \tag{6}
\end{equation*}
$$

In Ref. 12, the coefficient in Eq. 6 appeared as 0.41 instead of 0.42 suggested by Sparrow, et al. ${ }^{23}$ for $\operatorname{Pr}=1.86$, average $\operatorname{Pr}$ for all the present experiments.
where $\mathrm{Gr}(\mathrm{x})$ was the local Grashof number based on the average pool film density difference, and $\mathrm{Nu}_{\mathrm{N}}(\mathrm{x})$ now represented the local natural heat transfer correlation where x was measured along the heat transfer surface downward from the free surface. The forced convective component was represented by

$$
\begin{equation*}
\mathrm{Nu}_{\mathrm{C}}(\mathrm{x})=0.332 \operatorname{Re}(\mathrm{x})^{1 / 2} \cdot \mathrm{Pr}^{1 / 3} \tag{7}
\end{equation*}
$$

where $\operatorname{Re}(x)$ was the local Reynolds number based on the superficial vapor velocity at the pool surface. The method of modeling the combined natural/ forced convection from a volune-boiling pool consisted of correlating the ratio

$$
\begin{equation*}
\mathrm{Y}=\frac{\mathrm{Nu}}{\mathrm{Nu}} \tag{8}
\end{equation*}
$$

to the group

$$
\begin{equation*}
\mathrm{z}=\frac{\mathrm{Nu}_{\mathrm{C}}}{\mathrm{Nu} . \mathrm{i}} \tag{9}
\end{equation*}
$$

where Nu was the effective Nusselt number, either local or average, for the combined heat transfer process. Following a general correlation procedure, ${ }^{14}$ it was suggested that the functional form of the correlation should be

$$
\begin{equation*}
\mathrm{Y}=\left[1+\mathrm{z}^{\mathrm{n}}\right]^{1 / \mathrm{n}} \tag{10}
\end{equation*}
$$

where $n$ was determined by a best fit evaluation of the data (see Fig. 1). Alternate forms of Eqs. 5 and 7 were proposed for the case of turbulent heat

[^0]

Figure 1 Proposed Correlation Scheme for Matching Forced and Free Convection Components of Boundary Heat Transfer Relation. (BNL Neg. No. 9-591-76).

# transfer. The correlation was tested against the measured average heat transfer data from their tests. The best agreement was obtained using the laminar relations and a value of $\mathrm{n}=1 .^{12}$ 

## 3. ANALYTICAL MODELING

### 3.1 Heat Transfer

The data of Gustavson, et al. represent the first data available for local convective heat transfer coefficient from volume-heated boiling pools. The pres' nt authors conceived frons the available data that the mode of heat transf $r$, is stead of resembling mixed convection in which the effects were approximately superimposed, more closely approximated an enhanced mode of natural convection boundary layer flow and heat transfer. The phenomenon of boundary layer flow and heat transfer is depict ad in Fig. 2. It is assumed that the vapor rising through the pool causes a net liquid drift upward, which encounters the free surface and is forced so return downward along the cold boundary. In this case, the net bouvancy effect is due to the liquid-to-two-phase density difference. The heat transfer distribution from the volumetrically boiling pool to the bour dary exibits behavior not unlike a single-phase natural convection boindary layer, enhanced by the flow of net liquid recirculation due to upward vapor drag through the central liquid and downward along the walls. With this point of view in mind, single-phase natural convection boundary layer theory coupled with the 3 , andy effect of the two-rhase flow in the bulk liquid was used to attempt to correlate the Nusselt number to a modified Rayleigh number.


Figure $2 \quad \begin{aligned} & \text { Schematic of Boundary Layer Flow and Heat Tru.sfer From Volume } \\ & \text { Boiling Pool. (BNL Neg. No. } 9-369-76) .\end{aligned}$

$$
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$$

Assuming that $\alpha_{V_{\infty}} \ll(1-\alpha) \rho_{\ell_{\infty}}$, it has been shown that a modified Grashof number based on the void fraction may be defined as, ${ }^{15}$

$$
\begin{equation*}
G r^{*}(x, \alpha)=g_{e f f}\left[\rho_{W}-(1-\alpha) \rho_{\ell \infty}\right] \rho_{f} x^{3} / \mu_{f}^{2} \tag{11}
\end{equation*}
$$

Furthermore, if the boundary is inclined from the vertical by angle $\theta$ in such a manner that the boundary layer remains attached to the wall, the angle of inclination may be used to define the effective gravitational component in the direction of flow and Eq. 11 becomes

$$
\begin{equation*}
\mathrm{Gr}^{*}(\mathrm{x}, \alpha, \theta)=\mathrm{g} \cos \theta\left[\rho_{\mathrm{w}}-(1-\alpha) \rho_{\ell \infty}\right] \rho_{\mathrm{f}} \mathrm{x}^{3} / \mu_{\mathrm{f}}^{2} \tag{12a}
\end{equation*}
$$

If $\alpha \rho_{\ell \infty} \gg \rho_{W}-\rho_{\ell \infty}$, this may be reduced to the simple form below

$$
\begin{equation*}
G r^{*}(x, \alpha, \theta)=\frac{\mathrm{g} \cos \theta \alpha x^{3}}{v_{f}^{2}} \tag{12~b}
\end{equation*}
$$

The experimental data of Gustavson, et al. were correlated on the bases of modified single-phase natural convection theory of the forms below; a: Using average pool void fraction,

$$
\begin{equation*}
\mathrm{Nu}(\mathrm{x}, \bar{\alpha})=\mathrm{K}_{1}\left[\mathrm{Gr}^{*}(\mathrm{x}, \bar{\alpha}) \cdot \operatorname{Pr}\right]^{0.25} \tag{13a}
\end{equation*}
$$

and
b: Using locally measured void fraction,

$$
\begin{equation*}
\operatorname{Nu}(x, \alpha)=K_{2}\left[G r^{*}(x, \alpha) \cdot \operatorname{Pr}\right]^{0.25} \tag{13b}
\end{equation*}
$$

in which the properties used were the measured properties for the zinc sulfate solution at the appropriate film temperature. The value $\mathrm{K}_{\mathrm{i}}$ was determined from a log-log-linear least-squares fit to the data, and the
forms of the correlations are

$$
\begin{equation*}
a-N u(x, \bar{\alpha})=0.78\left[\mathrm{Gr}^{*}(\mathrm{x}, \bar{\alpha}) \cdot \operatorname{Pr}\right]^{0.25} \tag{14a}
\end{equation*}
$$

with a standard deviation in the correlation coefficient of $\pm .35$, and

$$
\begin{equation*}
\mathrm{b}-\mathrm{Nu}(\mathrm{x}, \alpha)=0.76\left[\mathrm{Gr}^{*}(\mathrm{x}, \alpha) \cdot \mathrm{Pr}\right]^{0.25} \tag{14b}
\end{equation*}
$$

with a standard deviation in the correlation coefficient of $\pm .56$. The data was visually interpreted to be in the laminar regime and for the most part, fell in the range $\mathrm{Ra}(\mathrm{x})<10^{11}$. The experimental data, as well as the $\log -\log -1$ linear least squares fit to the data for Eq. 14 a are shown in Fig. 3. The scatter in the correlation is basically the experimental scatter, and no finer structure was observed. The form of the correlation was insensitive to whether the average or local void fraction was used probably because the effects of the $1 / 4$ root of the void fraction over the measured range of $\alpha$ was lost in the scatter of the data. The use of the average void fraction is attractive since then the computation of local heat flux will not require knowledge of the local void distribution which is more difficult to measure and compute.

In order to use either correlation method to predict the effects of boundary heat transfer from volume-heated boiling pools, knowledge of the local or pool-average void fraction, as seen in Eqs. 13a,b is required.

### 3.2 Void Dynamics

The void distribution may be calculated based on a one-dimensional


Figure 3 Modified Laminar Natural Convection Correlation of Boundary Heat Transfer from Volumetric-Boiling Pools (Based on Data from Ref. 12). (BNL Neg. No. 1-1295-79).
two-phase drift flux model. ${ }^{15}$ Consider a volume-heated boiling pool in which the volumetric vapor source may be written as

$$
\begin{equation*}
r_{v}=\frac{\ddot{Q}_{B}^{\prime \prime \prime}(1-\alpha)}{h_{f g}} \tag{15}
\end{equation*}
$$

where the term ( $1-\alpha$ ) signifies that the local heat generation occurs only In the liquid and $r_{v}$ is the vapor source $\left(\mathrm{gm} / \mathrm{cm}^{3} \mathrm{~s}\right)$. For most low power boiling pools in which the evaporated liquid is "made-up," the liquid volume flux will be negligible in comparison to the vapor volume flux. The steady state vapor mass conservation equation may be written as

$$
\begin{equation*}
\frac{d j_{g}}{d x}=\frac{\ddot{Q}_{B}^{\prime \prime}(1-\alpha)}{\rho_{v^{h}} f g} \tag{16}
\end{equation*}
$$

The relation between the superficial vapor velocity, ${ }_{g}$, and the drift velocity, $\mathrm{V}_{\mathrm{gj}}$, may be written as, ${ }^{16}$

$$
\begin{equation*}
\frac{\left\langle j_{g}\right\rangle}{\langle\alpha\rangle}=C_{0}\langle j\rangle+V_{g j} \tag{17}
\end{equation*}
$$

where the notation $<>$ indicates a cross-sectional area average quantity. If we assume $<j \geqslant<j_{g}>$ and the distribution parameter $C_{0}=1.2$, this reduces to

$$
\begin{equation*}
\left\langle j_{g}\right\rangle=\frac{\langle a\rangle V_{g j}}{1-C_{o}^{\langle\alpha\rangle}} \tag{18}
\end{equation*}
$$

Assuming that the drift velocity, $\mathrm{V}_{\mathrm{gj}}$, can be represented as

$$
\begin{equation*}
V_{g j}=U_{\infty}(1-\alpha)^{n} \tag{19}
\end{equation*}
$$

-14- $\quad 494 \quad 060$
where $\mathrm{n}=0$ for churn-turbulent flow, and $\mathrm{n}=2$ for bubbly flow and dropping the bracket notation, Eq. 16 becomes

$$
\begin{equation*}
\frac{d}{d \xi}\left[\frac{\alpha(1-\alpha)^{n}}{1-C_{0} \alpha}\right]=K(1-\alpha) \tag{20}
\end{equation*}
$$

subject to the initial condition

$$
\begin{equation*}
\alpha=0 \quad \text { at } \quad \xi=0 \tag{21}
\end{equation*}
$$

where $\xi=x / H_{0}$ and $K=\frac{\dot{Q}_{B}{ }^{\prime \prime \prime} H_{0}}{\rho_{V^{h} f}{ }^{U}{ }_{\infty}}=j_{g^{\infty}} / U_{\infty}$. Equation 20 has been numerically integrated by two algorithms, an Euler predictor-corrector method and a fourth order Runge-Kutta method, with good agreement. The average void fraction, $\bar{\alpha}$, is defined as

$$
\begin{equation*}
\ddot{\alpha}=\frac{\xi_{\max }-1}{\xi_{\max }} \tag{22}
\end{equation*}
$$

The results of the local void fraction calculation were compared to the local void fraction data of Gustavson, et al., ${ }^{12}$ for four selected experimental runs for a value of $\mathrm{K}=1.75$. Although agreement between calculation and experiment was poor on a local basis, the average void fraction for the four runs, $\bar{\alpha}_{\text {meas }}{ }^{2} .40$, agreed quite well with the calculated average void fraction, $\bar{\alpha}_{\text {call }}{ }^{2} .41$.

The data of Gustavson, et al. ${ }^{12}$ tend to support the concepts of boundary layer heat transfer and one-dimensional two-phase drift flux vapor distribution modeling for pools in the bubbly flow regime. However, the uncertainty in the measurements performed make it difficult to differentiate the degree of agreement with the various models proposed, as well as to
identify the various flow regime transition criteria for hydrodynamic and heat transfer behavior. In particular, the conditions for transition from bubbly flow to churn-tabulent flow are not clear, nor are the changes in the associated hydrodynamic and heat transfer behavior. As a result, it is difficult to extrapolate these results to other heat transfer systems of interest, in particular the behavior of internally-heated boiling pools of nuclear fuel in an HCDA, which may exist at power levels beyond the range of the previous work. For these reasons, the experiment described herein was undertaken.

## 4. EXPERIMENTAL

### 4.1 Pool Description

A schematic view of the overall pool construction is seen in Fig. 4. The pool was rectangular in cross-section, 18 cm wide $\times 33.5 \mathrm{~cm}$ long. The electrodes were recessed into lexan walls and polished to eliminate surface nucleation. The electrodes, as well as the walls and base, could be supplied with cooling water flow to eliminate preferential surface nucleation If necessary. Evaporative and boiling vapor losses were recovered through a make-up water flow port connected to a constant level weir adjusted to $H_{o}$. The net makeup water flow rate was measured and converted to gross vaporization power. The makeup flow was introduced from the pool bottom into a baffled space to preheat the water to $\mathrm{T}_{\text {sat }}$ and prevent inlet subcooling effects. No boiling occurred in this space. The entire pool was

$$
-16-494 \quad 062
$$



Figure 4 Schematic View of Inclined-Wall Volumetric Boiling Pool Apparatus. (BNL Neg. No. 1-1385-79).
constructed of lexan with the exception of the copper electrodes and boron nitride test plate.

The boiling and nonboiling depths of the pool were measured with a voltage probe connected to a precision traversing mechanism. The conductor was lowered by the traversing mechanism until continuity was achieved and the voltmeter indicated the pool voltage. The pool was powered to the operating power and the probe was once again lowered until the operating voltage was again indicated. In this fashion, visual observations of pool depth were eliminated and more objective measurement of $H_{o}$ and $H_{B}$ was possible. The uncertainty in this measurement technique was essentially the fluctuations in pool height while boiling.

### 4.2 Test Plate

The test wall was constructed of lexan and was machined in such ? fashion that the base of the test wall was continuously inclinable from the vertical position to any inclined position. A schematic of the test wall is shown in Fig. 5. The test surface was composed of boron nitride sheet (I. 27 $\mathrm{cm} \times 30.5 \mathrm{~cm} \times 12.7 \mathrm{~cm})$, machined and recessed into the lexan wall with the pool-side surface flush with the lexan and in direct contact with the boiling pool. The material has heat transfer characteristics of an excellent thermal conductor, but is electrically insulating at the same time. These properties, along with low water absorption and thermal expansion and ease of machining, made BN an ideal material for these tests. In addition, no electrically-insulating covering was necessary, eliminating contact heat transfer resistance and temperature extrapolation. The rack


Figure 5 Schematic View of Inclined Wall Test Plate. (BNL Neg. No. 1-1384-79).
surface of the BN test plate was cooled by flowing water. A separate flow loop was designed to supply a continuous flew of water, $15-20 \mathrm{lpm}$, to remove the heat transfered to the wall. The flow rate was designed to be high enough that the convective resistance to heat transfer in the coolant loop was negligible. The entire back surface of the BN was exposed to the coolant flow. This eliminated channel coolant effects, as well as hot and cold spots from coil cooling techniques previously employed (see Fig. 5.). A picture of the assembled test pool may be seen in Fig. 6.

### 4.3 Test Plate Instrumentation

The BN was instrumented with chromel-alumel thermocouples for local heat transfer measurements. The thermocouples were 0.025 cm diameter stainless steel-clad microthermocouples, which were machined flat at the junction and electro-gold-plated with $\sim 0.003 \mathrm{~cm}$ of gold forming the hot junction across the isolated chrome and alumel leads. A schematic of the cross-sectionally polished and gold-plated microthermocouples is shown in Fig. 7. A photograph of the polished but unplated thermocouple tip may be seen in Fig. 8. The thermocouples were individually calibrated at the 'ce point and steam point taking local barometric pressure into account, and the average calibration data for each was compared to NBS type $K$ data. It was found that all the gold-plated thermocouples calibrated to within $\pm .07{ }^{\circ} \mathrm{C}$ from the steam to the ice point. The gold-plated microthermocouples were then cemented into 26 locations in the BN wall, 19 on the front at 1.27 cm intervals, and 7 on the back at 3.81 cm intervals at locations listed in Table 2. They were installed in such a manner that the measuring junction


Figure 6 Photograpt: of Assembled Test Pool. (BNL Neg. No. 3-1418-77).

$$
494 \quad 067
$$



Figure 7 Schematic View of Gold Plated Microthermocouple.
(BNL Neg. No, $2-376-77$ ).

$$
494 \quad 068
$$



Figure 8 End View of Unplated Thermocouple Cross Section.
(BNL Neg. No. 3-1422-77).

$$
494 \quad 069
$$

THERMOCOUPLE LOCATIONS IN TEST PLATE

was flush with the wall surface within an estimated $\pm .003 \mathrm{~cm}$ tolerance and cemented in place under a microscope. The gold-plated function thus comprised part of the test wall surface. Heat losses along the thermocouple sheath were negligible since the leads were immersed in the plate at least 50 diameters.

### 4.4 Data Acquisition

The thermocouples were connected to a $150^{\circ} \mathrm{F} \pm .2^{\circ} \mathrm{F}$ oven-type reference function along with a thermocouple in the bulk pool, and the data was then routed to the automated data acquisition system. The centralized data acquisition and analysis system was constr zed around an HP 9640 system, consisting of a 21 MX minicomputer with 112 kilowords of central memory, 7.5 megaword cartridge disk, and 9 track magnetic tape transport. Control of the system was accomplished by interactive software, which received transfer parameters from the experimenter and proceeded to scan the data channels upon command. The thermocouples were scanned by a 300 channel guarded crossbar scanner, which transferred data to an integrating digital voltmeter with microvolt resolution. Each thermocouple was sequentially sampled until the standard deviation of the output converged to a preset criterion or the maximum sample limit was exceeded. At this point, the scanner proceeded to the next thermocouple and repeated the same procedure until all 27 thermocouples had been integrated. The raw data was transferred to magnetic tape and preliminary engineering calculations were performed to convert the thermocouple output and system properties into local convective heat transfer coefficient and average pool void fraction.

A photograph of the entire inclined wall boiling pool test apparatus may be seen in Fig. 9.

All the measuring devices and their uncertainties are listed in Table 3.

## 5. EXPERIMENTAL RESULIS

### 5.1 Range of Experiments

The experiments described have been performed over a range of dimensionless vaporization power, $\mathrm{f}_{\mathrm{g} \infty} / \mathrm{U}_{\infty}$, up to 1.8 . Local heat fluxes along the inclined boundary were measured as indicated in Eq. 23.

$$
\begin{equation*}
h(x)=\frac{k_{\text {BN }}\left(T_{\text {front }}(x)-T_{\text {back }}(x)\right)}{a \cdot\left(T_{\text {pool }}-T_{\text {front }}(x)\right)} \tag{23}
\end{equation*}
$$

Accuracy of these measurements was estimated to be within $\pm 5$ percent. The tests reported herein do not have local void fraction measurements included, but rather have been correlated only on an overall average basis. The pool-average void fraction was measured as indicated in Eq. 24 with an estimated accuracy of $\pm 3$ percent.

$$
\begin{equation*}
\bar{\alpha}=\left(H_{B}-H_{0}\right) / H_{B} \tag{24}
\end{equation*}
$$

A complete error analysis is presented in Appendix A. Sample calculation of the heat transfer data is presented in Appendix B. For the boiling experiments presented here, the wall angles investigated were $90^{\circ}, 75^{\circ}$, and $60^{\circ}$ from horizontal with an accuracy of $\pm .5^{\circ}$. The flow regimes that were investigated are listed below and will be discussed in this order:


Figure 9 Photograph of Assembled Test Pool Facility.
(BNL Neg. No. CN2-798-78).

TABLE 3

## LIST OF MEASURING DEVICES USED AND THEIR UNCERTAINTY

| INSTRUMENT | UNCERTAINTY |
| :--- | :--- |
| Thermocouple, Gold-Plated, Type K | $\pm .07{ }^{\circ} \mathrm{C}$ |
| Digital Voltmeter, HP 3455A | $\pm 1 \mu \mathrm{v}$ |
| Reference Junction, REF-CEL 200 | $\pm .10{ }^{\circ} \mathrm{C}$ |
| Traversing Mechanism | $\pm .001 \mathrm{~m}$ |
| Make-Up Flow Meter System | $\pm .2 \mathrm{ml} / \mathrm{s}$ |
| Power Stats | $\pm 1$ percent |
| Cross-Bar Scanner, HP 2911A, B, |  |
| Hewlett-Packard Minicomputer, 21MX Series |  |
| Printer-Plotter, States 42 |  |

$494 \quad 074$

## 1. Nonboiling, single-phase

2. Boiling
a. Incipient boiling
b. Bubbly flow regime
c. Transition
d. Churn-turbulent flow regime

Tests were performed to determine if there were any measurable effects of the test wall coolant flow rate and makeup water temperature upon the boundary heat transfer distribution. The coolant flow rate was varied from 15-20 Rpm with no measurable effect upon the magnitude of the measured heat transfer coefficients. The make-up water temperature was observed to have no effect as long as it entered the pool from the baffled preheating space at or close to $\mathrm{T}_{\text {sat }}$.

### 5.2 Nonboiling Regime

Initial experiments were performed in nonboiling pools in order to perform operational checkout of the equipment and instrumentation. In addition, the nonboiling heat transfer to vertical and inclined boundaries was of interest in order to examine the nature of the boundary layer heat transfer. For these experiments, the total power applied to the pool was in the range of 1.0 to 2.5 kw . Boiling was not allowed to occur and heat transfer was single phase only. The profile of local boundary heat transfer behaved similar to single-phase laminar natural convection. The greatest magnitude of the local heat transfer coefficient was measured at or near the top of the test plate (i.e., the leading edge of boundary layer) and was in
the range $0.015-0.020 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}^{\circ} \mathrm{C}$. Correlation of this data indicated that the behavior agreed very well with established single phase laminar natural convection as expected and verified the ability of the equipment to measure local boundary heat flux from volume-heated pools accurately. Also the use of the effective gravitational component,

$$
\begin{equation*}
g_{\text {eff }}=g \cos \theta \tag{25}
\end{equation*}
$$

was verified for natural convection, and the effect of the internal heat source on the boundary layer thickness was found to be negligible as calculated by the correction method of Randall and Sesonske. ${ }^{17}$

### 5.3 Incipient Boiling Regime

As the power that was applied to the pool was increased, the regime changed as volumetric bubble nucleation in the bulk liquid began to appear. The onset of nucleation was determined solely by visual observation of the pool. This regime was called the incipient boiling regime. The behavior was characterized by bubble formation and rise with little measurable increase in average pool height. The pool average void fraction was in the approximate range 0.00 to 0.03 . The behavior of the boundary heat transfer was once again observed to resemble laminar natural convection as in the nonboiling tests. However, it was observed as in Fig. 10 that the maximum local heat transfer coefficient increased to approximately $0.08 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}$ ${ }^{\circ}$ C. This was an increase over the nonboiling case of approximately a factor of $4-5$. This indicated that although boundary layer-type heat transfer behavior was persisting, the superficial vapor velocity of the rising steam was causing a net recirculation of liquid which was rising through the


Figure 10 Profile of Local Boundary Heat Transfer Coefficient From Volume Boiling Pool - Run No. 9001 - Incipient Boiling. (BNL Neg. No. 4-844-79).
saturated pool. This liquid drift would encounter the free surface and turn towards the cold boundaries and flow downwards along the wall, enhancing the boundary layer heat transfer as evidenced in the magnitude of the convective coefficient. Boiling inception appeared to begin at a threshold value of $\mathrm{j}_{\mathrm{g}_{\infty} / \mathrm{U}_{\infty}}$ approximately equal to 0.2 . Below this value, the pool was volume-heated single phase and above this value, two-phase effects and volumetric boiling became evident. This value of dimensionless superficial vapor velocity indicated the magnitude of the evaporative power losses from the pool. Reducing the total vaporization power, $j_{g_{\infty}} / U_{\infty}$, by the evaporative losses, $\left(j_{g_{\infty}} / U_{\infty}\right)_{0}$, yields the net boiling power presented in dimensionless form below:

$$
\begin{equation*}
\left(j_{g_{\infty}} / U_{\infty}\right)^{*}=j_{g_{\infty}} / U_{\infty}-\left(j_{g_{\infty}} / U_{\infty}\right)_{0} \tag{26}
\end{equation*}
$$

It is recognized that this evaporative loss term, $\left(\mathcal{g}_{g_{\infty}} / U_{\infty}\right)_{0}$, will be system dependent and will diminish as the pool free surface area to volume ratio decreases. Analysis of the void distribution was performed on the basis of $\left(\mathrm{j}_{\mathrm{g}^{\infty}} / \mathrm{U}_{\infty}\right)^{*}$ as will be seen.

### 5.4 Bubbly Flow Regime

A further increase in power applied to the pool resulted in net production of vapor and a finite void fraction. This flow regime was characterized by a stable array of densely packed bubbles which formed initially in the upper region of the pool above an essentially quiescent single-phase region below. As boiling power was increased further, the thickness of the bubbly boiling region increased, penetrating downward through the nonboiling region. The bubbly flow regime is a liquid-continuous flow regime in which -32-


Figure 11 Profile of Local Boundary Heat Transfer Coefficient From Volume Boiling Pool - Run No. 9013 - Bubbly Flow. (BNL Neg. No. 3-1891-79).
the dispersed phase is the vapor. As verification of this assumption, the time trace of power vs time was examined to determine the effect of bubbly motion upon electrical coupling of the liquid. Any decoupling of the liquid from the applied electric field would be recognized by a transient fluctuation of the power trace. None was evidenced, indicating that the pool power was evenly applied and distributed through the continuous liquid phase creating a constant volumetric power density. The pool average void fraction was observed to be very sensitive to the vaporization power in the bubbly flow regime. A small increase in $j_{g_{\infty}} / U_{\infty}$ resulted in a rather large increase in $\bar{\alpha}$, as shown later in Fig. 14.

The maximum average void fraction achieved in these tests occurred for the bubbly flow regime just prior to flow regime transition and was approximately $0.55-0.60$. While in the bubbly flow regime, the pool was observed to swell periodically. This is believed to be caused by local subcooling effects due to reentry into the pool of cold liquid from recirculating boundary layer flow which would cool the pool temporarily below the satration temperature and induce partial void collapse.

The spatial profile of local boundary heat transfer coefficient maintaine its boundary layer nature as before. An example of this is shown in Fig. 11. The maximum local heat transfer coefficient was observed, in all cases, to be at or near the pool surface, and its magnitude was measured at approximately $0.20 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}{ }^{0} \mathrm{C}$. The coefficient was observed to decreise as depth along the heat transfer surface increased and it was observed to vary in magnitude along the test plate surface by approximately a factor of 3-5. Th average heat transfer coefficient for all the runs in the bubbly
flow regime was calculated to be approximately 0.10 ; this was greater than the single-phase tests by about a factor of 5 . The local heat transfer coefficient data was compared to the predictions of the models previously described in Eqs. $10^{12}$ and $14 a^{15}$, and the comparison is shown graphically, for Run 9013, in Fig. 11; the comparison on a local basis is available for all the runs in Appendix C and will be presented later on an average basis in Table 6. It is clear in Fig. 11 that the local convective heat transfer coefficients for these experiments exceeded the calculations of boch local heat transfer models available ${ }^{12,15}$ (derived from the local boundary heat transfer data from volume boiling pools reported in Ref. 12) by as much as a factor of 2 or more. This will be supported by Table 6 which will present a comparison of the average convective coefficients for all the experiments reported to the models referenced ${ }^{12,15}$ on an average basis.

### 5.5 Bubbly-Churn Turbulent Transition

As the pool power was increased further, a flow regime transition was observed to begin in the vicinity of $f_{g_{\infty}} / U_{\infty} \sim 0,8-1.0$. This flow regime was characterized by an increasing instability in bubble array order and the onset of bubble agglomeration; densely packed bubbles in a liquid continuous flow began to break down into large regions of liquid and large regions of vapor. The onset of this transition region appeared for the most part to coincide with full penetration of the boiling region to the pool bottom. The runs that characterize this region are runs 6009-6011 and 6014-6015 (hydrodynamic only). These runs demonstrated a partial pool collapse due to
the bubble agglomeration mechanism previously mentioned. Such behavior is observed in adiabatic bubble columns as previously shown. ${ }^{19,20,21}$ The previously observed good agreement between measured and calculated average void fraction based on $\left(\mathrm{f}_{\mathrm{g}_{\infty}} / \mathrm{U}_{\infty}\right)^{\text {* }}$ for the bubbly flow regime was no longer observed; instead the measured void fraction fell between the calculated values based on both the bubbly and the churn turbulent drift flux This will be seen in Table 6 where the calculated average void fraction is that based on the bubbly flow drift flux model, and the number in parentheses is that for the churn-turbulent drift flux model. In addition, the previously observed periodic pool swelling behavior diminished.

The spatial profile of local boundary heat transfer coefficient continued to maintain a strong boundary layer behavior as before. However, as seen in .jg. 12, a great deal of scatter appeared, and the variation along the test wall became less. As has been noted previously ${ }^{11}$, the magnitude of the heat flux became more nearly constant, and for this case, the ratio $\bar{h}_{\text {upper }} / \bar{h}_{\text {lower }}$ was approximately $4 / 3$. The average boundary heat transfer coefficient was measured for transition runs 6009-6011 only, and was found to be approximately 0.125 . This was greater by 25 percent over the average heat transfer coefficient in the bubbly flow regime and indicated that the hydrodynamic instability causing bubble asolomeration and flow regime transition was responsible for a corresponding increase in the boundary heat transfer coefficient. In spite of this apparent increase in the average boundary heat transfer, $t$ ie measured heat transfer coefficients were in the range of those for the highest power bubbly flow regime runs. Correlation of the transition region data was close to the bubbly flow data, as we will


Figure 12 Profile of Local Boundary Heat Transfer Coefficient From Volume Boiling Pool - Run No. 6009 - Transition. (BNL Neg. No. 4-845-79).
see in the next section, however, the scatter in the measurements was greater, indicative of the instability in void dynamics observed.

### 5.6 Churn-Turbulent Flow Regime

The churn-turbulent flow regime appeared to dominate for $\int_{g_{\infty}} / \mathrm{U}_{\infty} \geq 1.0$. This flow regime was characterized by a total breakdown in the well-ordered close packed bubble array observed for the bubbly flow regime. Instead, the hydrodynamic behavior appeared chaotic and highly "turbulent." Well-orderad flow patterns caused by upward vapor drift and downward boundary layer flow were no longer evident. In addition, the liquid-continuous flow hydrodynamics was destroyed by massive bubble agglomeration. This phenomenon appeared to be responsible for the creation of large regions locally which were entirely liquid or vapor. Large vapor flow paths appeared in the flow, allowing the escape of greater vapor mass flux than in the bubbly flow regime with considerably less liquid hold up. The result was a considerably lower average void fraction, as defined previously.

The flow regime transition from bubbi. o churn turbulent flow occurred suddenly and completely at a value of $\int_{g^{\infty}} / U_{\infty}$ approximately equal to one. At this point, the average void fraction suddenly collapsed from a value of 0.55-0.60 to approximately 0.4 . Simultaneously, the apparently reasonable assumption (born out at this time by visual observations only) of onedimensional flow for the bubbly flow regime and corresponding good agreement in average void fraction between experiments and one-dimensional drift flux void calculations appeared to no longer be valid for churn-turbulent flow. On the contrary, the flow appeared to become more three-dimensional in

Vehavior, and the applicability of one-dimensional drift flux modeling under these conditions is questionable. Nevertheless, comparison between experimentally measured average void fraction data and calculated values based on the one-dimensional drift flux model for churn-turbulent flow was good. The average void fraction measured for the cases of transition and churnturbulent flow was in the range of 0.40 and relatively insensitive to an increase in power for $j_{g_{\infty}} / U_{\infty}$ up to 2.0 .

During some of the bubbly flow runs, a thin but stable foaming layer was observed to form on the pool surface. The $\mathrm{ZnSO}_{4}$ electrolyte solution was frequently replaced to avoid the acidirion of unwanted contaminants, but no surface active chemicals were added to destroy this thin foam. The reasons for its formation are not well known, although its presence has been observed before. ${ }^{12,18}$ Regardless of its cause, the foam layer was invariably observed to completely and immediately disappear upon transition to the churn-turbulent flow regime, indicating that foams may not be an effective flow regime in such dynamic flow systems at dimensionless superficial vapor velocities in excess of unity, corresponding to vapor velocities greater than the bubble terminal rise velocity.

The heat transfer behavior also changed dramatically, and a sample of the local distribution of boundary heat transfer coefficient is demonstrated in Fig. 13. The apparent boundary layer nature of the heat transfer distribution seemed to disappear, replaced by a more uniform heat transfer coefficient along the boundary. The maximum local heat transfer coefficient was observed to be in the range $0.25-0.30 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}{ }^{\circ} \mathrm{C}$. The coefficient was observed to fluctuate temporally and spatially as well.


Figure 13 Profile of Local Boundary Heat Transfer Coefficient From Volume Boiling Pool - Run No. 7516 - Churn-Turbuient. (BNL Neg. Nr. 4-846-79).

The temporal fluctuations are evident in the standard deviation data of the local boundary temperature history. Whereas for the bubbly flow regime, the standard deviation of the discretely sampled instantaneous wall temperature distribution was found to be in the range $0.2-0.6{ }^{\circ} \mathrm{C}$, and for the transition flow regime, the standard deviation of the wall temperature was found to be in the range of $1.0^{\circ} \mathrm{C}$, a dramatic increase was observed for the churn-turbulent flow regime. The standard deviation of the local wall temperature averaging technique was found to be in the range $2.0-7.0^{\circ} \mathrm{C}$, an order of magnitude greater than previously observed for the well-ordered bubbly flow regime. Interpretation of this data concerning the standard deviation of the local wall temperature in churn-turbulent flow indicated that the standard deviation was nearly equal in most cases to the difference between the saturated pool temperature and the average wall temperature, i.e.,

$$
\begin{equation*}
\sigma_{w}=T_{p o o l}-\bar{T}_{w} \tag{27}
\end{equation*}
$$

This was interpreted to mean that, intermittently, free stream conditions were present at the boundary of the pool. This indicated that the wall boundary layer was periodically being destroyed by the highly chaotic three-dimenstonal hydrodynamic behavior of the churn-turbulent flow and subsequently being reestablished. This type of intermittent renewal of the boundary layer may account for the enhanced heat transfer observed.

Investigation of the time trace of power applied to the pool was used to evaluate the effective overall electrical coupling of the liquid to the applied electric field as before. The power trace was observed to experience high frequency fluctuations in contrast to the steady nature of the
bubbly flow regime. This was interpreted to mean that due to hydrodynamic fluctuations in the pool, the electrical rest chance was fluctuating and perhaps portions of the liquid were becoming electrically isolated from the electric field; under such conditions the pool could no longer be characterized by liquid-continuous concepts. It is not clear at what dimensionless superficial vapor velocity (power) joule heating becomes ineffective in supplying uniform power density per unit 1 iquid volume due to the observed electrical uncoupling mechanism in churn-turbulent flow. In the churnturbulent flow regime for the heat transfer runs presented, the average boundary heat transfer coefficient was measured to be approximately 0.15 $\mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}$ ${ }^{\circ} \mathrm{C}$. This represented an increase of 50 percent over $t$.e average heat transfer coefficient measured for the bubbly flow regime,

## 6. DATA ANALYSIS AND DISCUSSIONS

### 6.1 Comparison of Calculated and Measured Pool Void Fraction

For the experiments presented so far, the average void fraction, $\bar{\alpha}$, was measured and compared to the dimensionless superficial vapor velocity based upon total vaporization power, $j_{g_{\infty}} / U_{\infty}$. It was demonstrated that there existed a threshold velocity, $\left(j_{g_{\infty}} / \mathrm{U}_{\infty}\right)$, below which the pool would not boil. Subsequently, a net boiling superficial vapor velocity was defined, $\left(j_{g_{\infty}} / U_{\infty}\right)^{*}$, as seen in Eq. 26. It has been determined that the quantity $\left(f_{g_{\infty}} / U_{\infty}\right)_{0}$ was a system parameter and approximately equal to 0.2 for these tests.

A composite diagram of all the average void fraction data is plotted in -42-

Fig. 14 as a function of the dimensionless superficial vapor velocity for the incipient boiling, bubbly flow, transition and churn-turbulent flow regimes.

The incipient boiling bubbly flow data and the churn-turbulent data were compared to the predictions of the one-dimensional drift flux model based upon net vaporization power and the appropriate drift flux model. In addition, the transition data were compared to both bubbly and churnturbulent flow models. The comparisons are shown in graphical form in Fig. 15. The model was in fair agreement with the bubbly flow ard incipient boiling data for small values of $\mathrm{j}_{\mathrm{g}_{\infty}} / \mathrm{U}_{\infty}$ and improved considerably as this value increased. This behavior was not unexpected in view of the strong sensitivity of the void fraction in the bubbly flow regime to small changes in boiling power as demonstrated in Fig. 14.

The transition data, due to the nature of the onset of flow instability, demonstrated poor agreement with both the bubbly and churn-turbulent flow models. As expected, however, the measured values did all fall intermediate to the two model predictions.

For the churn-turbulent flow data, good overall agreement between experiment and analysis was achieved. As pointed out previously, the sensitivity of the void fraction in the churn-turbulent flow regime co the vaporization power was considerably less than in the bubbly flow regime. This means that fluctuations in the power are not strongly reflected in the measured or calculated void fraction. This behavior is evident in Fig. 14 where the churn-turbulent data demonstrated a flat profile.

These data indicated the following behavior:


Figure 14 Pool Average Void Fraction, $\bar{\alpha}$, vs Dimensionless Superficial Vapor Velocity Based on Total Vaporization Power. (BNL Neg. No. 4-993-79).

$$
\begin{array}{lll}
-44- & 494 & 090
\end{array}
$$



Figure 15 Comparison of Measured Average Void Fraction to Calculated Average Void Fraction Based on Net Boiling Power. (BNL Neg. No. 4-992-79).

1. Evaporative losses were substantial in the bubbly flow regime for this pool with a large surface to volume ratio. For pools with smaller surface to volume ratio, this dependence is expected to diminish. In the churn-turbulent regime, the void fraction was less sensitive to the power and uncertainty in the boiling power and boundary heat losses contributed smaller uncertainty in the measured and calculated void fraction as demonstrated.
2. The nonboiling portion of the pool in the bubbly flow regime was not taken into account by the drift flux model. In addition, the effect of wall angle on the superficial vapor velocity is not presently incorporated into the analysis presented.
3. For the experi -nts repnrted to date, the pool-average void fraction never exceeded 0.60 .
4. Transition from bubbly flow to churn-turbulent flow occurred for the total dimensionless superficial vapor velocity, $\mathrm{j}_{\mathrm{g} \infty} / \mathrm{U}_{\infty}$, in the range 0.8 to 1.0 . This transition was accompanied by an immedfate and sudden collapse in pool average void fraction from approximately $0.55-0.60$ to 0.40 .
5. The thin foaming layer which existed for some of the bubbly flow runs invariably was destroyed during transition to churn-turbul.ent flow. This was interpreted to mean that foaming flows were unstable for $\mathrm{f}_{\mathrm{g}_{\infty}} / \mathrm{U}_{\infty}$ greater than unity.

$$
-46-\quad 494 \quad 092
$$

### 6.2 Natural Convection Analysis of Previous Data ${ }^{12,11}$ For Heat Transfer From A Volume Boiling Pool to a Vertical Boundary

The correlation techniques described by Eqs. $13 \mathrm{a}, \mathrm{b}$ were applied to local convective boundary heat transfer data of Gustavson, et al. ${ }^{12}$ The assumption inherent in these equations is that the boundary layer is laminat, resulting in the assumed $1 / 4$ exponent on the Rayleigh number. The local heat transfer correlations derived from the local heat transfer data and the (a) averaged void fraction data, (b) local void fraction data were found to be

$$
\begin{equation*}
a-\operatorname{Nu}(x, \bar{\alpha})=0.78\left[\mathrm{Gr}^{*}(\mathrm{x}, \bar{\alpha}) \cdot \operatorname{Pr}\right]^{0.25} \tag{14a}
\end{equation*}
$$

and

$$
\begin{equation*}
b-\operatorname{Nu}(x, \alpha)=0.76\left[\operatorname{Gr}^{*}(x, \alpha) \cdot \operatorname{Pr}\right]^{0.25} \tag{14b}
\end{equation*}
$$

The standard deviations were found to be $\pm 0.35$ and $\pm 0.56$, respectively. It was observed that the convenience of utilizing the average void fraction, $\bar{\alpha}$, instead of the local void fraction, $\alpha$, resulted in little change in the correlation for heat transfer. The ratio of the correlation coefficients for the local heat transfer based on average vs. local void fraction was 1.03.

The same correlation technique was employed to test the existing data from 11 and 12 on an overall average basis. In this method, the average heat transfer coefficient and average void fraction were used. In conventional natural convection, it is assumed that the free stream density based on an equation of state is a constant since all properties outside the
boundary layer are evaluated at free stream temperature and pressure. This would be the case of using the average void fraction in the heat transfer correlation. For this case, it can be shown for laminar flow that direct Integration of the local heat transfer correlation yields the average heat tran ier correlation with $h$ replaced by $\bar{h}$ and $x$ replaced by $L$. The coefficient for the average correlation, $\mathrm{K}_{\mathrm{T}}^{\prime}$, is related to the coefficient for the local correlation from Eq. 13a as

$$
\begin{equation*}
z_{L}^{\prime}=1.33 K_{L} \text { for } \alpha(x)=\bar{\alpha} \tag{28}
\end{equation*}
$$

For the turbulent natural convection case, it can be shown similarly that the average turbulent correlation coefficient, $\mathrm{K}_{\mathrm{T}}^{\prime}$, is related to the local correlation coefficient, $K_{T}$, by

$$
\begin{equation*}
K_{L}^{\prime}=.83 K_{T} \text { for } \alpha(x)=\bar{\alpha} \tag{29}
\end{equation*}
$$

(See Appendix D for the derivation of Eqs. 28 and 29).
The average heat transfer data $r$ _ Fustavson, et al. ${ }^{12}$ were analyzed using the reported values for the superficial vapor velocity, average heat transfer coefficient, and average void fraction, as well as measured properties for $\mathrm{ZnSO}_{4}$ electrolytic solution.

The average heat transfer coefficients for the data of Gabor, et al. ${ }^{11}$ were not reported. They were calculated from the reported values of electrode heat flux, wall temperature, and pool temperature as

$$
\begin{equation*}
\overline{\mathrm{h}}=\frac{\dot{q}_{\text {electrode }}^{\prime}}{\left(\mathrm{T}_{\text {pool }}-\overline{\mathrm{T}}_{\mathrm{W}}\right)} \tag{30}
\end{equation*}
$$

The superficial vapor velocity, $j_{g^{\infty}}$, was calculated from the reported value for the boiling heat flux as

$$
\begin{equation*}
j_{g \infty}=\frac{q_{B O I L}^{\prime}}{\rho_{v} f_{f g}} \tag{31}
\end{equation*}
$$

The average void fraction data was reported. Only runs with an average void fraction greater than or equal to 0.05 were analyzed.

The results of the analysis of the data from Ref. 12 and Ref. 11 are presented in tabular form in Tables 4 and 5, and in graphical from in Figs. 16 and 17. It was found that the average natural convection correlation of the da: $:$ of Gustavson, et al. ${ }^{12}$ was

$$
\begin{equation*}
\overline{\mathrm{Nu}}=\frac{\overline{\mathrm{h} L}}{k}=1.07\left[\frac{g \bar{\alpha}{ }^{3} \mathrm{Pr}}{v_{\mathrm{f}}^{2}}\right]^{0.25} \tag{32}
\end{equation*}
$$

with a standard deviation of $\pm .30$. The exponent was assigned from inspection of the data. The standard deviation of the correlation coefficient was found to be 0.30 or 28 percent, indicative of the scatter in tic data. The majority of the data fell in the range of $\mathrm{Ra}^{*}<10^{12}$. As indicated in Fig . 16 , there was no noticeably different trend observed for the foam or dense data. Examination of the magnitude of the average heat transfer coefficient and superficial vapor velocity indicated that most of the data fell in the bubbly flow region previously identified. The ratio of the correlation coefficients, $K^{\prime} / K$, was 1.37 , tending to reinforce the use of the average void raction in correlating the local heat transfer data as well, as indicated in Fig. 3.


Figure 16 Natural Convection Correlation of Average Heai Transfer Data of Gustavson, et al. ${ }^{12}$ (BNL. Neg. No. 4-1304-79).

TABLE 4
AVERAGE HEAT TRANSFER AND VOID FRACTION DATA OF GUSTAVSON, ET AL. ${ }^{12}$ AND COMPARISQN TO EXISTING MODELS


[^1]It was found that the average natural convection correlation of the data of Gabor, et al. ${ }^{11}$ was

$$
\begin{equation*}
\overline{\mathrm{Nu}}=\frac{\overline{\mathrm{h}}}{\mathrm{k}}=1.58\left[\frac{\mathrm{~g} \overline{\alpha L}^{3} \mathrm{Pr}}{v_{\mathrm{f}}^{2}}\right]^{0.25} \tag{33}
\end{equation*}
$$

with $=$ standard $2 v i a t i o n$ of $\pm 0.33$. There were no local void fraction or heat flux measurements available for this work to perform a similar comparison of the local and average heat transfer correlation coefficients as performed for the data in Ref. 12. Once again, the exponent was assigned after examining the data. The mafor'ty of the data fell in the range $\mathrm{Ra}^{*}$ < $10^{12}$ as did Gustavson's data ${ }^{12}$. If the data were divided into two groups at $R a^{*}=10^{11}$, the correlation of the data on the basis of laminar and turbulent behavior would result in the set of correlations below,

$$
\begin{equation*}
a-\overline{N u}=(1.42) \mathrm{Ra}^{*} \cdot 25 \quad \mathrm{Ra}^{*}<10^{11} \tag{34a}
\end{equation*}
$$

with a standard deviation of $\pm 0.25$ for the laminar data and

$$
\begin{equation*}
\mathrm{b}-\overline{\mathrm{Nu}}=(.0309) \mathrm{Ra}^{*} .40 \quad \mathrm{Ra}^{*}>10^{11} \tag{34b}
\end{equation*}
$$

with a standard deviation of $\pm .0058$ for the turbulent data. Although scatter in the data makes the determination of laminar vs. turbulent boundary layer behavior tenuous, both sets of correlations for Gabor's data are plotted in Fig. 17. This will be discussed in greater depth in the next section. Examination of the magnitudes of the average heat transfer coefficients and dimensionless superficial vapor velocities indicated that the majority of this data was also expected to fall in the bubbly flow regime.


AVERAGE HEAT TRANSFER AND "OED FRACTTON DATA OF GABOR, ET AL. 11

|  | RUN | $\begin{aligned} & \text { WALL TEMP } \\ & \text { DEG } \\ & \text { C } \end{aligned}$ | VOID FRACTION EXP | AVER EXP |  | $\begin{aligned} & \text { TRANS } \\ & \text { DEG C } \\ & \mathrm{EO} 10 \\ & \mathrm{~N}=: 1.0 \end{aligned}$ | COEFF $\begin{aligned} & \mathrm{EO} 10 \\ & \mathrm{~N}=0.7 \end{aligned}$ | SVV* | DFPTH <br> CM | NUX | RAX |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1 | 59.7 |  |  |  |  |  |  |  |  |  |
|  | 2 | 59.7 60.3 55 | .41 .42 | .098 .097 | . 072 | .073 .072 | .093 .091 | 1.34 1.28 | 19.39 19.79 | 1174. 1186. | $.5034 \mathrm{E}+12$ $.5502 \mathrm{E}+12$ |
|  | 3 | 55.6 | . 31 | . 085 | . 069 | . 06.5 | . 682 | . 74 | 16.50 | 872.0 | $.5502 \mathrm{E}+12$ $.2326 \mathrm{E}+12$ |
|  | 4 | 55.9 | . 51 | . 088 | . 072 | . 061 | . 077 | . 67 | 23.30 | 1275. | . $1079 \mathrm{E}+13$ |
|  | 5 | 49.7 | . 16 | . 065 | . 061 | . 056 | . 076 | . 41 | 13.64 | 551.0 | . $6592 \mathrm{E}+11$ |
|  | 6 13 | 50.3 67.0 | . 17 | . 068 | . 062 | . 056 | . 678 | . 49 | 13.70 | 580.6 | . $7086 \mathrm{E}+11$ |
|  | 13 15 | 67.0 | .06 .29 | . 058 | . 051 | . 046 | . 059 | . 25 | 12. 16 | 426.7 | . $1887 \mathrm{E}+11$ |
|  | 15 | 69.1 72.1 | . 29 | . 083 | .066 | . 064 | . 081 | . 67 | 14.39 | 733.9 | . $1050 \mathrm{E}+12$ |
|  | 17 | 72.1 | . 211 | .089 .136 | . 067 | . 0662 | .079 .082 | . 56 | 14.46 | 791.7 | $.1143 \mathrm{E}+12$ |
|  | 18 | 74.7 | . 35 | . 116 | . .973 | . 6.078 | . 088 | . 66 | 17.80 | 1428. | . $3777 \mathrm{E}+12$ |
|  | 19 | 77.9 | . 42 | . 136 | . 074 | .077 | . 099 | 1.61 | 19.70 | 1653. | $3465 \mathrm{E}+12$ $6001 \mathrm{E}+12$ |
|  | 20 | 77.4 | . 58 | . 137 | .074 | .069 | . 089 | 1.35 | 27. 16 | 2230. |  |
|  | 21 | 76.0 | . 45 | . 133 | . 074 | . 972 | . 092 | 1.20 | 26.76 | 1699. | . $7422 \mathrm{E}+12$ |
|  | 23 | 70.2 | . 05 | . 061 | . 049 | . 644 | . 056 | . 22 | 12.99 | 452.5 | . $1560 \mathrm{E}+11$ |
|  | 24 | 76.9 | .27 | . 092 | . 076 | . 953 | . 081 | . 58 | 15.60 | 885.3 | . $1915 \mathrm{E}+12$ |
|  | 25 | 81.3 | . 45 | . 130 | . 075 | . 1067 | . 086 | . 85 | 29.70 | 1658. | . $7619 \mathrm{E}+12$ |
|  | 26 | 83.2 | . 53 | .156 | . 075 | . 076 | . 098 | 1.80 | 24.30 | 2334. | . $1465 \mathrm{E}+13$ |
|  | 27 | 82.8 | . 05 | . 964 | . 050 | . 042 | . 953 | . 15 | 12.00 | 473.0 | . $1661 \mathrm{E}+11$ |
|  | 28 | 83.1 | . 16 | . 089 | . 058 | . 052 | . 956 | . 31 | 12.70 | 625.7 | . $3945 \mathrm{E}+11$ |
|  | 29 | 88.9 | . 45 | . 137 | . 076 | . 061 | . 077 | . 46 | 20.70 | 1744. | . $7868 \mathrm{EE}+12$ |
|  | 39 | 88.1 | . 39 | . 148 | . 075 | . 072 | . 092 | 1.67 | 18.76 | 1702. | . $5030 \mathrm{E}+12$ |
|  | 37 | 53.0 | . 06 | . 679 | . 049 | . 050 | . 063 | . 38 | 12.16 | 527.2 | . $1750 \mathrm{E}+11$ |
| u | 38 | 52.3 | . 09 | . 063 | . 054 | . 051 | . 065 | . 36 | 12.50 | 496.3 | . $2882 \mathrm{E}+11$ |
| i | 39 | 38.3 | . 23 | . 08.1 | . 065 | . 064 | . 08 ! | . 69 | 14.80 | 772.1 | . $1264 \mathrm{E}+12$ |
|  | 49 | 57.1 | .23 | . 081 | . 066 | . 064 | . 981 | . 68 | 14.80 | 744.8 | . $1256 \mathrm{E}+12$ |
|  | 46 | 50.4 | . 09 | .077 | . 054 | . 061 | . 076 | . 74 | 12. 50 | 599.8 | .2851E+11 |
|  | 52 | 44.5 50.8 | . 05 | . 060 | . 047 | . 049 | . 061 | . 40 | 12.90 | 456.8 | . $1355 \mathrm{E}+11$ |
|  | 55 | 50.8 | . 10 | .076 | . 055 | . 062 | .078 | .77 | 12.79 | $501.4$ | $.3330 \mathrm{E}+11$ |
|  | 56 | 58.1 | . 56 | .163 | .072 | . 067 | . 084 | 1. 14 | 25.90 | 1657. | $.1647 \mathrm{E}+13$ |
|  | 56 58 | 48.4 57 | . 06 | . 072 | .649 | . 059 | .073 | . 75 | 12.10 | 543.4 | $.1705 \mathrm{E}+11$ |
|  | 58 92 | 57.7 58.7 | .37 .99 | 195 .123 | .071 .063 | .073 .059 | .093 .087 | 1. 30 | 18.19 | 1191. | $.3707 \mathrm{E}+12$ |
|  | 95 | 59.2 | . 17 | + 130 | . 073 | -073 | . 098 | . 50 | 7.04 | 534.6 | . $5244 \mathrm{E}+10$ |
|  | 96 | 61.2 | . 32 | . 146 | . 081 | . 083 | . 105 | . 86 | 9.40 | 621.4 | . $1322 \mathrm{E}+11$ |
| 5 | 97 | 60.4 | . 34 | . 146 | . 982 | . 083 | . In5 | . 87 | 9.76 | 878.8 | $4576 \mathrm{E}+11$ $.5326 \mathrm{E}+11$ |
| $\stackrel{\square}{\square}$ | 102 | 65.3 | . 11 | . 087 | . 049 | . 043 | . 0.55 | . 43 | 25.80 | 1390. | -.3325E+12 |
|  | 103 | 70.8 | . 20 | .121 | . 056 | .053 | . 068 | . 91 | 28.80 | 2154. | . $8652 \mathrm{~L}+12$ |
|  | 105 | 64.9 | . 99 | . 085 | . 047 | . 0.42 | . 054 | . 44 | 25.30 | 1332. | $.2560 \mathrm{~L}+12$ |
|  | 196 | 71.2 | . 25 | . 129 | . 058 | . 953 | . 067 | . 83 | 30.70 | 2447. | . $1313 \mathrm{E}+13$ |
|  | 109 | 63.8 | . 14 | . 106 | . 055 | . 049 | . 063 | . 45 | 19.90 | 1307. | . $1926 \mathrm{I} \mathrm{E}+12$ |
| $D$ | 112 | 64.2 | . 16 | .113 | . 057 | . 050 | . 053 | . 43 | 20.40 | 1429. | . $2377 \mathrm{~F}+12$ |
|  | 117 | 66.3 52.4 | . 23 | . 132 | . 061 | . 66.0 | . 076 | . 96 | 22.20 | 1814. | .4452E+12 |
|  | 118 | 53.8 | . 13 | . 694 | . 058 | - 064 | -986 | . 43 | 6.80 | 347.1 | . $34955 \mathrm{E}+10$ |
|  | 121 | 56.3 | . 21 | . 099 | . 975 | . 675 | . 094 | . 55 | 8.10 | 4.88 .4 | $.8710 \mathrm{E}+10$ $1871 \mathrm{E}+11$ |
|  | 122 | 53.6 | . 18 | . 084 | .073 | . 679 | . 089 | . 44 | 7.80 | 497.8 | . $1406 \mathrm{E}+11$ |
|  | 123 | 49.2 | . 05 | . 068 | . 054 | . 956 | . 076 | .27 | 6.70 | 284.1 | $.2423 \mathrm{E}+10$ |

SVV indicates the dimensionless superficial vapor velocity.

### 6.3 Correlation of Present Data

In a similar fashion to the natural convection correlation procedure employed to analyze the data from Refs. 11 and 12 , the average heat transfer data for the present tests were likewise analyzed. The local distribution of boundary heat transfer coefficient, $h(x)$, was integrated to determine the average heat transfer coefficient. The average void fraction was determined as described in Eq. 24. The superficial vapor velocity, $\mathcal{I}_{\mathrm{go}}$, was determined by converting the flow rate of make-up water into an average vapor flux as

$$
\begin{equation*}
j_{g_{\infty}}=\frac{G_{\rho_{\ell}} H_{o}}{V_{01} \cdot \rho_{v}} \tag{35}
\end{equation*}
$$

where $G$ is the make-up flow rate $\left(\mathrm{cm}^{3} / \mathrm{sec}\right), H_{0}$ is the nonboiling pool depth $(\mathrm{cm})$, and Vol. is the total pool volume $\left(\mathrm{cm}^{3}\right)$.

In a straight-forward fashion, the average heat transfer correlation was determined as indicated below for laminar bubbly flow,

$$
\begin{equation*}
\overline{\mathrm{Nu}}=(1.54) \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}, \theta)^{0.25} \mathrm{Ra}^{*} \leq 1.865 \times 10^{11} \tag{36a}
\end{equation*}
$$

with a standard deviation of $\pm 0.08$, and for turbulent bubbly flow

$$
\begin{equation*}
\overline{\mathrm{Nu}}=(0.0314) \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}, \theta)^{0.40} \quad \mathrm{Ra}^{*}>1.865 \times 10^{11} \tag{36b}
\end{equation*}
$$

with a standard deviation of $\pm 0.0016$. These data are available in Table 6 , and the two correlations are plotted in Fig. 18. The scatter in the data is seen to be less by a factor of $4-6$ than previously observed as can be seen in the standard deviation in the correlation coefficients and the transition from laminar to turbulent behavior is more evident. It was on


TABLE 6
COMPARISON OF MEASURED AND CALCUTATED AVERAGE VOID FRACTION AdD HEAT TRANSFER COEFFICIENT FROM VOLUMETRIC BOILINC POOL


[^2]** No heat transfer data were recorded for these runs.
the basis of these observations that the laminar-turbulent correlation of the data in Fig. 17 was analyzed. Although there have been no boundary layer measurements to substantiate the claim of turbulent boundary layer transition, the assigm ent of the laminar exponent (i.e., 0.25) and the turbulent exponent (i.e., 0.40 ) to the data correlation, similar to a singlephase natural convection, was done on the justification of observation of the marked change in the behavior of the data in the vicinity of $\mathrm{Ra}^{*} \sim 1-2 \mathrm{x}$ 16 ${ }^{11}$. This observation was made possible due to the elimination of the majority of the scatter in the data which was present in previous work. This is demonstrated by the relative scatter in the correlation coefficient which is 5 percent for both the laminar and curbulent cases. This is in sharp contrast to the 28 percent and 21 percent standard deviation in the correlations of the data of Gustavson ${ }^{12}$ and Gabor ${ }^{11}$ presented here, respectively.

The data in the transition region between bubbly and churn-turbulent flow exhibited more scatter than the bubbly flow data. However, correlation of this data behaved similar to the turbulent bubbly flow data as indicated in Fig. 18.

The churn-turbulent regime data, however, deviazed sharply from the above observed behavior. For the same Rayleigh number, the churn-turbulent data was observed to lie significantly above the correlation for bubbly flow. There is insufficient data at this point to make any quantitative statements to correlate the data to particular model assumptions. However, the magnitude of the temporal fluctuations in the wall temperature, as well as the significantly higher boundary heat transfer coefficient were interpreted to indicate that the multi-dimensional hydrodynamic nature of the
boiling pool was interfering with the formation of the wall boundary layer, if not destroying it.

The correlations derived from the data of Gustavson, et al. ${ }^{12}$ and Gabor, et al. ${ }^{11}$, as well as the present data, are summarized in Table 7. The local correlations of the data of Gustavson, et al. ${ }^{12}$ indicated little sensitivity to the use of either the average or local void fraction. The ratio of the local and average correlation coefficients supported the use of the average void fraction in the correlation of the local heat transfer data. The correlations of the data of Gabor, et al. ${ }^{11}$ and of the present work have been performed on an average basis only. Examination of the correlations derived revealed that the data agreed within the standard deviation of Gabor's data. ${ }^{11}$ However, both exceeded the correlation of Gustavson's data ${ }^{12}$ by a factor of approximately 1.5 . This is in agreement with observations that the local heat transfer data exceeded the calculations of the previous existing models ${ }^{12,15}$ derived from the data of Ref. 12 by a wide margin. The local heat transfer data of this work will be analyzed on a local basis in the future.

## 7. SUMMARY AND CONCLUSIONS

### 7.1 Bubbly Flow Regime

For volume-heated boiling pools characteristic of the kind investigated here and in the bubbly flow regime, the following conclusions can be made: 7.1.1 Hydrodynamics
(1) The bubbly flow regime persisted for a value of $j_{g_{\infty}} / U_{\infty}$

TABLE 7

SUMMARY OF LOCAL AND AVERAGE CORRELATIONS FOR HEAT TRANSFER FROM VOLUME BOILING POOLS

|  | AUTHOR | WALL hingle | LOCAL OR AVERAGE heat transfer | LAMINAR OR TURBULZNT | CORRELATION | $\begin{aligned} & \text { STANDARD } \\ & \text { DEV 'ATION } \end{aligned}$ | RANGE OF RAYLEIGH NUMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Gustavson, et al. ${ }^{12}$ | Vertical | Local <br> Local <br> Average | Laminar <br> Laminar <br> Laminar | $\begin{aligned} & \mathrm{Nu}(\mathrm{x})=.78 \mathrm{Ra}^{*}(\mathrm{x}, \bar{\alpha})^{0.25} \\ & \mathrm{Nu}(\mathrm{x})=.76 \mathrm{Ra}^{*}(\mathrm{x}, \alpha)^{0.25} \\ & \overline{\mathrm{Nu}}=1.07 \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}) 0.25 \end{aligned}$ | $\begin{aligned} & \pm .35 \\ & \pm .56 \\ & \pm .30 \end{aligned}$ | $\begin{aligned} & \mathrm{Ra}^{\star}<10^{12} \\ & \mathrm{Ra}^{*}<10^{12} \\ & \mathrm{Ra}^{\star}<2 \times 10^{12} \end{aligned}$ |
|  | Gabor, et al. ${ }^{11}$ | Vertical | Average <br> Average <br> Average | Laminar <br> Laminar <br> Turbulent | $\begin{aligned} & \overline{\mathrm{Nu}}=1.58 \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha})^{0.25} \\ & \overline{\mathrm{Nu}}=1.42 \mathrm{Ra}^{*}(\mathrm{~L}, \bar{\alpha}) \\ & \overline{\mathrm{Nu}}=.0309 \mathrm{Ra}^{\star}(\mathrm{L}, \bar{\alpha})^{0} \end{aligned}$ | $\begin{aligned} & \pm .33 \\ & \pm .25 \\ & \pm .0058 \end{aligned}$ | $\begin{aligned} & \mathrm{Ra}^{*}<2 \times 10^{12} \\ & \mathrm{Ra}^{*}<10^{11} \\ & \mathrm{Ra}^{\star}>10^{11} \end{aligned}$ |
|  | Present Work | $\begin{gathered} 90^{\circ}, 75^{\circ} \\ 60^{\circ} \end{gathered}$ | Average <br> Average | Laminar <br> Turbulent | $\begin{aligned} & \overline{\mathrm{Nu}}=1.54 \mathrm{Ra}^{\star}(\mathrm{L}, \bar{\alpha}, \theta)^{0.25} \\ & \overline{\mathrm{Nu}}=.0314 \mathrm{Ra}^{*}\left(\bar{L}_{\boldsymbol{\mu}}, \theta\right)^{0.40} \end{aligned}$ | $\begin{aligned} & \pm .08 \\ & \pm .0016 \end{aligned}$ | $\begin{aligned} & \mathrm{Ra}^{\star}<1.865 \times 10^{11} \\ & \mathrm{Ra}^{\star}>1.865 \times 10^{11} \end{aligned}$ |

$\rightarrow 0$
$\xrightarrow{\square}$
up to unity. In this flow regime, the pool underwent periodic swelling possibly due to subcooling from the returning cold boundary layer fluid into the pool bottom. The pool exhibited a stratified state with a boiling region over an essentially nonboiling single phase region below. The depth of the nonboiling region decreased as the volumetric vaporization source increased such that the nonboiling region was confined to the conditions where $\int_{g_{\infty}} / U_{\infty}<0.2$.
(2) The maximum average void fraction observed in the bubbly flow regime was in the range 0.55 to 0.60 at $\int_{\mathrm{g}^{\infty}} / \mathrm{U}_{\infty}$ approximately unity. In this range of power, transition to a churn-turbulent flow regime was observed in which boiling penetrated to the pool bottom, and a sudden collapse in average pool void fraction was observed from approximately 0.55-0.60 to 0.40 . While it might be coincidental, the upper limit in bubbly flow cor responds approximately with the packing density of spheres at $\mathrm{j}_{\mathrm{g}_{\infty}} / U_{\infty} \sim 1$.
(3) In the bubbly flow regime, it was observed that surface evaporative losses were non-negligible for pool geometry utilized herein having a large surface-to-volume ratio, and that a significant fraction of the volumetric power density went into these losses. The net boiling power was defined as the total vaporization power minus the evaporative losses. It was found
that calculation of the pool-average void fraction by means of a one-dimensional drift flux model based on the net boiling power agreed well with the experimental data Independent of the wall angle. Agreement between calculated and measured average void fraction improved for increasitig power.
(4) The average void fraction in the bubbly flow regime was found to be very sensitive to the volumetric boiling power. Small changes in $\mathrm{J}_{\mathrm{g}^{\infty}} / \mathrm{U}_{\infty}$ were observed to cause large variations in the poolaverage void fraction.

### 7.1.2 Heat Transfer

(1) Boundary heat transfer from volume-boiling pools in the bubbly flow regime behaved similar to natural convectiontype boundary layer heat transfer. The spatial variation in the local heat transfer coefficient was as great as a factor of 3-5 along the wall, with the greatest heat transfer at or near the pool surface. The data reported here for local convective heat transfer coefficient exceeded those previously reported by Gustavson, et al. ${ }^{12}$ by a factor of 2 or more but agreed with the earlier average pool data of Gabor, et al. ${ }^{11}$ within the scatter in their data.
(2) For boundary layer-type heat transfer from volume-boiling pools in the bubbly flow regime, the effect of small angle
of inclination of the boundary from vertical was modeled by defining an effective gravitational component along the wall as indicated below;

$$
\begin{equation*}
g_{e f f}=g \cos \theta \tag{37}
\end{equation*}
$$

where $\theta$ is the angle of inclination from the vertical. For the data described herein with inclinations up to $30^{\circ}$, this correlation proved adequate.
(3) Correlation of average heat transfer based on the average void fraction indicates laminar flow behavior up to Rayleigh number of $1.865 \times 10^{11}$. The correlation is of the approx mate form

$$
\begin{equation*}
\overline{\mathrm{Nu}}=1.54 \overrightarrow{\mathrm{Ra}}^{*}(\mathrm{~L}, \bar{a}, \theta)^{0.25} \text { for } \mathrm{Ra}^{*}<1.865 \times 10^{11} \tag{38a}
\end{equation*}
$$

For higher Rayleigh number, the data behaves similar to turbulent natural convection and the correlation for the range $\mathrm{Ra}>1.865 \times 10^{11}$ is

$$
\overline{\mathrm{Nu}}=0.0314 \widehat{\mathrm{Ra}}^{*}(\mathrm{~L}, \bar{\alpha}, \theta)^{0.40} \text { for } \mathrm{Ra}^{\star}>1.865 \times 10^{11}
$$

The consistency of the data presented herein represents a significant improvement over previously reported data of Gabor, et al. ${ }^{11}$ and Gustavson, et al. ${ }^{12}$. The standard deviation in both cases was found to be 5 percent in marked contrast to the previous data having standard deviations of 21 to 28 percent.

### 7.2 Churn-Turbulent Flow Regime

For volume-heated boiling pools characteristic of the kind investigated here and in the c.surn-turbulent flow regime, the following conclusions can he made:

### 7.2.1 Hydrodynamics

(1) Flow regime transition from bubbly flow to churnturbulent flow was observed to occur in the vicinity of $\mathrm{f}_{\mathrm{g}_{\infty}} / \mathrm{U}_{\infty}$ equal to one. Flow regime transition was accompanied by a marked collapse in the pool average void fraction from $0.55-0.60$ to 0.40 , similar to the bubble column observations in adiabatic flow as observed by Zuber and Hench and others. ${ }^{19,21}$
(2) The pool hydrodynamics appeared not to behave in a one-dimensional fashion any longer. Three-dimensional circulations appeared to dominate and caused largescale bubble agglomeration, responsible for the lower void fraction even at higher vapor generation rates than in bubbly flow. Periodic swelling behavior of the pool ended.
(3) The ifquid-continuous nature of the bubbly flow regime began to break down due to the three-dimensional nature of the flow and appearance of large vapor pockets in the flow. This behavior was indicated by fluctuations recorded in the power trace of the pool.
(4) Surface evaporative losses were found to be less significant for this regime than for bubbly flow. The void fraction appeared to be somewhat insensitive to increases in pool power in the range $1.0<\mathrm{j}_{\mathrm{g}_{\infty}} / \mathrm{U}_{\infty}<2.0$. The average void fraction in this range was measured $t$ be approximately 0.40 . The measured and calculated average void fraction data agreed well for the range of superficial velocity investigated.

### 7.2.2 Heat. Transfer

(1) The average heat transfer coefficient was approximately $0.15 \mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}^{\circ} \mathrm{C}$. Large fluctuations wert observed in the standard deviation of the incal wall temperature fluctuations. In some instances, the fluctuations were of the same magnitude as the difference between the pool temperature and the time-averaged wall temperature, indieating pavtial or complete local destruction and renewal of the wall boundary layer. It is this mechanism that is believed responsible for the increased boundary heat transfer coefficient.
(2) The profile of local heat transfer coefficient was more uniformly distributed along the wall, exhibiting large fluctuations spatially. The maximum local heat transfer coefficient was observed to be in the range $0.25-0.30$ $\mathrm{cal} / \mathrm{cm}^{2} \mathrm{~s}^{\circ} \mathrm{C}$.

## 8. ACKNOWLEDGEMENTS

The authors wish to express their appreciation to Dr. John D. Gabor of Argonne National Laboratory, Professor Muf£d S. Kazimi of Massachusetts Institute of Technology, and Professor John C. Chen of Lehigh University for their valuable discussions and insights into the work presented. The technical assistance of George A. Zinmer, James H. Klein, and Laurie Sweeney, and the skillful preparation of the manuscript by Nancy Schneider are acknowledged and greatly appreciated.

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## APPENDIX A

## Error Analysis:

1) Uncertainty in Heat Transfer Coefficient, $h(x)$ :

It was shown that the heat transfer coefficient is represented by the relation

$$
\begin{equation*}
h(x)=\frac{k_{\text {BN }}\left(T_{\text {front }}(x)-T_{\text {back }}(x)\right)}{a \cdot\left(T_{\text {pool }}-T_{\text {front }}(x)\right)} \tag{23}
\end{equation*}
$$

The total uncertainty in $h(x)$ may be computed by taking the total differentrial of Eq. 23 as follows:

$$
\begin{equation*}
d h=\frac{\partial h}{\partial k_{B N}} d k_{B N}+\frac{\partial h}{\partial \Delta T_{f b}} d \Delta T_{f b}+\frac{\partial h}{\partial a} d a+\frac{\partial h}{\partial \Delta T_{p f}} d \Delta T_{p f} \tag{A-1}
\end{equation*}
$$

where

$$
\begin{equation*}
\Delta T_{f b} \equiv T_{\text {front }}(x)-T_{\text {back }}(x) \tag{A-2}
\end{equation*}
$$

and

$$
\begin{equation*}
\Delta T_{p f} \equiv T_{\text {pool }}-T_{\text {front }}(x) \tag{A-3}
\end{equation*}
$$

The most probable mean square error, $\varepsilon_{h}^{2}$, is written as

$$
\begin{equation*}
\varepsilon_{h}^{2}=\left(\frac{d h}{h}\right)^{2} \tag{A-4}
\end{equation*}
$$

and this reduces to

$$
\begin{equation*}
\varepsilon_{h}=\left[\varepsilon_{k_{B N}}^{2}+\varepsilon_{\Delta T_{f b}}^{2}+\varepsilon_{a}^{2}+\varepsilon_{\Delta T}^{2}{ }_{p f}^{2}\right]^{1 / 2} \tag{A-5}
\end{equation*}
$$

The magnitude of each quantity in Eq. A-5 and its uncertainty are listed below:

| QUANTITY | MAGNITUDE | UNCERTAINTY |
| :---: | :---: | :---: |
| $\mathrm{k}_{\mathrm{BN}}$ | $0.041 \mathrm{cal/cm} \mathrm{~s}{ }^{\circ} \mathrm{C}$ | 0.001 |
| $\mathrm{~T}_{\mathrm{fb}}$ | $20^{\circ}{ }^{\circ} \mathrm{C} \quad *$ | 0.2 |
| $\mathrm{~T}_{\mathrm{pf}}$ | $20{ }^{\circ} \mathrm{C} \quad \star$ | 0.2 |
| a | 1.27 cm | 0.002 |
| Nominal Value |  |  |

The result is that the most probable uncertainty in $h(x)$ is approximately $\pm 3$ percent.
3) Uncertainty in Average Void Fraction, $\bar{\alpha}$ :

The average void fraction may be written as

$$
\begin{equation*}
\bar{\alpha}=\left(H_{B}-H_{0}\right) / H_{B} \tag{24}
\end{equation*}
$$

This may be written in differential form as shown 'efore to represent the mean square error in the average void fraction as

$$
\begin{equation*}
\varepsilon_{\bar{\alpha}}=\left[\left(\frac{H_{0}}{H_{B}\left(H_{B}-H_{0}\right)} d H_{B}\right)^{2}+\left(\frac{1}{H_{B}-H_{0}} d H_{0}\right)^{2}\right]^{1 / 2} \tag{A-6}
\end{equation*}
$$

where $\mathrm{dH}_{\mathrm{B}} \sim 0.5$ and $\mathrm{dH}{ }_{0} \sim 0.1 \mathrm{~cm}$. One can readily see that for small $\bar{\alpha}$ (ie., $\left(H_{B}-H_{0}\right)$ small), the fractional error will be large and approaching $\infty$ as $\left(H_{B}-H_{0}\right) \rightarrow 0$.

For $H_{0}=20 . \mathrm{cm}$ and $H_{B}=30 . \mathrm{cm}$, the result is the most probable error in $\bar{\alpha}$ was approximately 3 percent. For $H_{0}=25 . \mathrm{cm}$ and $H_{B}=30 . \mathrm{cm}$, the error was found to be approximately 8 percent. The value 3 percent is being used for the remaining calculations.
3) Uncertainty in Boundary Layer Coord tate, $x$

The coordinate from the free surface along the boundary layer, x , may be shown to be represented by

$$
\begin{equation*}
x_{i}=H B O I L / \cos \theta-H O N E+(i-1) E P S I \tag{A-7}
\end{equation*}
$$

in centimeters where HBOIL is the boiling depth measured from the pool bottom to the boiling free surface, HONE is the distance along the test wall from the base to the first thermocouple, and EPSI is the spacing along the test plate between thermocouples. The errors are independent and the incremental part, (i-1) (EPSI), does not accumulate as may be expected. The reason for this is that each thermocouple location was sited with respect to the same reference point and not with respect to the previous thermocouple along the plate. In this fashion, the positional uncertainty was not accumulative.

In a similar fashion to $\mathrm{A}-2$, the linearly independent uncertainties in $x$ may be shown to be

$$
\begin{equation*}
\varepsilon_{\mathrm{x}}=\left[\varepsilon_{\mathrm{HBOIL}}^{2}+\varepsilon_{\mathrm{HONE}}^{2}+\varepsilon_{\mathrm{EPSI}}^{2}\right]^{1 / 2} \tag{A-8}
\end{equation*}
$$

The situde of each quantity in Eq. $\mathrm{A}-7$ and its uncertainty are listed below:

| QUANTITY | MAGNITUDE $(\mathrm{cm})$ | UNCERTAINTY |
| :---: | :---: | :---: |
| BOIL | $30.0 \star$ | .5 |
| HONE | 31.4 | .2 |
| EPSI | 1.27 | .05 |

*Nominal Value

The result is that the most probable uncertainty in $x$ is approximately 5 percent.
4) Uncertainty in $\mathrm{Nu}(\mathrm{x})$ :

The Nusselt number is written as

$$
\begin{equation*}
N u(x)=\frac{h(x) x}{k_{f}} \tag{A-9}
\end{equation*}
$$

where the uncertainty in $k_{f}$ is negligible. Similar to $A-1$ and $A-2$, it may be shown that the most probable error in $\mathrm{Nu}(\mathrm{x})$ is

$$
\begin{equation*}
\varepsilon_{\mathrm{Nu}}=\left[\varepsilon_{\mathrm{h}}^{2}+\varepsilon_{\mathrm{x}}^{2}\right]^{1 / 2} \tag{A-10}
\end{equation*}
$$

From the previous sections, it is clear that che most probable error in $\mathrm{Nu}(\mathrm{x})$ is 6 percent.
5) Uncertainty in $\operatorname{Ra}^{*}(\mathrm{x})$ :

The Rayleigh number may be shown to be represented as

$$
\begin{equation*}
\mathrm{Ra}^{*}(\mathrm{x})=\frac{\operatorname{gax}^{3} \cos \theta \mathrm{Pr}}{v_{\mathrm{f}}^{2}} \tag{A-11}
\end{equation*}
$$

The uncertainty is $\bar{\alpha}$ has been shown previously to be approximately 3 percent. The uncertainties in $g, \cos \theta, \operatorname{Pr}$, and $v_{f}$ are negligible. This then reduces to

$$
\begin{equation*}
\varepsilon_{\mathrm{Ra}}{ }^{*}=\left[\frac{\varepsilon_{-}^{2}}{\alpha}+9 \varepsilon_{\mathrm{x}}^{2}\right]^{1 / 2} \tag{A-12}
\end{equation*}
$$

From previously c mpured results, the most probable error in $\mathrm{Ra}(\mathrm{x})$ is shown to be 15 percent.
6) Uncertainty in Laminar Correlation Coefficient, K:

For laminar bubbly flow, the correlation coefficient may be shown to be

$$
\begin{equation*}
\mathrm{K}=\mathrm{Nu} / \mathrm{Ra}^{* 1 / 4} \tag{A-13}
\end{equation*}
$$

It is readily shown that the most probable error in $K$ is

$$
\begin{equation*}
\varepsilon_{\mathrm{K}}=\left[\varepsilon_{\mathrm{Nu}}^{2}+\frac{1}{16} \varepsilon_{\mathrm{Ra}}{ }^{2}\right]^{\frac{2}{2}} \tag{A-14}
\end{equation*}
$$

The result is that for laminar bubbly flow, the most probable error in the correlation coefficient is expected to be 7 percent. Recall that the standard deviation in the correlation coefficient was computed to be 5 percent, further substantiating this result.
7) Uncertainty in Turbulent Correlation Coefficient, K:

For turbulent bubbly flow, the correlation coefficient may be shown to be

$$
\begin{equation*}
\mathrm{K}=\mathrm{Nu} / \mathrm{Ra}^{* 0.4} \tag{A-15}
\end{equation*}
$$

It may be readily shown that the most probable error in $K$ is

$$
\begin{equation*}
\varepsilon_{K}=\left[\varepsilon_{\mathrm{Nu}}^{2}+(.4)^{2} \varepsilon_{\mathrm{Ra}}{ }^{2}\right]^{1 / 2} \tag{A-16}
\end{equation*}
$$

The result is that for turbulent bubbly flow, the most probable error in the correlation coefficient is expected to be 8 percent. Recall that the standard deviation in the correlation coefficient was computed to be 5 percent, further substantiating this result.

All the quantities and their calculated probable errors are tabulated in Table A-1.

## TABLE A-1

SUMMARY OF EXPERIMENTAL UNCERTAINTIES

| QUANTITY | FRACTIONAL UNCERTAINTY |
| :--- | :--- |
| Heat Trans cer Coefficient | $3 \%$ |
| Boundary-Layer Coordinate | $5 \%$ |
| Average Void Fraction | $3-8 \%$ |
| Nusselt Number <br> Rayleigh Number <br> Laminar Correlation <br> Coefficient | $6 \%$ |
| Turbulent Correlation <br> Coefficient | $75 \%$ |

$494 \quad 121$

## APPENDIX B

## Sample Calculation (Run 9001)

For each experimental run, 27 thermocouple readings are shown in Appendix A corresponding the the 27 local heat transfer measurements. The locations of all the test plate thermocouples are listed in Table 2. Let $\mathrm{TC}_{i}, \mathrm{TF}_{i}, \mathrm{~TB}_{i}$ and $\mathrm{T}_{\text {POOL }}$ represent the actual thermocouple output, test plate front surface temperature, test plate coolant-side surface temperature, and pool temperature, respectively. Then the values $\mathrm{TC}_{i}(i=1-27)$ are mapped as follows:

$$
\begin{array}{ll}
\mathrm{TF}_{i}=\mathrm{TC}_{i} & i=1,19 \text { on } 1.27 \mathrm{~cm} \text { centers } \\
\mathrm{TB}_{i}=\mathrm{TC}_{i+19} & i=1,7 \text { on } 3.81 \mathrm{~cm} \text { centers } \\
\mathrm{T}_{\mathrm{POOL}}=\mathrm{TC}_{27} & \tag{B-3}
\end{array}
$$

The test plate coolant-side has only seven thermocouples; the temperature distribution is filled out to 19 values by linearly interpolating two values between each pair of back-side thermocouples as follows.

An array TBFUL is defined and equivalent to $T B$ by the assignment indicated below:

$$
\begin{equation*}
\mathrm{TBFUL}_{i}=\mathrm{TB}_{j} \tag{B-4}
\end{equation*}
$$

for $i=1,4,^{-}, 10,13,16,19 \quad j=7,6,5,4,3,2,1$. Next, $\operatorname{TBFUL}(2,3, \ldots 17,18)$ are in nearly in rpolated between the measured values TBFUL ( $1,4,7,10,13,16$, 19). This fills out the front and coolant side temperature distributions to 19 points each. For $i=1$, the data point is at the test plate top nearest to the free surface; for $i=19$, the data point is at the test plate bottom furthest from the free surface.

$$
-76-
$$

The following quantities are required for the actual calculation of the heat transfer data:

HBOIL is the boiling pool depth ( cm ) measured from the pool bottom.
$\theta$ is the wall angle inclination from vertical.
$k_{\text {wall }}$ is the boron nitride ( $B N$ ) test plate thermal conductivity $\left(=0.041 \mathrm{cal} / \mathrm{cm} \mathrm{s}^{\circ} \mathrm{C}\right)$.
a is the $B N$ thickness $(=1.27 \mathrm{~cm})$.
$\mathrm{T}_{\mathrm{f}}$ is the average film temperature for calculating boundary layer
properties $\left(=\left(\bar{T}_{\text {front }}+T_{\text {pool }}\right) / 2\right)$ where $\bar{T}_{\text {front }}=\frac{\sum \mathrm{TF}_{i}}{19}$.
$\bar{\alpha}$ is the average void fraction $\left(=\left(H_{B}-H_{0}\right) / H_{B}\right)$.
$\operatorname{Pr}$ is the Prandtl number evaluated at $T_{f}(=1.94$ for Run 9001).
$\mathrm{k}_{\mathrm{f}}$ is the film thermal conductivity evaluated at $\mathrm{T}_{\mathrm{f}}(=.00162 \mathrm{cal} / \mathrm{cm} \mathrm{s}$ ${ }^{\circ}$ C) for Run 9001 .
$v_{f}$ is the film kinematic viscosity evaluated at $T_{f}(=.3124$ cs for Run 9001).

$$
\mathrm{T}_{\text {film }}=91.7^{\circ} \mathrm{C} \text { for Run } 9001
$$

The local coordinate, $x_{i}$, local heat transfer coefficient, $h_{i}$, local Nusselt number, $N i_{1}$, and the local Rayleigh number, $R a_{1}$ are calculated according to the following formulae:

$$
\begin{equation*}
x_{i}=\mathrm{HBOIL} / \cos \theta-31.4+(i-1) 1.27 \tag{B-5}
\end{equation*}
$$

$$
\begin{equation*}
h_{i}=\frac{\left(T F_{i}-T B F U L_{i}\right) k_{\text {wall }}}{a \cdot\left(T_{\text {pool }}-T F_{i}\right)} \tag{B-6}
\end{equation*}
$$

$$
\begin{equation*}
N u_{i}=h_{i} x_{i} / k_{f} \tag{B-7}
\end{equation*}
$$

$$
-77-
$$

$$
\begin{equation*}
R a_{i}=g \bar{c} x_{i}^{3} \cos \theta \mathrm{Pr} / \nu_{f}^{2} \tag{B-8}
\end{equation*}
$$

for $i=1,19$ where $g$ is the gravitational acceleration coefficient. The procedure is identical for all the runs; $\bar{\alpha}$, HBOIL, and $T_{f}$ may be different for each run.

EXAMPLE: RUN 9001, $i=1$ (top most heat transfer channel, nearest pool surface).

$$
\begin{aligned}
& \text { HBOIL }=33.0 \mathrm{~cm} \\
& \cos \theta=1.0 \\
& \mathrm{x}_{1}=33.0-31.4=1.60 \mathrm{~cm} \\
& \mathrm{~h}_{1}=\frac{(89.8-59.5) .041}{1.27(101.1-89.8)}=.0866 \mathrm{cal} / \mathrm{cm} \mathrm{~s} \\
& \\
& \mathrm{Nu}_{1}=(.0866)(1.60) /(.00162)=85.4 \\
& \mathrm{Ra}_{1}=(980 .)(.03)(1.60)^{3}(1.94) /(.003124)^{2}=.2394 \times 10^{8}
\end{aligned}
$$

## APPENDIX C

In this appendix are listed the local heat transfer data for all the experiments performed. They are compiled in numerical order with the vertical wall $\left(90^{\circ}\right)$ data first, the $75^{\circ}$ data second, and the $60^{\circ}$ data last. The flow regime for each run is listed in Table 6 . The detailed thermocouple readings are indicated in the figures, and their locations on the test plate are tabulated in Table 2.

$$
494 \quad 125
$$








$494 \quad 132$

$494 \quad 133$
-87-



$494 \quad 134$


$494 \quad 136$


## RUN NJMBER. . . 9607

average void fraction. . . . . . . . . 12
INITIAL POOL. DEPTH(CH) . . . . . . . 29.0
BOIL.ING POOL. DEPTH(CW) . . . . . . . 33.8
VOL. POWER DENS. (CAL/CM3 SEC) . ,08
SUPERFICIAL VAPOR VEI.OCITY. .. . . 30
angle of Inclinat Ion (begirees) . .00
POOL VOL.UME (CM * * 3) . . . . . . . . . . . . 17435
PRANDTL NUMBER . . . . . . . . . . . . . . . 1.90
TOTAL PONER (KW) . . . . . . . . . . . . . . 9.24
average surface temi ideg Cl . . 85.9

| 6 |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| i | $\begin{aligned} & \text { DEPTH } \\ & \text { (CM) } \end{aligned}$ | LOCAL HEAT TRANSFER COEFF EXPT (CAL/CM2 SEC DEG C) |  |  |  | NUSSEI.T NUMBER | MODIFIED | RAYLEIG | NUMBER |
|  | 1.600 | . 1757 | . 0776 |  |  |  |  |  |  |
|  | 2.870 | . 1660 | . 0670 | . 0675 | .0960 0734 | 173.6 |  |  |  |
| 5 | 4. 140 5.410 | . 1599 | . 0612 | . 6490 | . 0734 | 293.6 |  | $.9742 \mathrm{E}+68$ $.5623 \mathrm{E}+69$ |  |
| - 0 | 6.680 | -1369 | . 0572 | . 0449 | .0591 | 467. |  | . $1688 \mathrm{E}+10$ |  |
|  | 7.959 | . 1143 | . 0543 | .0429 | . 0550 | 455.6 |  | 3766E+10 |  |
| + | 9.220 |  | . 0520 | . 0397 | . 0519 | 469.0 |  | .7090E+18 |  |
|  | 10.49 |  | . 0485 | . .0378 | . 0493 |  |  | 1195E+ 11 |  |
|  | 11.76 |  | . 0471 | . 0350 | . 0455 |  |  |  |  |
|  | 14.39 | . 0430 | .0-649 | . 0339 | . 0449 |  |  |  |  |
| 5 | 15.57 | . 9512 | . 0439 | . 0329 | . 0426 | 378.7 |  |  |  |
|  | 16.84 18.11 | . 0426 | .0431 | . 0313 | .0415 .0494 | 499.5 |  | 6955E+11 |  |
| $\infty$ | 19.38 | . 04446 | . 0423 | . 0396 | .0395 | 441.9 |  | $8977 \mathrm{E}+11$ $1136 \mathrm{E}+12$ |  |
|  | 29.65 | . 04.43 | . 0416 | . 02929 | . 0386 | 497.5 543.1 |  | $1413 \mathrm{E}+12$ |  |
|  | 21.92 | . 0441 | -0403 | . 02288 | . 0378 | 546.3 |  | 1731E-12 |  |
|  | 23.19 24.46 | .0451 | 0398 | . 0283 | . 03371 | 59.4 .6 |  | 2094E* 12 |  |
|  |  |  | . 0392 | . 6279 | .0358 | 642.8 |  | 2966E+12 |  |




$494 \quad 141$


$494 \quad 143$

## RUN NUMBER. . . 9016

| AVERAGE VOID FRACTION. | . 18 |
| :---: | :---: |
| INITIAL. POOL. DEPTHICM) | 27.0 |
| BOILING POOL. DEPTII (CM) | 33.0 |
| VOL. POWER DENS, (CAL /CM3 SEC). | . 13 |
| SUPERFICIAL VAPOR VLLOCITY.... | . 45 |
| ANGIE OF INCLINATION (DEGREES). | . 00 |
| POOL VOLUME CCM** | 16232. |
| PRANDTL. NUMBER | 1.87 |
| TOTAL POWER (KW) | 10.6 |
| AVERAGE SURFACE TEMP (DEG C) | 88.8 |


| $\rightarrow$ | $\begin{aligned} & \text { DEPTH } \\ & \text { (CN) } \end{aligned}$ | LOCAL IEAT TRANSFER COEFF (CAI. /CN2 SEC DEG C) |  |  |  | NUSSELT NIMBER | MODIFIED | RAYIEIGH | NUMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1.690 | .2108 | . 0862 |  |  |  |  |  |  |
|  | 2.879 | . 1963 | .0745 | . 06.37 | . 10848 | 267.4 346.4 |  | $.1479 \mathrm{E}+09$ |  |
|  | 4. 149 | . 1833 | . 0679 | .0563 | . 0746 | 346.4 466.6 |  | $.8535 \mathrm{E}+09$ |  |
|  | 6. 688 | +.1986 | . 06635 | . 6515 | . 0681 | 660.5 |  | +.5717E. 10 |  |
|  | 7.950 | . 1536 | . 0577 | . 0454 | . 06597 | 713.9 751.0 |  | . 1076E+ 11 |  |
|  | 9.220 |  | . 0556 | . 0432 | . 0567 | 751.0 |  | .1814E*11 |  |
|  | 11.76 |  | .0539 | 0415 $0-909$ | . 0543 |  |  |  |  |
|  | 13.63 | 0741 | . 0519 | . 0387 | -0505 |  |  |  |  |
|  | 15. 57 | .9775 | . 0498 | . 03775 | . 0489 | 651.1 |  | 1056E+12 |  |
|  | 16.84 | .0521 | . 0478 | . 0365 | 0.6476 0.64 | 741.5 |  | $1363 \mathrm{E}+12$ |  |
|  | 18.11 19.38 | . 0561 | . 0478 | . 0348 | . 0452 | 639.6 |  | $1724 \mathrm{E}+12$ |  |
|  | 20.65 | . 0505 | -0462 | . 0341 | . 0442 | 649.4 |  | 2144E*12 |  |
|  | 21.92 | . 0458 | . 0448 | . 0334 | . 0433 | 642.7 |  | . $3179 \mathrm{E}+12$ |  |
|  | 23. 19 | . 0426 | . 0442 | . 0322 | .0-417 | 617.7 |  | 3803E+12 |  |
|  | 24.46 |  | . 0436 | . 0317 | . 9410 | 6 . |  | 4503E +12 |  |






$494 \quad 149$
-103-

$494 \quad 151$
-105-
RUN mumber.

7561
AVERAGE VOIT FRACTION......... . . 11
INITIAL POOL. DEFTH(CM) . . . . . . . 29.9
BOILING POOL DLPTHiCM) ......... 32.5
VOL. PONER DFNS. (CAL/CM3 SEC). . 12
SUPERFICIAI. YAPOR VELDCITY.... . 46
ANGE OF INCL.INATION (DEGREES) - 15.00
POOL VOLUME $\left(C H^{*+3) ~ . ~ . ~ . ~ . ~ . ~ . ~ . ~ . ~ . ~ . ~ . ~} 19463\right.$.
PRANDTL. NUMBER. ................... . . . 1.91
TOTAI. POWER (KW) . . . . . . . . . . . . . . 11.0
average surface teup ideg C)... 84.7



$494 \quad 154$

$494 \quad 155$
-109-


[^3]

$494 \quad 159$-113-


$494 \quad 161$

$494 \quad 162$


494163 -117-


| $\begin{aligned} & \text { DEPTH } \\ & \text { (CX) } \end{aligned}$ | $\begin{aligned} & \text { LOCAL } \\ & \text { Expt } \end{aligned}$ | $\begin{aligned} & \text { IEAT TR TR } \\ & \text { CM2 } \\ & D_{0} 14= \end{aligned}$ | $\begin{gathered} \text { ANSFER } \\ E C \text { DEG } \\ E 010 \\ N=1.0 \end{gathered}$ | $\begin{aligned} & \text { COEFF } \\ & \text { C) } \\ & \text { E010 } \\ & \text { Noe.7.7 } \end{aligned}$ | nUSSELT NUMBER | MODIFIED | Rayleign | NUMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2.246 | . 1963 | . 0868 | . 6734 | . 0976 |  |  |  |  |
| 3.516 4.786 | - 17562 | . 07776 | -.6631 | . 0836 | 378.7 |  | $.5718 \mathrm{E} \cdot 69$ $.2193 \mathrm{E} \cdot 10$ |  |
| 6.056 | . 1576 | . 0677 | -0528 | . 0759 | 466.6 586.8 |  | . $5331 \mathrm{E} \cdot 10$ |  |
| 7.326 | -1372 | . 0646 | . 0496 | . 0658 | 653.8 |  | -984E* 11 |  |
| 8.596 | . 1357 | . 6621 | . 0471 | . 6616 | 717.6 |  | 3204E.1i |  |
| 9.866 |  | . 9609 | .0451 | . 0588 |  |  | -320-4.-1i |  |
| 11.14 |  | . 0582 | . 0434 | .as6s |  |  |  |  |
| 12.41 |  | -0566 | . 0419 | .0545 |  |  |  |  |
| 13.68 |  | .0553 | . 6406 | .e528 |  |  |  |  |
| 14.95 |  | .0541 | . 0395 | . 0512 |  |  |  |  |
| 16.22 | -1090 | .0530 | .0385 | .0199 | 995.3 |  |  |  |
| 17.49 | . 6778 | . 0580 | . 0376 | . 0486 | 836.7 |  |  |  |
| 18.76 | .0783 | .0511 | . 0368 | . 0475 | 992.6 |  | . 266971.12 |  |
| 20.63 | . 06875 | ${ }^{-6503}$ | -9360 | . 0465 | 995.2 |  | 4051E. 12 |  |
| 22.57 | . 6730 | -0488 | . 03547 | - $0+4 \times 8$ | 884.7 1012.2 |  | . $4872 \mathrm{E}+12$ |  |
| 23.84 | . 0774 | . 6481 | .0342 | . 0439 | 1134.2 |  | - 6836 E (12 12 |  |
| 25,11 |  | . 0475 | . 0336 | . 0432 |  |  | .6831E. 12 |  |


$494 \quad 165$



494 iol

## RUN NUMBER. . . 7509

average vold fract ION. ..... 32
INITIAL POOL DEPTH(CH) ..... 22.0
BOLI.ING POOL. DEPTIICM) ..... 32.5
VOL. PONER DENS, (CAL/CM3 SEC) ..... 21
SUPERFICIAL YAPOR VELOCITY. . ..... 59
ANGLE OF INCLINAT ION (DEGREES) . 15.00POOL VOLUME (CM**3) . . . . . . . . . . . . 14394(1.85.
PRANDTL NUMBER. ..... 1.85

1. TAL. POWER (KW) ..... 13. 44AVERAGE SURFACE TEMP (DEG C) . . . 90.8

| (Cl) | LOCAL HEAT TRANSFER COEFF <br> (CAL/CM2 SEC DEG C) |  |  |  | NUSSELT NUMBER | MODIFIED | RAYLEIGH | NUMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | EXFT | EO $14 a$ | $\begin{aligned} & E 0,10 \\ & N=1,0 \end{aligned}$ | $\begin{aligned} & E 010 \\ & \mathrm{~N}=0.7 \end{aligned}$ |  |  |  |  |
| 2246 | . 1917 | . 6917 | .0735 | . 1061 |  |  |  |  |
| 3.5 .6 | . 1770 | . 0820 | . 6683 | . 1061 | 264.6 382.5 |  | . 7096E*99 |  |
| 4. 736 | . 1608 | . 0759 | . 0616 | , 0815 | 482.5 |  | . $2722 \mathrm{E} \cdot 19$ |  |
| 6.056 | . 1647 | . 6715 | . 0569 | .0751 | 613.1 |  | . $6863 \mathrm{E}+18$ |  |
| 7.326 | . 1517 | . 0682 | . 0535 | . 0704 | 683.0 |  | -1390E+11 |  |
| 8.596 | . 1378 | .0655 | . 0508 | . 0667 | 727.9 |  | . $2467 \mathrm{E}+11$ |  |
| 9.866 |  | . 0633 | . 0486 | . 0636 | 727.9 |  | .3976E* 11 |  |
| 11. 14 |  | . 0614 | . 0467 | . 0611 |  |  |  |  |
| 12.41 |  | . 0598 | . 0451 | . 0589 |  |  |  |  |
| 13.68 |  | . 0584 | . 0437 | . 0570 |  |  |  |  |
| 14.95 16.22 |  | . 0571 | . 0425 | . 0554 |  |  |  |  |
| 16.22 17.49 | .1643 .0854 | . 05559 | .0414 | .0539 .0525 | 1039.1 917 |  | - $26698 \cdot 12$ |  |
| 18.76 | . 0882 | . 0539 | . 03045 | . 0525 | 917.4 1816.7 |  | . $3347 \mathrm{E} \cdot 12$ |  |
| 20.03 | . 0976 | . 0531 | . 0387 | . .9502 | 1816.7 1261.3 |  | $.4130 E+12$ $.5027 \mathrm{~F}+12$ |  |
| 21.39 | . 0781 | .0522 | . 0380 | . 0492 | 1201.3 1021.5 |  | $.5027 \mathrm{E}+12$ $.6045 \mathrm{E}+12$ |  |
| 22. 57 | . 0839 | . 0515 | . 03373 | . 0483 | 1151.4 |  | . 6045 E + 12 |  |
| 23.84 25.11 | .0881 | . 0508 | .0367 | . 0474 | 1290.1 |  | . $8477 \mathrm{E}+12$ |  |
|  |  | . 9501 | . 6361 | . 0466 |  |  |  |  |







## RUN NUMBER... 7512

average void fraction. . . . . . . . ..... 38
INITIAL POOL DEPTH(CM) ..... 20.0
BOIL.ING POOL DEPTII(CM) ..... 32.5
VOL. PONER DENS, (CAI/CM3 SEC). ..... 25
SUPERFICIAL VAPOR VELOCITY. ..... 65
ANGLE OF INCLINATION(DEGREES) ..... 15. 00
POOL VOLUME (CM**3) 12989
Prandtl number ..... 1.84
TOTAL POWER (KW) ..... 14.0
avERage Surface temp (LEG C)... ..... 91.5



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494 \quad 175
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$494 \quad 176$


## RUN NUMBER... 7514

AVERAGE VOID FRACTION......... . . . 52
INITIAL POOL. DEPTH(CM) . . . . . . . 15.0
BOILITK POOL. DEPTH(CM) . ........ 31.0
VOL. PONER DENS. (CAL/CN3 SEC). 38
SUPERFICIAL Vapor Velocity.... . 74
ANGLE OF INCLINATION(DEGREES) . 15.00
POOL. VOLUME (CH**3) . . . . . . . . . . . . 9560.6
PRANDTL NUMBER . . . . . . . . . . . . . . 83
TOTAL PONER (KW) . . . . . . . . . . . . . . . 20.5
AVERAGE SURFACE TEMP (DEG C) ... 93.2


RUN NUMBER... 7515
AVERAGE VOID FRACTION......... . . 52
NITIAI. POOL DEPTH(CM)......... is.
MOHLING POOL DEPTH(CM) ..... 31.6
VOL. PONER DENS, (CAL/CH3 SEC) ..... 42
SUPERFICIAL VAPOR VETOCITY....
ANGLE OF INCL.INATION(DEGREES) . $15.6 e$ ..... 15.60
POOL VOLUME (CK**3)
PRANDTL NUMBER. ..... 1.83
TOTAL PONER (KW) ..... 20.5
AVERAGE SURFACE TEMP (DES C) . . . ..... 92.7

| $\begin{aligned} & \text { DEPTH } \\ & \text { (CM) } \end{aligned}$ | LOCAL HEAT TRANSFER COEFF (CAL/CA2 SEC DEG C) |  |  |  | NUSSELT NUMBER | MODIFIED | RAYLEIGH |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| . 694 | . 2394 | . 1386 | . 1382 | . 1859 | 192.0 |  | . $3364 \mathrm{E}+08$ |
| 1.964 | . 1946 | . 1969 | . 0954 | . 1276 | 234.6 |  | .7635E*09 |
| 3.234 | . 1602 | . 0943 | . 0803 | . 1070 | 318.2 |  | . $3410 \mathrm{E}+16$ |
| 4.504 | . 2569 | . 0868 | . 0718 | . 0953 | 710.6 |  | .9211E+18 |
| 5.774 | . 1954 | . 0816 | . 0661 | . 0874 | 692.9 |  | . $1941 \mathrm{E}+11$ |
| 7.044 | . 1550 | . 0777 | . 0619 | . 0816 | 670.4 |  | . $3524 \mathrm{E}+11$ |
| 8.314 |  | . 0745 | . 0586 | .0771 |  |  |  |
| 9.584 |  | . 0719 | .0559 | . 0735 |  |  |  |
| 16.85 |  | . 0697 | .0537 | . 6765 |  |  |  |
| 12. 12 |  | . 0678 | .0518 | . 0679 |  |  |  |
| 13.39 |  | . 0661 | . 0502 | . 0656 |  |  |  |
| 14.66 | . 1335 | . 0646 | . 0487 | . 0637 | 1202.8 |  | . $3180 \mathrm{E} \cdot 12$ |
| 15.93 | . 1179 | . 0633 | . 04775 | . 0619 | 1145.4 |  | . $4079 \mathrm{E} \cdot 12$ |
| 17.29 | . 1298 | . 0621 | . 0463 | . 0664 | 1371.3 |  | . $5135 \mathrm{E}+12$ |
| 18.47 | . 1611 | . 0618 | . 0453 | . 0589 | 1827.5 |  | . $6358 \mathrm{E}+12$ |
| 19.74 | . 0888 | . 0609 | . 0443 | . 0576 | 1077.3 |  | .7761E•12 |
| 21.01 | . 1264 | .0591 | . 6435 | . 6565 | 1631.9 |  | .9357E+12 |
| 22.28 | . 18907 | .0582 | 0427 | . 0554 | 2473.1 |  | . $1116 \mathrm{E} \cdot 13$ |
| 23.55 | . 1636 | .0574 | .0419 | . 6544 | 2757.3 |  | 1318E* 13 |



$494 \quad 182$


|  |  |  |  |
| :---: | :---: | :---: | :---: |
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## RUN NIJMBER. . . 7518

| average void fraction. | . 34 |
| :---: | :---: |
| INITIAL POOL DEPTH(CW) | 13.6 |
| BOILING POOL. DEPTH(CW) | 29.6 |
| VOL. POWER DENS. (CAL/CM3 SEC). | . 73 |
| SUPERFICIAL VAPOR VELOCITY.... | 1.79 |
| ANGILE OF INCLINATION (DEGREES) . | 15.60 |
| POOL VOLLIME (CH**3) . . . . . . . . . . | 12293. |
| PRANDTL. NUMBER | 1.82 |
| TOTAL POWER (KW) | 44.6 |
| AVERAGE SURFACE TEMP (DEG C) | 94.5 |

 $\begin{array}{llll}{ }^{(C A L / C N 2 ~ S E C ~ D E G ~} \\ \text { EXPT } \\ & \text { E014a } \\ & & \mathrm{EO} 10 & \mathrm{E}=1.0 \\ & \mathrm{~N}=0.7\end{array}$

NUSSELT NUMBER
MODIFIED RAYLEIGA NUMBER
---

| 1.163 | .1184 | .1163 | .1323 | .1780 |
| :--- | :--- | :--- | :--- | :--- |
| 2.433 | .1802 | .0917 | .0998 | .1345 |
| 3.703 | .0973 | .0 | .0854 | .1150 |
| 4.973 | .1443 | .0 | .0766 | .1031 |
| 6.243 | .1913 | .0725 | .0766 | .0949 |
| 7.513 | .1434 | .0692 | .0660 | .0587 |
| 8.783 | .1161 | .0666 | .0624 | .0838 |
| 10.65 | .1241 | .0643 | .6595 | .0798 |
| 11.32 |  | .0625 | .0571 | .0765 |
| 12.59 |  | .0608 | .0550 | .0736 |
| 13.86 |  | .0594 | .0532 | .0711 |
| 15.13 | .1171 | .05 .51 | .0516 | .0683 |
| 16.40 | .2361 | .0569 | .0501 | .0676 |
| 17.67 | .0559 | .0489 | .0653 |  |
| 18.94 |  | .0549 | .0477 | .0637 |
| 20.21 |  | $.054 \theta$ | .0467 | .0622 |
| 21.48 |  | .0532 | .0457 | .0609 |


| 84.5 | $.1666 \mathrm{E}+09$ |
| ---: | ---: |
| 269.2 | $.3760 \mathrm{E} * 09$ |
| 221.3 | $.3441 \mathrm{E}+16$ |
| 449.7 | $.8334 \mathrm{E}+16$ |
| 733.3 | $.1649 \mathrm{E}+11$ |
| 661.7 | $.2874 \mathrm{E}+11$ |
| 626.9 | $.4591 \mathrm{E}+11$ |
| 766.3 | $.6885 \mathrm{E}+11$ |
|  |  |
|  |  |
| 1088.5 | $.2349 \mathrm{E}+12$ |
| 2562.4 | $.3741 \mathrm{E}+12$ |



|  |  | MODIFIED RAYLEIGH NUMBER |  |
| :---: | :---: | :---: | :---: |
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$494 \quad 191$

$494 \quad 192$




| AVERAGE VOID FRACTION. | .17 |
| :---: | :---: |
| INITIAL POOL. DEPTH(CM) | 25.9 |
| BOIL.ING POOL DEPTH(CW) | 30.0 |
| VOL. POWER DENS. (CAL/CM3 SEC). | . 15 |
| SUPERFICIAL VAPOR VELOCITY | . 48 |
| ANGIE OF InClinat Ion (DEGREES). | 30.09 |
| POOI, VOL UME (CY**3) . . . . . . . . . . . | 18278. |
| PRANDTL NUMBER | 1.88 |
| TOTAL. POWER (KW) . . . . . . . . . . . . . | 14.6 |
| AVERAGE SURFACE TEAP (DEG C) | 88.5 |


| i | $\begin{aligned} & \text { DEPTH } \\ & \text { (CM) } \end{aligned}$ | LOCAI. HEAT TRANSFER COEFF <br> (CAL/CM2 SEC DEG C) |  |  |  | NUSSEL.T NUMBER | MODIFIED | RAYLEIGH | GEMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 3. 241 | . 1783 | . 0706 | . 0596 | . 0795 | 355.4 |  | .9731E-69 |  |
|  | 4. 511 | .1242 .1795 | . 0650 | .0532 | .0768 | 344.5 |  | . $2624 \mathrm{E}+19$ |  |
|  | 7.051 | . 1271 | . 0581 | . 0458 | . 0666 | S50.9 |  | - 1602 E. 11 |  |
|  | 8.321 | . 1646 | . 0558 | . 0434 | . 0573 | \$35.9 |  | . $16-47 \mathrm{E}+11$ |  |
| $\xrightarrow{\sim}$ | 9.591 10.86 | . 1150 | .0538 | . 0414 | .0545 |  |  |  |  |
| 0 | 12.13 | . 1256 | . 0508 | . 0383 | . 0503 | 768.1 |  | $.3662 \mathrm{E} \cdot 11$ $.5103 \mathrm{E} \cdot 11$ |  |
| $\stackrel{\square}{7}$ | 13.49 | . 0712 | . 0495 | . 0371 | . 0487 | 586.8 |  | . 6879 E - 11 |  |
|  | 14.67 15.94 | . 0775 | . 04848 | . 0369 | .0472 .0459 | 699.0 |  | $.9026 \mathrm{E}+11$ |  |
|  | 17.21 | . 0762 | .0-165 | . 0342 | . 0447 | 793.1 |  | . $1158 \mathrm{E}+12$ |  |
|  | 18.48 | . 0819 | . 0457 | . 0335 | . 0437 | 930.7 |  | .1804E.12 |  |
|  | 19.75 21.02 | . 0729 | .0449 .0443 | .0328 | . 0427 | 874.6 |  | . $2202 \mathrm{E}+12$ |  |
| B | 22.29 | . 9758 | . 0436 | . 0315 | . 0416 | 1038.5 |  | $.2655 E+12$ $.3156 E+12$ |  |
|  | 23.56 24.83 | . 0834 | .0439 .0424 | .0319 .0305 | 0403 0396 | $12 / 3.9$ |  |  |  |
|  | 26.10 |  | .0419 | . 0300 | .0389 |  |  | .4376E+12 |  |



$494 \quad 198$

REX NUMBER. . . 6004
aVERAGE VOID FRACTION. ..... 11
INITIAL. POOL. DFPTH(CM) ..... $25 . \theta$
BOIL ING PGOL DEPTH(CM). ..... 30.0
VOL. POWER DENS. (CAL CCM3 SEC) .....  15
SUPERFICIAL VAPOR VEIOCITY... ..... 48
ANGLE OF INCLINATION(DECBEES)
OOL VOLUNE (CH**3) ..... 18278.
PRANDTL NUMBER. ..... 1.89
TOTAL POWER (KV) ..... 12.6
AVERAGE SURFACE TEMP (DEG C). ..... 87.8

| $\begin{aligned} & \text { DEPTH } \\ & \text { (CK) } \end{aligned}$ | $\begin{array}{r} \text { LOCAL } \\ \text { (C) } \end{array}$ | beat <br> JCN2 | ANSFER <br> EC DEC | COEFF | NUSSELTT NHMBER | MODIFIED | Rayleigu | NUMBER |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | EXPT | E0 14a | $\begin{aligned} & E 0 \quad 10 \\ & \mathrm{EO}=1.0 \end{aligned}$ | $\begin{aligned} & \text { EO } 10 \\ & \mathrm{~N}=0.7 \end{aligned}$ |  |  |  |  |
| 3.241 | . 1852 | .0705 | . 0595 | . 8794 | 369. |  |  |  |
| 4.511 | . 1511 | . 0649 | . 0532 | .8767 | 869.2 |  | .969 IE+09 |  |
| 5.781 | . 1162 | . 0618 | . 0489 | . 0649 | 419.3 413.3 |  | . $2613 \mathrm{E}+10$ |  |
| 7.651 | . 1186 | .0581 | . 0458 | . 0606 | 413.3 514.5 |  | . S509E+ 10 |  |
| 8.321 | . 60552 | -0557 | . 6433 | .0572 | 514.5 |  | .9979E+ 16 |  |
| 9.591 | . 1167 | . 05538 | . 0413 | . 0545 | 683.5 |  | 1640E, 11 |  |
| 10.86 | . 0870 | .0521 | .0397 | .0527 | 683.2 |  | . $2511 \mathrm{E}+11$ |  |
| 12.13 |  | . 0597 | . 0383 | . 0503 | 581.2 |  | +3647E+11 |  |
| 13.49 | -0565 | $0-495$ | . 6371 | . 0486 |  |  |  |  |
| 14.67 15.94 | .0762 .0763 | . 0484 | . 0364 | . 0472 | 5487.6 |  | .6851E* 11 |  |
| 15.94 17.21 | .6763 .19757 | + 0.0474 | . 0350 | . 0459 | 748.5 |  | .8989E* 115 |  |
| 18.48 | -9774 | -0456 | . 0.0342 | +0447 | 801.4 |  | $1451 \mathrm{E}+12$ |  |
| 19.75 | . 9716 | . 6449 | . .0327 | -0437 | 879.7 |  | $1797 \mathrm{E} \cdot 12$ |  |
| 21.02 | . 0769 | . 0442 | . 0321 | . 0418 | 869. |  | $2193 \mathrm{E} \cdot 12$ |  |
| 22.29 | . 0754 | . 9436 | . 0315 | . 0418 | 1033.6 |  |  |  |
| 23.56 24.83 | .0794 | -0430 | . 0309 | . 0462 | 1151.0 |  | - $3753 \mathrm{E} \cdot 12$ |  |
| 24.83 26.10 | .6827 | .0424 .0419 | . 0304 | .0395 .6389 | 1262.7 |  | $4358 \mathrm{E} \cdot 12$ |  |




494203






## RUN NUMшER . . . 6009

average void fraction.......... . . 37
... ..... 19.8
MOILING POOL DEPTHICH
VOL. PONER DENS (CAL/CM3 SEC) 37
SUPERFICIAL. VAPOR VFIOCITY.... . 92
ANGLE OP INCLINATION(DEGREES) . 30.60
POOL VOLIME (CH**3) . . . . . . . . . . . . 13299 ..... 3299
PRANDTL NUABER ..... 1.85
TOTAL POWEK (KW)average surface terp (DEG C)... 91.
ver II
(CM)
$\cdots . .$.
3.241
4.51
5.78
7.05
8.32
9.59
16.86
12.13
13.4
14.6
15.9
17.2
18.4
19.7
21.6
22.29
23.5
24.8
26.10

LOCAL HFAT TRANSFER COEF
EXPT ECAL EO 2 SEC DEG C)
$\begin{array}{llll}\text { EXPT } & \text { EO } & 14 \mathrm{a} & \mathrm{FO} 10 \\ \mathrm{~N}=1.0 & \mathrm{EOP10} \\ \mathrm{~N}=\mathrm{a} .7\end{array}$

| .1984 | .0863 | .0765 | .1025 |
| :--- | :--- | :--- | :--- |
| .1965 | .0795 | .0682 | .0911 |
| .1835 | .0747 | .06266 | .0835 |
| .1291 | .9711 | .0585 | .0779 |
| .1192 | .0682 | .0553 | .0735 |
| .1167 | .0658 | .0527 | .0699 |
| .1227 | .0638 | .05066 | .0670 |
| .1107 | .0621 | .0488 | .0645 |
| .0969 | .0605 | .0472 | .0623 |
| .1076 | .0592 | .0458 | .0604 |
| .1197 | .0580 | .0445 | .0587 |
| 1172 | .0569 | .0434 | .0572 |
| 1167 | .0559 | .0424 | .0558 |
| 1095 | .0549 | .0415 | .0546 |
| 1319 | .0541 | .0407 | .0534 |
|  | .0533 | .93999 | .0524 |
|  | .0526 | .03992 | .0514 |
|  | .0519 | .0512 | .0376 |
|  | .0505 |  |  |
|  |  | .0497 |  |

NUSSELT NUMBER MODIFIED RAYLEIGH NIMBEA
$\qquad$
$\qquad$
395.2
528.2
651.8
559.6
609.6
652.6
819.3
825.5
799.5
979.2
1172.5
1239.6
1257.6
1329.3
1793.6
$.2166 \mathrm{E}+10$
$.5841 \mathrm{E}+10$
$.1229 \mathrm{E}+11$
$+22291 \mathrm{E}+11$
$-2266 \mathrm{H}^{2}$
$.3666 \mathrm{E} \cdot 11$
$.5614 \mathrm{E}+1$
$-8152 \mathrm{E}+1$
$.8152 \mathrm{E}+1$
$.1136 \mathrm{E} \cdot 12$
$-1531 \mathrm{E}+12$
$.153 \mathrm{E}+12$
$.2009 \mathrm{E} \cdot 12$
$.2009 \mathrm{E} \cdot 12$
$.3248 E+12$
$.4016 \mathrm{E} \cdot 12$
$.4903 \mathrm{E}+12$
. $5916 \mathrm{E}+12$





## RUN NUMBER... 6012

avERAGE VOID FRACTION ..... 19
INITIAL. POOL DEPTHICA ..... 21.0
BOILING POOL DEPTIICM ..... $26 . \theta$
VOL POWER DENS. (CAI/CM3 SEC) ..... 60
STPEREICIAL VAPOR VETACITY.... 1.64ANGLE OF IACLINATION(DEGREES) . 30.00POOL. VOLUME (C) * * 3) . . . . . . . . . . . 14917PRANUTL. NUMBER . . . . . . . . . . . . . . . 1.85
TOTAL. POWER(SW) . . . . . . . . . . . . . . . . 45.6AVERAGE SURFACE TEMP (DEG C $1 . . .92 .2$
1.OCAI HEAT TRANSFER COEFF EXPT EO 14a EO $10 \quad \mathrm{E} 010$
$\qquad$

NODIFIED RAYIFIGI NUMBER


|  | 1. 162 |  | . 6950 | . 1193 | . 1602 |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | 2.432 |  | . 0790 | . 0895 | . 1295 |
|  | 3.792 | . 1583 | . 9711 | . 0763 | . 1628 |
|  | 4.972 | . 1560 | . 0661 | . 0683 | . 0929 |
|  | 6.242 | . 1889 | . 9624 | . 0628 | . 0845 |
|  | 7.512 | . 2082 | . 0596 | . 6587 | .0790 |
|  | 8.782 | . 1375 | . 0573 | . 0558 | .0745 |
| $\checkmark$ | 19.05 | . 1423 | .0554 | .0528 | . 6769 |
|  | 11.32 | . 1366 | . 0538 | . 0505 | . 0679 |
|  | 12.59 | . 1227 | . 0524 | . 0486 | . 6653 |
|  | 13.86 | . 1952 | . 0511 | . 0477 | . 0631 |
|  | is. 13 | . 0976 | 6506 | . 0456 | . 0611 |
|  | 16.40 |  | . 0490 | . 0443 | . 0593 |
|  | 17.67 |  | . 0481 | . 6431 | . 6578 |
| $\sigma$ | 18.94 |  | . 0.473 | 0421 | . 6564 |
|  | 20.21 |  |  | 0411 | .0551 |

### 341.9 476.5 <br> 476.5 724.7 <br> 961.2 742.1 <br> 742.1 878.8 <br> 878.8 950.3 <br> 950.3 949.5 <br> 949.5 1662.5

907.2


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494 \quad 217 \quad-171-
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## APPENDIX D

## Algebraic Development of Equations 28 and 29.

In single phase natural convection heat transfer along a vertical flat plate, it has been determined that for Pr equal 2.0 , the local free convection heat transfer correlations for laminar and turbulent flow are

$$
\begin{equation*}
\mathrm{Nu}(\mathrm{x})=0.42\left[\mathrm{Gr} \cdot \mathrm{Pr}^{1 / 4}\right. \tag{D-1}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{Nu}(\mathrm{x})=0.0295 \mathrm{Gr}^{2 / 5} \mathrm{Fr}^{7 / 15}\left[1+.494 \mathrm{Pr}^{2 / 3}\right]^{-2 / 5} \tag{D-2a}
\end{equation*}
$$

respectively. Equation D-2 may be reduced to the following approximate form by grouping $(\mathrm{Gr} \cdot \mathrm{Pr})^{2 / 5}$ and letting $\mathrm{Pr}=2.0$;

$$
\begin{equation*}
\mathrm{Nu}(\mathrm{x})=0.0245[\mathrm{Gr} \cdot \mathrm{Pr}]^{2 / 5} \tag{D-2~b}
\end{equation*}
$$

In order to compare these results to average experimental heat transfer data, it has been necessary to convert these local heat transfer coefficlients into the corresponding average values along the surface. By introducing the expression for the Grashof number into the above equations, it can be seen that the local heat transfer coefficient is proportional to the distance from the leading edge along the laminar boundary layer to the $-1 / 4$ power ${ }^{(25)}$, i.e.,

$$
\begin{equation*}
h(x) \sim x^{-1 / 4} \tag{D-3a}
\end{equation*}
$$

and for the turbulent case, proportional to the distance from the leading edge along the boundary layer to the power $0.2^{(26)}$, 1.e.,

$$
h(x) \approx x^{0.2}
$$

It may be readily shown by integrating equations $D-1$ and $D-2 a, b$ that the average heat transfer coefficient over the length of the plate, $L$, is related to the local coefficient at $x=L$ as follows: for the laminar boundary layer (Eq. D-1),

$$
\begin{equation*}
\bar{h}=\frac{1}{L} \int_{0}^{L} h(x) d x=\left.\frac{4}{3} h(x)\right|_{x=L} \tag{D-4a}
\end{equation*}
$$

and for the turbulent boundary layer (Eqs D-2 $a, b$ ),

$$
\begin{equation*}
\bar{h}=\frac{1}{L} \int_{0}^{L} h(x) d x=\left.\frac{h(x)}{1.2}\right|_{x=L} \tag{D-4b}
\end{equation*}
$$

Thus, the local Nusselt relations (Eqs. $D-1$ and $D-2 a, b$,) may be modified to represent the average heat transfer behavior by substituting $\overline{\mathrm{h}}$ for $\mathrm{h}(\mathrm{x})$, L for x , and multiplying the correlation coefficient, K , by the factor 1.33 (for laminar case) or .83 (for turbulent case). When this is done, Eqs. $\mathrm{D}-1, \mathrm{D}-2 \mathrm{a}$, and $\mathrm{D}-2 \mathrm{~b}$ are transformed into the average correlations below, respectively,

$$
\begin{align*}
& \overline{\mathrm{Nu}}=0.56[\mathrm{Gr} \cdot \mathrm{Pr}]^{1 / 4}  \tag{D-5}\\
& \overline{\mathrm{Nu}}=0.0246 \mathrm{Gr}^{2 / 5} \mathrm{Pr}^{7 / 15}\left[1+.494 \mathrm{Pr}^{2 / 3}\right]^{-2 / 5}  \tag{D-6a}\\
& \overline{\mathrm{Nu}}=0.0210[\mathrm{Gr} \cdot \mathrm{Pr}]^{2 / 5} \tag{D-6~b}
\end{align*}
$$

These correlations are in excellent agreement with empirically derived relations in the literature.

The above conversion procedure is valid only in the case that the free stream properties are constant. In the case of boundary layer heat transfer

$$
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$$

from volume-heated boiling pools, it has been shown that there exists a vertical distribution of vapor fraction which is defined by the solution to Eq. 20. Thus, the integration demonstrated in Eqs. $D-4 a, b$ will depend not only on the explicit spatial dependence in the Grashof number, but upon the implicit spatial variation $\quad$ a the void fraction distribution as well.

The forms of the local void distribution that will be considered are listed below:
a. $\alpha(x)=\bar{\alpha}$
b. $a(x)=2 \bar{\alpha}(x / L)$
c. $\alpha(x)=1.5 \bar{a}(x / L)^{1 / 2}$
d. $\alpha(x)=3 \bar{a}(x / L)^{2}$

The procedure for determining the ratio of the average heat transfer correlation based upon $\bar{\alpha}$ to the local heat transfer correlation based upon $\alpha(x)$ is identical to that previously described. The results of the intergration are presented in Table D-1 for both the laminar case ( $\mathrm{Nu} \approx \mathrm{Ra}^{1 / 4}$ ) and the turbulent case $\left(\mathrm{Nu} \sim \mathrm{Ra}^{2 / 5}\right)$.

The local heat transfer data from Ref. 12 has been correlated by both local and average void fraction with little observed sensitivity in the result, as demonstrated by Eqs. $14 \mathrm{a}, \mathrm{b}$. On this justification, the local heat transfer data from this work was correlated on the basis of the average void fraction. As yet, only preliminary correlations of the data have been performed for laminar bubbly flow heat transfer. The results indicate quite close agreement with case d. Final results for local correlation of both laminar and turbulent bubbly flow data will be reported in the future.

## SUMMARY OF LOCAL-TO-AVERAGE HEAT TRANSFER CORRELATION

 CONVERSION UPON FREE STREAM VOID DISTRIBUTION|  | $\mathrm{K}_{\mathrm{ave}} / \mathrm{K}_{\text {local }{ }^{*}}$ |  |
| :---: | :---: | :---: |
| VOID DISTRIBUTION FUNCTION | LAMINAR | TURBULENT |
| $\alpha(\mathrm{x})=\bar{\alpha}$ | 1.33 | .83 |
| $\alpha(\mathrm{x})=2 \bar{\alpha}(\mathrm{x} / \mathrm{L})$ | 1.19 | .82 |
| $\alpha(\mathrm{x})=1.5 \bar{\alpha}(\mathrm{x} / \mathrm{L})^{1 / 2}$ | 1.26 | .84 |
| $\sim(x)=3 \bar{\alpha}(x / \mathrm{L})^{2}$ | 1.05 | .78 |

${ }^{*} K_{\text {ave }}$ is the average heat transfer correlation coefficient,
$\overline{\mathrm{Nu}}=\mathrm{K}_{\text {ave }}\left(\mathrm{Gr}^{*} \cdot \mathrm{Pr}\right)^{\mathrm{n}}$ where $\mathrm{n}=1 / 4$ for laminar flow, $\mathrm{n}=2 / 2$ for turbulent flow.
${ }^{*}{ }_{\text {ducal }}$ is the local heat transfer correlation coefficient, $\mathrm{Nu}(\mathrm{x})=$ $\mathrm{K}_{\text {local }}\left(\mathrm{Gr}^{*} \cdot \mathrm{Pr}_{\mathrm{r}}\right)^{\mathrm{n}}$ where $\mathrm{n}=1 / 4$ for laminar flow, $n=2.5$ for turbulent flow.

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BNL RSP Division Heads
BNL RSP Group Leaders
BNL RSE Person: 1
P. Abramson, ..... ANL
A. Alter, ..... DOE
L. Baker, ANL
S. G. Bankoff, Northwestern Uni.
D. Basdekas, ..... NRC
J. Boudreau, ..... LASL
I. Catton, Uni. of California
J. C. Chen, Lehigh Uni.
R. Coates, Sandia Laboratory
R. T. Curtis, NRC
A. Dukler, Uni, of Houston
D. T. Eggen, Northwericern Uni.
E. Epstein, ANL
J. Gabor, ANL
W. Gammill, NRC
R. Henry, ..... ANL
H. H. Hummel, ANL
W. Y. Kato. BNL
M. S. Kazimi, MIT
C. N. Kelber, NRC
H. J. Kouts, BNL
T. Kress, ORNL
F. Kulacki, Ohio State Uni.
J. T. Larkins,
J. Martin, HEDL
R. Ostensen, Sandia Laboratory
A. Reynolds, Uni, of Virginia
M. Silberberg, ..... NRC
R. Stein, ..... ANL
M. Stevenson, LASL
H. Todosow, ..... BNL
T. G. Theofanus, Purdue Uni.
J. C. Walker, Sandia Laboratory
R. W. Wright, NRC
U.S. NRC Division of Technical Information and Control


[^0]:    ${ }^{* *}$ In Ref. 12, the exponent of the Prandti number appeared as $1 / 2$ instead of $1 / 3$ as suggested by Rays ${ }^{24}$

[^1]:    * SVV indicates the dimensionless superficial vapor velocity.

[^2]:    * SVV indicates the dimensionless superficial vapor velocity

[^3]:    494
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