

International Agreement Report

Relevant Results Obtained in the Analysis of LOBI/MOD2 Natural Circulation Experiment A2-77A

Prepared by F. D'Auria, G. M. Galassi

Universita' Degli Studi Di Pisa Dipartimento Di Costruzioni Meccaniche E Nucleari Pisa, Italy

Office of Nuclear Regulatory Research U.S. Nuclear Regulatory Commission Washington, DC 20555

April 1992

Prepared as part of The Agreement on Research Participation and Technical Exchange under the International Thermal-Hydraulic Code Assessment and Application Program **(ICAP)**

Published by U.S. Nuclear Regulatory Commission

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and

National Technical Information Service Springfield, VA 22161

NUREG/IA-0084 NT 163 (90)

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Abstract

The present document describes the activities carried out at Pisa University to assess the RELAP5/MOD2 performance in the application to the natural circulation test A2-77A performed in LOBI/MOD2 facility.

Sensitivity calculations have been performed in this context, with the attempt to distinguish the code limitations from the uncertainties of the measured conditions.

The characterization of instabilities in two-phase natural circulation and the evaluation of the user effect upon the code results are special goals achieved in the frame of the A2-77A analysis. Both of these are discussed in sect. 5.

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TABLE OF CONTENTS

LIST OF FIGURES IN THE TEXT Page

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 $\mathcal{A}^{\mathcal{A}}$

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vii.

Page

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LIST OF TABLES IN THE TEXT Page

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ACKNOWLEDGEMENTS

Part of the work has been carried out in the framework of the European Community Shared Cost Action Program (contract 2996-06 EL ISP I).

The CEC contribution is gratefully acknowledged.

The authors wish to acknowledge the stimulating discussions with LOBI researchers Mrs. G. De Santi, L. Piplies, **H.** Stadtke, and B. Worth and with Prof. Mazzini of DCMN.

Mr. G. Fantappiè contributed to the analysis with basic models, and Mr. P. Lombardi performed most of the work related to the evaluation of user effect on the results.

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1. INTRODUCTION

In the frame of the system codes assessment activities $/1/2/$, the DCMN of Pisa University performed the post test analysis, by RELAP5/MOD2 code, of the A2-77A test carried out in the LOBI/MOD2 experimental facility.

LOBI/MOD2 is a PWR experimental simulator installed at JRC of Ispra.

A2-77A is a natural circulation test including 1-phase, 2-phases and reflux condensing modes of natural circulation.

Actually the experiment consists of two parts wich differentiate owing to the pressures of primary and secondary sides (9.0 and 7.5 HPa for the PS pressure in two cases, respectively). Having been recognized that measured phenomena are essentially the same in the two parts of the test, it has been decided to analyze only the former part of the test (PS pressure equal to 9.0 MPa) in the present framework.

The purpose of the analysis is to assess the capabilities of an advanced code in predicting the various phases of natural circulation in a typical PWR situation; in particular mechanical non equilibrium phenomena (CCFL, stratification, etc.) are important in the test together with asymmetries in geometric and boundary conditions of the loops connected with the RPV. Besides, the values of fluid velocities and of temperature differences among the various zones of the loop, are relativelly small (much lower than typical values characterizing the nominal condition of the loop with pumps in operation) making very critic the achievement of steady conditions in the code calculations.

The instabilities occurring in the main quantities during the twophase period and the influence of user upon the code results are discussed into detail.

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2. EXPERIMENT DESCRIPTION

2.1 - LOBI/MOD2 test facility

The LOBI facility simulates the KWU PWR plant of Biblis (FRG). The facility (LOBI-MODI) was built in the frame of a R&D contract between BMFT and CEC; at present it is operated by CEC.

The primary circuit is approximately $1/700$ scale model of the four loops of Biblis reactor and consists of the vessel and of two loops: the triple (intact) loop representing three loops and the single loop (broken), representing one loop of the reference plant. The facility is schematically shown in Fig. **1.**

The core simulator is constituted by 64 directly heated rods, arranged in a 8x8 square matrix, having the same geometrical dimensions of that in real plant: nominal heating power is 5.3 Mw.

The operating conditions (PS and SS pressures and temperatures, fluid velocities inside the bundle and the SG tubes, etc.), are typical of PWR systems.

HPIS and accumulator injections are provided in both HLs and CLs; AFW and further plant specific features (PRZ sprays and heaters, etc.), are included in the facility.

Almost 50 experiments, comprising LBLOCAs, SBLOCAs and Special Transients, have been performed up to now, roughly in ten years of operation.

A2-77A is one of the characterization tests of the facility.

2.2 - Test objectives and system configuration

The main objectives of the test are */3/:*

- a) to investigate the LOBI-MOD2 test facility characteristics during natural circulation conditions;
- b) to analyze how the different energy transport modes of natural circulation settle in LOBI i.e:
	- subcooled single-phase natural circulation;
	- satured single-phase natural circulation;
	- two-phase natural circulation;
	- reflux condenser mode;
- c) to determine the transitions between the various modes as a function of primary system mass inverntory;
- d) to investigate the instrumentation capability to detect low natural circulation mass flow rates, small differential pressures and temperatures;
- e) to compare LOBI results with that of other natural circulation experiments, such as those performed in PKL, Semiscale, Flecht.

Test A2-77A was performed with the integral two-loop configuration.

The main data on the system configuration are summarized in Tab. I.

In particular it should be noted that PRZ was used only to achieve the initial conditions; afterward it was valved out. Similarly, **UH** was removed mainly to avoid condensation in structures due to heat losses during the draining periods.

MCP were at zero speed during the whole test; no pump seal water was applied during the test. In order to limitate the asymmetries between the two loops, a "locked rotor resistance simulator" was inserted in the BL, downstream the MCP. The hydraulic characteristics of this simulator are shown in Tab. II.

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Fig. 1 - Scheme of LOBI-MOD2 facility

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Tab. I - LOBI/MOD2 system configuration for A2-77 test

The resistance simulates together with the resistance of the pump at zero speed, the total locked rotor resistance. Information on resistance. characteristics of -this additional "Locked Rotor Resistance Simulator" is given in the following. Δp -measurements were performed with water (t = 22 to 25^o $= 997 \text{ kg/m}^3$ The results are: ζ Measurement No. | Re 1 $0.716*10₅⁵$ 73. 1 $0.716*10^5$ 73.1
2 $1.13*10^5$ 69.7 3 1.55 **10** 68.3 $4 \t\t 1.77 *10⁵₅ \t 66.3$ 5 \vert 1.99 \cdot 10⁵ \vert 65.3 With $\Delta p = \zeta \frac{\zeta}{a} w^2$ Re = $\frac{w * p}{2}$

... mean velocity in the broken loop

Di ... inner diameter of broken loop (46.1 mm)

Tab. II - Hydraulic characteristics of BL "locked rotor resistance simulator"

6.

2.3 **-** Initial and boundary conditions

The test started with PS pressure at 14. MPa and SS pressure at 8.65 MPa. Draining of water from LP was necessary to evaluate the loop behaviour in several natural circulation conditions. In particular, the SS pressure led to a PS pressure of about 9.0 MPa when saturation conditions where reached in PS (Fig. 2).

After each draining step the primary system was allowed to stabilize at the new conditions. The primary pressure decreased rapidly and after few draining steps reached the foreseen value of 90 bar. The test continued through saturated single-phase and two-phase natural circulation and it terminated with reflux condenser mode. The test was finished when dry-out phenomena in the uppermost sections of the core occurred, due to low level (52% of primary inventory) in the RPV.

The complete list of initial conditions for the test is given in Tab. III.

Further boundary conditions are ad follows:

- a) the heating power remained constant at about 183 Kw; the axial distribution is given in Fig. 3.
- b) the triple and single loop SGs were isolated throughout the test; "feed and bleed" procedure was adopted to control the secondary pressure and steam generator water level; the water level in the SG steam dome was regulated to remain costant at the initial nominal elevation by using the AFW system; the SS pressure relief valves acted at about 8.65 MPa for Part **1.** The pressure was kept costant during each single part of the test via the secondary relief valves; the PS pressure stabilized out at about 9.0 MPa;
- c) as already mentioned, transition between the various modes of natural circulation (single-phase, two-phase, reflux condensation) was obtained by reducing stepwise the primary inventory. Discrete amounts of water were drained from the vessel lower plenum for each step, condensed and measured in a catch tank. The system mass inventory was varied in increments of about 1% to 3% of the total initial system mass. After each draining step sufficient time was given to the primary system to stabilize out. The time duration of each cycle (including the draining period, the stabilization period and the steady state period) varied between 25 and 30 minutes. The draining mass flow was about $1 \ \text{kg/min}$ which represented a compromise between technical requirements of a slow draining process and a limitation of the test time duration.

The PS was affected by a small amount of continuous fluid leakage. This leakage could not be visually observed during the test.

To quantify the total amount of fluid lost by the primary loop, the primary system was completely refilled after the end of the test. Two independent methods were used to quantify this originally unknown mass leak *131.* On this basis the overall mass inventory decrease from the PS (drained water plus leaks) was estimated and the results are shown in Tab. IV $/4$; in particular the sum of the values included in columns A and B of Tab. IV have been used as input for code calculation (17 time steps).

Fig. 2 - Primary and Secondary Pressure during Part **¹**

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Tab. III **-** Initial conditions for test A2-77

Fig. 3 - Axial Power Distribution prior to Initiation of Test

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Tab. IV - History of drainage from LP during A2-77 test \vdash

 $-\cdot$ \sim \sim 100 minutes and the second contract of the \sim

3. ADOPTED NODALIZATIONS **AND** PERFORMED SENSITIVITY ANALYSES

3.1 - Nodalization

The basic nodalization is the one adopted for ISP 18 post test calculation */5/* (Fig. 4). A further nodalization has been specifically set up including two different U-tubes in order to assess the influence of differential elevations of U-tubes in SGs, considering some conclusions from experimental data analysis 16/, /7/. The main characteristics of the nodalizations are summarized in Tab. V.

Few aspects are common to the performed sensitivity calculations:

- 17 main "draining steps" are included in the code input: during each of them the same amount of water is drained from the lower plenum as in the experiment; nevertheless the time duration of each "steady state" period is lower than in the test owing to the need to save CPU time; obviously it has been checked that a reasonable steady situation was reached in the code calculation after the end of the various "steady state" periods (the observation of some experimental boundary conditions, e.g. vessel wall temperature, also suggested this choice);
- additional "draining steps" were added in the code input in order to arrive at dryout conditions in core rods, so that the overall time duration of the calculated test was more than 4000 seconds; in Tab. VI the draining steps used ad input to the code are specified (the comparison between data in Tabs. IV and VI allows one to fix the corresponding periods in the experiment and in the calculation as far as the PS mass inventory is concerned);
- **-** difficulties were encountered in achieving a satisfactory steady state situation (see also below);
- valves are included in various zones of the nodalization to isolate PRZ and UH from PS during the test.

3.2 - Analyzed cases

Several attempts were made to achieve initial steady state situation in the code calculation.

Afterward, in agreement with discussions had with Lobi researchers /8/, the cases reported in Tab. VII have been analyzed.

The main purposes of the sensitivity calculations are as follows:

case A): to achieve steady state conditions adopting fluid temperature distribution measured in the experiment \int 3 \int boundary and initial conditions defined in Tabs. I, II and II, and other conditions as in ISP 18 calculation /5/ including in particular heat losses to environment and localized pressure losses distribution;

case B): case A) led to inconsistencies between temperature distribution and flowrate in the various zones of PS, so the aim of case B) was to achieve the same flowrates as in the experiment by varying some localized pressure losses coefficients;

case C): to check the influence of pipes roughness on flowrate;

case D): to check the influence of incondensable gases in natural circulation;

case E): to check the influence of considering two different heights of Utubes in the nodalization.

Fig. 4 - RELAP5/MOD2 nodalization for A2-77 test **(1** U-tube, cases A to D)

Tab. V **-** Details of the nodalization used for RELAP5/MOD2 calculations (R5-A and R5-E models)

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Tab. VII - Significant characteristics of A2-77 calculations $\begin{array}{c} \n\ddots \n\end{array}$

4. COMPARISON BETWEEN MEASURED AND CALCULATED TRENDS

4.1 - Achievement of steady state

The initial part of the activity was devoted to the search of a steady state situation in the trends predicted by the code; in particular the effect was checked of:

1) fluid temperature distribution in SG SS;

2) heat losses in PS;

3) heat losses in SS;

4) initial temperature imbalance between fluid and structures in PS.

In all cases PRZ pressure and liquid level, SGs SS pressure and DC levels were maintained costant through proper time dependent volumes and junctions: the DC level in SS was maintained costant during the whole test by the AFW flowrate assumed at a temperature of 25 °C (same as in the experiment).

The coupling between natural circulation flowrate and fluid temperature distribution in PS including heat transfer from PS to SS was studied for each of the boundary conditions **1)** to 4).

The consideration of subcooling in the DC of SGs (case **1))** leads to an heat exchange between PS and SS much greater than the core power, thus causing much larger natural circulation in the PS loops (the resulting power imbalance was partially compensated by the TMDPVOL connected with PRZ). The assumption of (roughly) initial saturation conditions in the DC of SGs led to the decrease of PS flowrate and to a better agreement with experimental data. In this case (case A in Tab. VII) the PS flowrate was roughly two times the measured value.

Doubling the values of heat losses in PS and SS (cases 2) and 3)) caused variations of the PS flowrate roughly around 10% of the experimental value (thus insufficient to match the measured values).

The influence of the initial temperature distribution in the structures was much stronger and proper tuning (t **1°C** in almost all structures) could have had been used to match the measured values of flowrate in PS. This choice would have had unknown influence in the subsequent part of the transient, so the decision was taken to assume all the structures in equilibrium with the fluid in the steady state conditions (apart from the consideration of heat losses).

In the following paragraphs cases A, B, C, D and **B** of Tab. VII are discussed. It should be noted that no result is available for case D due to the failure of the code in evaluating non condensible gases. Still case B must be considered as the reference calculation.

4.2 - Case A (Initial calculation)

The analyses performed to achieve the initial conditions (sect. 4.1) led to the following decisions concerning unspecified or unclear boundary and initial conditions:

- **-** heat losses in PS and SS as in ISP 18;
- initial fluid temperature in SS of **SG** assumed at saturation;
- **-** HL fluid temperatures in both BL and IL equal to the measured values in IL;
- CL fluid temperatures in BL and IL as specified from the experiment;
- localized pressure loss coefficients and pipes roughness same as in ISP **18.**

Significant results from steady state are shown in App. A. It can be noted that both IL and BL flowrate values are more than 50% higher than the experimental values. This result suggested to not continue the calculation

without adjusting some input parameters.

4.3 - Case B (Reference calculation)

The results of case A demonstrated that some tuning in input data was necessary in order to make meaningfull the comparison between calculated and measured trends. Various possibilities can be used in this regard even remaining within the uncertainties limits of measured parameters (e.g. pressure drops, fluid temperatures, structures temperatures, DC level in SG SS, losses in the PRZ isolation valve, etc.) $(*)$.

For the sake of simplicity it was decided to vary, with respect to case A, only the localized pressure loss coefficient at - core inlet;

- HLs inlet in BL and IL;

- SGs PS U-tube inlet in BL and IL.

The achieved steady state situation is compared with experimental values in Figs. 5, 6, 7 and 8 and in Tab. VIII. Further information about steady state can be obtained from App. B where several variables trends are reported for the whole transients. Still in relation to steady state, a CNTRLVAR was set up as the (algebric) sum: CORE POWER - POWER EXCHANGE ACROSS U-TUBES OF SGs - HEAT LOSSES FROM PS (EXCLUDING PRZ).

It has been checked that the absolute value of this quantity was less than **10** Kw, in steady situation, at the beginning of the simulated test.

The resulting values of flowrate as a function of PS residual mass are compared with experimental data in Figs. 9 and **10** with reference to IL and BL, respectively. The PS pressure trends are compared in Fig. **11;** a list of significant measured and calculated events or quantities during A2-77 is reported in Tab. IX (for brevity in this Table some results from case E have also been included). In particular some characteristics of the oscillations of fluid velocities and densities measured when PS mass inventory equals 285 Kg are also compared with calculated data.

The analysis of the reported trends allows one to conclude that the code predicts quantitatively well the important phenomena occurring during the test. With reference to the oscillations measured during the test, the reported figures (see also App. B) demonstrate that oscillatory behaviour also results from calculations: the related frequency has the same order of magnitude as the measured one. Nevertheless a more in depth analysis has been performed in relation to the nature of these oscillations also including the evaluation of the measurement error (sect. 5.1).

4.4 - Case C (Piping roughness changed)

The objective of calculation C is to assess whether tuning on roughness has the same effect on the results of variations in localized pressure losses coefficients **/8/.** To this aim taking as reference the case A, the pipe roughness already utilized in ISP 18 analysis was multiplied by two. Significant results are shown in App. C.

⁽M) Later analyses **191,/10/** of the experimental data demostrated that the most plausible explanation for the discrepancies between measured and calculated trends in the single phase natural circulation, is the occurrence of reverse flow in some U-tubes (flow from outlet to inlet plenum of the SGs). This creates additional resistances in the loop that cannot be accounted for in a nodalization with only one U-tube. Nevertheless a two-dimensional model for the SG plena appears necessary to simulate this kind of phenomenon.

Fig. 5 - Primary Fluid Temperatures along Triple Loop at Steady State No.1. Comparison between measured and calculated trends

Fig. 6 - Primary Fluid Temperatures along Single Loop at Steady State No.1. Comparison between measured and calculated trends

		$R5 - B$	EXP.	Units
Primary System				
Mass flow:	- Intact Loop	0.98	1.0	kg/s
	- Broken Loop	0.31	\cdot 3	kg/s
Pressure:	- Upper Plenum	14.0	14.0	MPa
Fluid temperatures:				
Vessel outlet	- Intact Loop	320	321	$^{\circ}$ C
	- Broken Loop	320	319	\circ C
Vessel inlet	- Intact Loop	300	300	\circ C
	- Broken Loop	296.	296	\circ C
Pressurizer:		336.	336	$^{\circ}$ C
Core Power:		0.183	0.183	HW
Liquid mass without pressurizer:		373.	c. 370	kg
Liquid volume without pressurizer:		0.524	0.522	$\binom{3}{m}$
Pressurizer water level:		2.52	c. 2.5	w
Secondary System				
Auxiliary Feedwater mass flow:				
	- Intact Loop			
	- Broken Loop	--	--	
Pressure:	- Steam Line	8.65	8.65	MP _{i4}
Temperatures:				
Steam Generator	- Intact Loop	25.	25	۰c
Inlet	- Broken Loop	25.	25.	\circ C
Steam Generator	- Intact Loop	300	300	\circ C
Outlet	- Broken Loop	300	300	۰C
Downcomer Water	- Intact Loop	8.58	c. 8.6	
Level:	- Broken Loop	8.9	c. 8.9	m. m

Tab. VIII - Test conditions at initiation test A2-77: comparison between measured and calculated data

Fig. 9 - Mass Flow Rate in Triple Loop C.L.

Fig. 10 - Mass Flow Rate in Single Loop C.L.

Fig. 11 - Comparison between calculated and measured trends for UP pressure (case B). Calculated points have been superimposed to the experimental ones considering the PS mass inventory

Fig. 12 - Friction coefficient in the core and in IL HL for tests A2-77 and A2-81 (initial/nominal conditions)

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(*) Unidirectional probes

Tab. IX - Comparison between significant parameters in the experiment and in code calculations

The observation of loop flowrates demonstrated that tuning on roughness has the same effect as tuning in localized pressure losses and appears to be more physically based owing to uncertainties in the value of roughness also caused by corrosion/erosion of pipes internal walls.

To better understand the results, in Fig. 12 the plotter of the wall friction coefficient (Colebrook formula) implemented in RELAP5/MOD2 code is reported as a function of Reynold number.

The working points in HL and inside core are also reported with reference to initial (measured) conditions in A2-81 and A2-77A. It should be noted that the working points when going from A2-81 to A2-77A move in a transition zone of the diagram where Reynold number is important in the evaluation of the friction coefficient. So a further explanation of the unsatisfactory results obtained in case A (e.g. apart from experimental uncertainties in evaluating roughness) could be the inadequancy of Colebrook model in the zone.

4.5 - Case D (Presence of incondensible)

The objective of case D was to study if incondensible gases possibly present in the loop at the beginning of the test $8/$ could have the same effect as an increase of localized pressure losses. To this aim a mass of gas was given as input in the code equivalent to the volume occupied by the gas at ambient conditions in the U-tubes. The gas was assumed to be localized in the top of U-tubes. Unfortunatly the code did not work when the incondensible gas option was requested.

4.6 - Case E (Effect of parallel U-tubes)

The objectives of case E, that is the introduction in the nodalization of two parallel U-tubes characterized by different heights (difference in heights is 0.25 m) are essentially two:

- to observe if the lowest U-tubes allows larger natural circulation flowrates when steam appears in the top of the highest U-tubes;

- to observe possible links in the oscillations of velocities and densities inside the U-tubes.

The first problem which occurred in setting up the nodalization was the choice of the criteria by which separate the U-tubes. At least three possibilities did exist:

a) height, but almost each of the U-tubes has a different height;

b) position with respect to HL connection to SG PS **(3D** effect);

c) possible differences in inlet pressure losses (no data available).

As already mentioned the choice a) was made considering two U-tubes having equal characteristics apart from the height. The case B input was modified.

Significant results are shown in Tab. IX and in App. D. The main outcoming of the analysis is that natural circulation is slightly increased (as expected) with respect to the reference case (case B), when void formation occurs in the top of highest U-tubes. On the other side no interaction (parallel tube oscillations) appears to exit between the two Utubes in both SGs.

Also in this case better conclusions can be achieved only if a 3D model is available for the inlet plenum of **SG,** and a more detailed characterization is available from the experiment (measurement of localized pressure losses of each tube, diameter etc.).

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5. RELEVANT RESULTS ACHIEVED IN THE FRAME OF THE A2-77A TEST ANALYSIS

The LOBI A2-77A experiment constituted the basis of more in depth analysis at University of Pisa toward three main directions:

- a) study of the instabilities during the two-phase natural circulation period /11/,/12/;
- b) evaluation of the user influence on the setting up of the modalization and in the evaluation of the comparison between measured and calculated trends **/13/;**
- c) evaluation of the possibility to extrapolate measured and calculated natural circulation scenarios to plant systems /14/.

The main outcomes of the studies at items a) and b) are reported in the sections 5.1 and 5.2, respectively.

5.1 - Characterization of instabilities

Oscillations in almost all measured signals were detected for several thousands seconds into the transient, starting roughly when Primary System (PS) mass inventory was equal to 64% of the initial nominal value (including pressurizer and upper head inventories). They lasted up to the time when "stable" reflux condensation occurred (primary loop-residual mass roughly equal to 55% of the initial value). The highest amplitude of the peaks was observed just at the beginning of the unstable period; 400 s in this phase are considered hereafter.

The trends of differential pressures, fluid densities, velocities and temperatures in few points of the loop are shown in Figs. 13-16.

It should be noted that measured values of fluid velocity cannot be related specifically to either of the two phases being the output of a full-flow turbine meter. So the local void faction values should be considered in interpreting the signals. Even the comparison of liquid and steam velocities with calculated terms should be made with caution.

Different interpretations of the roots of the observed oscillations are possible when one examines the experimental data:

- **1.** Formation of relatively large steam slugs in the horizontal part of cold legs
- 2. Periodic clearing of loop seal
- 3. Formation of liquid level in the ascending legs of U-Tubes owing to entrainment of droplets and flooding at the inlet (Draining occurs when the liquid level reaches the top of the U-Tube)
- 4. Same scenario as in case 3, but with formation of liquid level due to condensation in U-Tubes
- 5. Combination of the above phenomena.

In cases 3 and 4, draining of tubes could be due to both a sort of syphon effect and the rise of the flooding and/or CCFL point from the bottom of U-Tubes toward the top.

The frequency and magnitude of differential pressure oscillation differed from tube to tube; they also took place out of phase. The interaction between the broken loop and intact loop, whose main variables oscillate out of phase, also can lead to energy exchange between the two steam generators. The relatively small pressure drop measured between the inlet and outlet of U-Tubes, with the possible presence of liquid in both
rising and/or descending legs of U-Tubes, does not permit the rising and/or descending legs of U-Tubes, does not permit the identification of the flow direction.

Still, the relatively low temperature difference between primary and secondary sides (a few kelvins) leads to a quite large range of consistent values of the condensation heat transfer across the U-Tubes.

Finally, the eventual fill and dump of U-Tubes, or at least the fluid

Fig. 13 - Differential pressure in intact-loop steam generator U-tube ascending leg (three U-tubes)

Fig. 15 - Fluid velocity at the in let of intact-loop steam generator and pump

Fig. 14 - Fluid density (diametral) in intact-loop and brokenloop hot legs

Fig. 16 - Fluid temperature in intact-loop and brokenloop cold legs

velocity oscillation in various zones of primary loop, creates variations in the heat transfer rate across steam generators, giving rise to temperature (and velocity) oscillations in the secondary side. This can cause feedback on the oscillations of primary side variables.

As far as the accuracy of the measurements is concerned, a very detailed procedure is adopted by the LOBI team to arrive at the definition of uncertainty bands for any measured signal $/4/$; these include consistency checks and calibrations before and after the test run. Reported values of accuracy should therefore be retained state-of-the-art values at least as far as integral facilities are concerned. Exemplary values are \pm 0.07 m/s, ± 15 Kg/m ³ , ± 10 KPa, and ± **1** K, with reference to velocity, densities, pressure drops and fluid temperatures, respectively.

5.1.1 - Code-predicted scenario

The application of the tuned nodalization (case B) as already mentioned produced a good agreement with the experimental data trends particularly with respect to the prediction of period and amplitude of oscillations. Comparison of measured and calculated data of differential pressure between the hot and cold legs of the intact loop and between the inlet and outlet of the U-Tubes in the broken-loop steam generator are presented in Figs. 17 and 18, respectively. Comparison of measured and calculated densities in the cold leg and velocities at the inlet of steam generator in broken loop are shown in Fig. 19 and 20, respectively.

The following main aspects can be outlined from the analysis of the corresponding couple of data:

- The mean values of the reported variables are essentially the same for both the experiment and the calculation.
- Liquid downflow appears in the calculation at U-Tube inlets (Fig. 20), whereas only steam flow can be detected from the experiment.
- The period of measured oscillations ranges around 130 s, while the calculated value is less than **100** s.
- Phase opposition occurs in the experiment (Fig. 14) between similar variables in the two loops, whereas essentially no phase shift is calculated by the code.
- Phase opposition still occurs in the experiment between hot and cold legs, whereas lower values of phase shift are calculated by the code **/11/.**

The analysis of predicted trends of void fractions inside the U-Tubes and of heat transfer between primary and secondary sides of steam generators permitted us to conclude that condensation in the ascending legs is the driving force of oscillations. In particular, the scenario depicted in Fig. 21 is the outcome of the calculation. At $t = t_1$ the steam-liquid mixture enters the ascending legs of U-tubes; condensation occurs on the walls owing to the lower temperature of the secondary system. This creates a rising mixture level in the ascending leg; in the descending leg the same phenomenon may occur owing to steam passing from the top of U-tubes and coming from the outlet plenum of the steam generator. Flooding, which occurs essentially at the inlet of the ascending leg, prevents the draining of the condensed liquid into the inlet plenum $(t = t₂)$. At $t = t₃$ the mixture level reaches the top of U-tubes, and the liquid flow begins toward the descending leg, the loop seal, and again to the vessel (siphon effect). The draining of the ascending leg leads to a new cycle. It should be pointed out that the scenario in Fig. 21 represents only a rough estimation of the reality owing to the one-dimensional nodalization of the U-Tubes (only one equivalent U-Tube is nodalized).

Fig. 17 - Comparison of measured and calculated differential pressures between hot and cold legs of intact loop

Fig. **¹⁹** Comparison of measured and Fig. 20 calculated densities in broken-loop cold leg

Fig. 18 - Comparison of measured and calculated differential pressures between inlet and outlet of steam generator broken loop

Fig. 20 - Comparison of measured and calculated velocities at the broken-loop steam generator inlet

Fig. 21 - Scenario of siphon condensation foreseeable from code calculation

the code calculation coincides with void formation in the horizontal part of the cold legs between the pump and the main vessel. In particular, the periodic collapse of the slug in this zone further enhances the oscillatory behaviour of the loop.

5.1.2 - Results from the phenomenological study

Owing to limitations of the code model and of the experimental data base, two main problems arise from the above analysis. These copncern

- Identification of the roots or at least of the parameters affecting the oscillations and evaluation of the realism of the proposed scenario
- Evaluation of the possibility to extrapolate the above phenomena to real plant situations.

The two problems are dealt with hereafter by applying independent models and performing sensitivity calculations with the code.

Behaviour of Parallel Tubes in the Case of Flooding

The behaviour of parallel U-Tubes in case of flooding has generally been experimentally investigated by various authors in the past with airwater mixtures.

With reference to the typical curve, giving the channel pressure drop as a function of the steam superficial velocity, four configurations that can occur simultaneously were measured by Wallis et al. */15/* (Fig. 22). It should be noted that situation **D** corresponds to pure gas flow in the experiment performed by Wallis et al., while condensation can be considered in the present analysis. The experiment was carried out with straight tubes.

This simple experiment shows that different configurations may occur simultaneously for different parallel tubes having the same imposed pressure drop when countercurrent flows of liquid and steam are involved.

Furthermore it should be emphasized that parallel U-tubes have even a larger degree of freedom than straight tubes, thus allowing, potentially, a larger number of simultaneous configurations with an imposed pressure difference between inlet and outlet.

CCFL Occurrence in LOBI Steam Generator Broken-Loop U-Tubes

The Hawighorst et al. **1161** correlation was applied to evaluate the CCFL occurrence in the steam generator broken-loop U-Tubes of LOBI facility. The correlation is:

$$
(J_{\mathbf{z},\mathbf{CR}}^4 \rho_{\mathbf{z}}^2)^{0.125} = C[\sigma g(\rho_f - \rho_{\mathbf{z}})]^{0.125}
$$
 (1)

The coefficient C distinguishes this correlation from the Kutateladzetype correlations.

Assuming the measured velocity and pressure in the hot leg of intact loop and the geometry of the system as boundary conditions, Eq. **(1)** demonstrates that two to four tubes (out of eight) should be in stalled conditions (case A in Fig. 22) at a given time, depending upon the chosen value for the velocity in hot leg $(0.2 \text{ to } 0.6 \text{ m/s})$. In particular, the measured value of gas flow rate corresponds to a Jg value lower than the value at the CCFL point if one considers all the U-Tubes. In the remaining U-Tubes, situations B, C or D can occur, allowing a rise in level in the ascending part.

Fig. 22 - Behavior of parallel straight tubes

Fig. 23 - Sensitivity calculations related to LOBI: trends of heat transfer across U-tubes when core power is varied

Fig. 24 - Sensitivity calculations related to LOBI: trends of velocities at the inlet of intact-loop steam generator U-tubes when the L/D ratio is varied

Condensation in U-Tubes of LOBI Steam Generator Broken Loop

The integral mass and energy balance of the fluid volume inside the ascending leg of U-Tubes, taking for simplicity the situation D in Fig. 22 as reference, yields for the time necessary to fill up a tube

$$
\bar{t} = \frac{Dh_{fg}\rho_f}{4H_{\text{cond}}DT} \tag{2}
$$

In order to have a rough estimation, the Nusselt correlation for the condensation heat transfer coefficient and the measured values for the requested quantities have been assumed. In this situation, t varies between 100 and 250 s depending upon the value of DT (ranging between 8 and 3 K).

The range of variation of t is consistent with the period of oscillations, demonstrating that condensation is a possible mechanism of instability.

CCFL Breakdown in U-Tubes of LOBI Steam Generator Broken Loop

It appears worthwhile to investigate the possibility of steam entering the U-Tubes to sustain a liquid column in the ascending leg. Looking at the previously considered time into transient, a pressure difference across the ascending legs of the broken-loop U-Tubes of about 0.03 ± 0.01 MPa is obtained from the experimental data/4/. This means an equivalent collapsed level ranging between 6 m and 3 m.

As a consequence of the above result, in order to have liquid at the top of U-Tubes, presence of steam bubbles and/or movement of the location of CCFL from the U-Tube inlet must be assumed.

Specific Sensitivity Calculations

Sensitivity calculations were performed by RELAP5/MOD2 code. In principle, several parameters affect the characteristics of the oscillations. Some of these are inherent to the structure of the code (nodalization, time step, empirical models, etc.), others depend upon the manner in which boundary conditions are fixed (e.g., pressure control system in secondary side, draining modes), and others actually depend upon the system configuration (distribution of heat losses and of form loss coefficients, geometry, power, etc). Consideration of the entire set of parameters in sensitivity calculations requires a very extended work; attention is focused hereafter on the last class inherent to system configuration.

Sensitivity analyses were performed with reference to roughness, local loss coefficients, the length/diameter ratio (L/D) of both hot and cold legs, and core power. The last two parameters were found to be the most significant, and the related results are outlined here. The values assumed by the above quantities are reported in Table X and compared with typical values of PWR plants and of LSTF facility.

It should be noted that in sensitivity calculations **El** and E2 the length of horizontal parts of cold leg in the LOBI facility were varied by the same amount as those for the hot legs, maintaining the original diameter.

Significant results are shown in Figs. 23 and 24, where calculated system variables are plotted for the considered time span at constant value of primary-loop mass inventory. Comparison with the nominal case leads to the following observations:

*Test ST-NC-002

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Tab. X - Values Assumed by Some Parameters in Sensitivity Calculations performed in order to characterize the oscillations

- The amplitude of oscillations decreases while the frequency increases when core power is reduced.
- The reduction of the L/D ratio creates new frequencies for the oscillations that are superimposed to the original ones, and an increase in L/D leads to an increase in amplitude.

The above analysis shows a close link between the characteristics of the siphon condensation phenomenon and the system geometry.

Consideration of U-Tube Height

As already mentioned (case E in sect. 4) the difference in heights has been considered in the literature as a possible reason for differences in

the behaviour of U-Tubes that potentially contributes to the oscillations.
Taking this into account, parallel pipes were nodalized wit Taking this into account, parallel pipes were nodalized with RELAP5/MOD2, splitting into two egual parts (with reference to flow area) the original pipe simulating the U-Tubes. The new pipes are characterized by different heights and lengths.

The resulting void fractions at the top of the inverted U-bend of the intact-loop steam generator are compared in Fig. 25 with the mean trend of the nominal calculation. No appreciable differences can be seen from the above trends or from a more detailed comparison. On the contrary, arbitrary variations of inlet loss coefficient cause substantial differences in the behaviour of the two U-Tubes.

These results demonstrate that the splitting of U-Tubes into parallel pipes in a nodalization suitable for a one-dimensional code is meaningful only if a two-dimensional model is available for the inlet plenum. This would allow differences in behavior among the U-tubes caused by fluid dynamics inside the plenum itself. Without the availability of such a model, differences in the behaviour of U-Tubes can be caused by different heights (this has been shown to be irrelevant), by different values of the inlet/outlet form loss coefficient (to be fixed almost arbitrarily), or by different number of U-Tubes in each group (to be fixed arbitrarily).

Presumed Scenario

The above results and considerations made it possible to define a more complex scenario for the ensemble of U-Tubes than that reported in Fig. 21. In particular, the following additional items are taken into account:

- Groups of tubes may exhibit substantially different behaviour with fixed boundary conditions: some tubes must be in stalled conditions in order to have CCFL in others.
- The height of tubes has little influence on the global evolution of the phenomenon; the two-dimentional fluid dynamics inside inlet plenum mainly differentiate the U-Tubes behaviour;

- The CCFL front is likely to advance.

The scenario in Fig. 26 was contrived with these considerations in mind.

5.2 - Influence of the code user on the results

The post-test analysis of the A2-77A experiment has been carried out independently, in the frame of an European Community research Program, by six different users utilizing five advanced codes. An outline of codes and users is given in Tab. XI /13/. A detailed descriptions of the codes can be found in ref. **117/.**

Fig. 25 **-** Calculated trends of void fraction at the top of U-tubes using a single pipe (nominal case) and two parallel pipes to nodalize the U-tubes

Fig. 26 - Presumed scenario for explaining siphon condensation. At each time t_1, t_2 , and t_3 , U-tubes are in configurations (a), (b), and (c), respectively.

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Tab. XI - Involved users and codes

5.2.1 - Comparison among input parameters

The lack of reflection on the input parameters may result as the most significant limitation of a current code analysis: the first specific objective pursued hereafter is to consider this gap, for a specific thermalhydraulic problem, by means of the comparison of the input decks developed by several users. The existence of different approaches pursued by the involved users, dealing with a wide variety of choices requiring subjective judgements, will be demostrated.

The starting point in setting up the nodalization is the knowledge of the following subjects:

- **1)** engineering of the facility and its instrumentation;
- 2) test specifications;
- 3) experiment scenario when dealing with post-test calculation as in the case here considered.

The final product is a "code model", that is a compromise between the user knowledge of the code performance, of the facility hardware and of the simulated transient scenario; the code cabilities and the required CPU time also play a role in defining the specifications of the input model.

Taking into account of the above, in order to realize a critical comparison of the nodalizations, twenty items related to the input decks (Tab. XII, Figs. 27 to 30) are compared among each other. These have been split into three groups:

- **-** the elements characterizing the degree of detail of the code model (e.g. number of nodes etc.; items 1 to 5);
- the elements characterizing the "nodalization fidelity" to the geometrical data of the facility (e.g. value of the overall volume of the facility; items 6 to **11);**
- the elements characterizing the interface between hydraulics and geometry (essentially pressure drop coefficients; items 12 to 20).

From the comparison of the above data, the following considerations can be made:

1) elements characterizing the nodalization detail.

In relation to the number of nodes, on one hand RFIAP and ATILET codes, with roughly 150 nodes, on the other hand the CATHARE code, with more than 300 nodes, and the TRAC code, with 250 nodes, can be distinguished.

This choice is only partially due to the user; instead the code numerical structure plays an important role for establishing the degree of detail of the code model. As an example, RELAP code, owing to the Courant limit, needs nodes having length greater than few tens of centimeters for making possible the simulation of typical LOCA transients, while CATHARE code has not such a constraint, allowing a greater freedom. Still, a geometrical discontinuity cannot be modelled as it is by the CATHARE code, but needs several components (nodes) of small dimension in the flow direction which are characterized by different areas.

In principle, the best results for a physical simulation should be given by a nodalization with a number of nodes ad large as possible, but this idea is abruptly nullified by the majority of the current system codes. Otherwise, an optimal number of nodes can be recognized for each code for a given simulation problem. Directions for the attainment of this number are not available in any code manual: only the user experience can achieve this parameter, considering the phenomena to be analyzed, in line with the available resources (CPU, computers, etc.) and the goals of the study, e.g. sensitivity analyses aiming at the interpretation of physical phenomena, licensing calculations, etc. It should be noted that a large

Tab. XII - Items related to the input decks

Fig. 27 - Number of nodes chosen by various users for nodalizing the LOBI facility

Fig. 28 - Overall number of structures mesh points chosen by various users for nodalizing the LOBI facility

Fig. 29 **-** Mass inventory in broken loop steam generator secondary side resulting from steady-state calculation by various users and comparison with experimental data

Fig. 30 **-** Forward pressure drop coefficient at intact loop connection with pressure vessel as selected by various users

amount of sensitivity analyses can bring substantial improvements of nodalization parameters; in this way coarse nodalization, with few elements, can produce better results than fine nodalization with much larger number of elements. This is valid for items **I** to 5.

Concerning the overall number of structures mesh points (item 5), the capital influence of this parameter on the heat transfer mechanism must be stressed. In particular, the heat release from structures is strongly affected by the number of meshes. The experience on system code assessment demonstrate that mesh thickness of few millimeters for piping and flanges, few tenths of millimeters for fuel rods, should be used for a suitable schematization.

2) "nodalization fidelity" to the geometrical data

The geometrical fidelity of the nodalization to the system hardware is a code model peculiarity that should be considered very carefully. As a difference from the nodalization detail, that is essentially an user choice, the agreement betwen input and actual system related parameters is an objective goal to be pursued. With reference to the overall volume of the facility, that can be considered the most important parameter in this group, the following approximations or inadequacies affect the agreement between actual system value and code input value

- imperfect knowledge of the plant data (drawings inadequacies, - imperfect knowledge of the plant data (drawings consideration of thermal dilatation of structures, etc.);
- presence of dead ends (nozzle, instrumentations lines, etc.);
- need to simulate a three dimensional configuration with a one-dimensional one.

The experience demonstrates, for example, the acceptability of few percent error on the nominal value of the overall facility volume.

Another aspect of the nodalization fidelity is related to the active heat transfer areas, i.e. sources and sinks. In relation to this, the most important characteristics to be preserved are the total core area and the overall steam generator heat transfer surface. Usually, these parameters are respected with an error less than 1%.

Two further problems, typically encountered by system code user, arise during the development of a nodalization:

- a) with reference to a PWR typical plant, the choice of the hydraulic channels number in the steam generator and in the core (for a BWR plant the same problem may occur in relation to the number of jet pumps and still to the number of core channels $\{8\}$. A schematization with only few "pipes" can preserve the overall thermal energy balance, but is blind towards a nonuniform behaviour of the various channels (e.g. nonuniform flow distribution in the steam generator U-tubes, or in the lower plenum of reactor vessel or in the steam generator plena, channel oscillations, etc.);
- b) passive structures of the plant. The consideration of all the structures constituting vessel, piping and internal wall, as well as flanges, valves and pump casing is almost impossible owing to limitations of computer memory. Approximations are needed and are usually done. Errors of the order of few percent and less than 20% of the nominal values of passive heat transfer areas and volumes, respectively, appear to be acceptable.

3) interface between hydraulics and geometry of the code model

This topic refers, essentially, to local and distributed pressure drops coefficients. With reference to the friction factor, a substantial difference among the various basic models implemented in the codes can be observed, e.g. the dependency of friction factor upon the Reynolds number

is not considered in the same way in all codes: for example Colebrook correlation is used in RELAPS/MOD2 and peculiar correlations are adopted by TRAC and CATHARE. Still roughness is only accounted by the RELAP input deck, while it has no influence on calculated pressure drops in CATHARE and TRAC.

Finally, with regard to the form loss coefficients, items 12 to 18 in Tab. XII the following remarks can be made:

- a) there is no theoretical model siutable to calculate this parameter in the wide variety of configurations encountered in modelling a typical nuclear plant or simulator; still no relaptionship gives the dependency of these factors upon Reynolds number and local void fraction;
- b) experimental uncertainties are often connected to this parameter, usually derived from pressure drop measurements;
- c) loss coefficients values have to account for the three-dimensional effects that cannot be modelled by a one-dimensional code;
- d) in definitive, these factors are a sort of "magic numbers", that, in order to have a good code prediction, must counterfeit many of the thermalhydraulic model deficiencies.

The data reported in Tab. XII and in Fig. 31 to 34 give an idea of typical dispersion range associated with assumed values of these parameters. It should be noted that "all" the above values lead to "reasonable agreement" with experimental data at least as far as the initial steady state in concerned (next section).

5.2.2 - Significant results from the calculations

The first natural circulation steady state of the test is executed at nominal mass inventory (100% of Residual Mass excluding upper head). During this period, a subcooled single-phase natural circulation flow takes place between core and steam generator.

This presents a particular interest because:

- a) it constitutes the first way to test the computer codes as well as the nodalizations during natural circulation flow;
- b) this is the situation that comes out in the primary side of a PWR after scram and main circulating pumps trip, without loss of coolant.

Relevant results from steady state calculations by the various codes are summarized in Tab. XIII and in Figs. 35 to 38. At least two inherent limitations of the available experimental data base should be pointed out:

- a) core power is 183 kW, while heat losses to environment from primary loop are of the order of 60 kW, with at least **10** kW uncertainty;
- b) the operation of feedwater and steam line in the secondary side is not entirely specified: involved flowrates are sufficient to maintain constant the pressure in the secondary side and to remove the residual power coming from primary loop, but the recirculation flowrates and the temperature distribution in the riser of the secondary side steam generator are not fully specificed.

The above two facts, together with the different possibilities of assigning the initial temperature distribution inside the wall structures in the zones where heat losses are concentrated, lead to indetermination in the calculation of loop energy balance. As an example, **I** K of uncertainty in secondary side mean temperature in the subcooled riser region may correspond to about 10 kW of indetermination in the power exchange between primary and secondary side.

Still, a variation of about 10% around the nominal experimental value of loop flowrate also corresponds to variation of few tens kW in the heat transfer across steam generators.

Finally, the pressure differences over the steam generators (less than

 $\frac{1}{2}$

Fig. 31 - Pressure drop over steam generator intact loop primary side resulting from steady-state calculation by various users and comparison with experimental data

Fig. 32 **-** Pressure drop over steam generator broken loop primary side resulting from steady-state calculation by various users and comparison with experimental data

Fig. 33 **-** Mass flowrates in downcomer of steam generator intact loop resulting from steady-state calculation by various users $\lambda=1.4$

Fig. 34 - Power losses to environment considered by various users and comparison with experimental data

ITEM	TRAC/PF1	CAT1 V1.3	RS/M2	ATHLET	R5/M1-EUR	RS/M2	EXP	
21. destabution of fluid temperature in H.	CEA	CENG	DCMN	GRS	JRC	UKAEA		notes
	Fig. 12	Fig. 12	Fig. 12	$\overline{\text{Fig. 12}}$	Fig. 12	Fig. 12	Fig. 12	enclosed
22. distribution of floid temperature in RIN.	Fig. 13	Fig. 13	Fig. 13	Fig. 13	Fig. 13	Fig. 13	Fig. 13	enclosed
23. distribution of fluid temperature in IT SG. ŧП	Fig. 14	Fig. 14	Fig. 14	Fig.14	Fig. 14	Fig. 14	Fig. 14	enclosed
24. distribution of fluid temperature H. SC. 55. RE	Fig. 15	Fig. 15	Fig. 15	Fig. 15	Fig. 15	Fig. 15	Fig. 15	cnclosed
25 Directed (1994).	-0.45	Ω	-0.70	-0.7	-0.928	-1	0.12	
26 Diesembt (O.Ba)	-0.40	Ω	-0.70	-0.6	-0.988	-1.05	-0.8 ± 0.8	
27. Director (2011) (O.Pa).	-1.3	Ω	-0.50	-0.7	-0.400	-0.9	-1.0 ± 1.0	
28. Director SG BE (FEE).	-1.3	Ω	-0.40	-0.6	-0.400	-0.8	-0.75 ± 1.0	
29 Do over Il PU (194)	-5.9	-6.7	-5.8	-8.0	-6.204	-8.02	-8.0 ± 2.0	
30 Do over 1st PU (kPa)	-4.9	-2.9	-5.7	-6.0	$-5.757/-1.80$	-5.08	-2.7 ± 5.0	
31. mass flowrate in B. Cl., (kp/x).	$\mathbf{1}$	1.03	0.98	1.1	1.02	1.1°	1.0	
32 mass flowrate in HE (Fp/s).	0.3	0.33	0.31	0.38	0.32	0.36	0.3	
33 pressure in UP (MPa)	14.0	14.3	14.0	14.0	14.03	14.0	14.0	
M. temperature in PRZ (K)		609	609	609	609.8	609	609	
35 level (collapsed) in PRZ (m)	\blacksquare	2.7	2.52	2.5	4.9	1.9	2.5	
36 pressure in SCs SS (MPa).	8.66/8.62 \bullet	8.65	8.65	8.65	8.66	8.65	8.65	۰ пли.
37. EW temperature at IL, & TIL. (K)	298	298	298	298	298	303.15	298	
36. temperature at IL & BL SG outlet. (K)	574	$563/571$ °	573	573	573.75	573.62	573	$\overline{6}$ II'll $\overline{6}$
39. DC collapsed fevel in SG IL. (in).	8.19	8.6	8.58	8.6	6.44	8.7	8.6	
40 TX collapsed fevel in SG III. (m)	8.25	8.9	8.90	8.9	6.56	8.9	8.9	
41. mass flowrate in LK' of SG II. (kg/s).	1.18	0.56	2.35	3.9	0.0234	2.6	NΛ	
42 mass flowrate in DC of SG BE (kp/k).	0.27	0.52	0.52	1.0	-0.11	0.8	NA	
43. power losses IS ensitonment. (EW).	35.6	40	55	45	30.9	48.4	40	exit PU and 0.0167
44 power losses SG II, environment. (kW).	8.9	22	14	12.4	16.83	24	$29 + 5$	
45. power losses SG BL environment. (kW).	10.4	20	12	14.3	12.1	18	$21 + 5$	
46 power exchange across SG II, (kW).	117.8	96	105	105	72.9	NA	NA	
47. juillet exchange actoss SG III, (kW).	32 ₂	47	31	32	22.3	NA	NA.	
48. core power (EW).	183	183	183	183	178.5	183	183	
49. UP DC bypain flow (kp/n).	\bullet	0.013	0.023	0.039	0.018	0.064		
50. HEATH INSTANTION (Ep/s)	\blacksquare	0.025	0.013	0.039	0.037	0.032	\bullet	
51 core hypass. (kg/s).	$\qquad \qquad \blacksquare$	0	\bullet .	0.052	$\mathbf 0$	\blacksquare	\blacksquare	

Tab. XIII - Items related to the initial steady state

Fig. 35 - Distribution of fluid temperature along intact loop resulting from steady-state calculation by various users and comparison with experimental data

Fig. 36 - Distribution of fluid temperature inside reactor pressure vessel resulting from steady-state calculation by various users and comparison with experimental data

Fig. 37 - Distribution of fluid temperature along U-tubes of steam generator intact loop resulting from steady-state calculation by various users and comparison with experimental data

Fig. 38 - Distribution of fluid temperature in intact loop steam generator secondary side riser resulting from steady-state calculation by various users and comparison with experimental data

zero, that is with a pressure gain, in this case; items 27 and 28 and Fig. 31) have an error band of the order of 200% of the measured value. This uncertainty is very critical because inherent to a quantity strongly affecting the U-tubes behaviour.

The above considerations partially explain the differences between initial values of parameters shown in Tab. XIII and in Fig. 27 to 38. In definitive, each user interprets in a subjective way the uncertainties characterizing the experimental data; then the importance of the user quantification is once again highlighted.

An exaustive report of the comparison between measured and predicted results of each calculation is not the objective of the present document. In the attempt to be more complete and, partly, to give an idea of the consequences of the choices discussed in the previous section, the calculated flowrates during the whole transient are compared with the nmeasured values in Fig. 39 and 40 as a function of the overall residual mass of the primary side. Reference is made to the intact and the broken loop, respectively.

The qualitative evaluation of the transient is well predicted by each code, that is the various codes capture the basic physics of the involved phenomena, but the quantitative values of the physical variables are nor satisfactory. The codes limitations already discussed before also apply in this case.

Fig. 39 - Comparison between measured and calculated trends of steadystate mass flowrate in intact loop during the whole transient as a function of primary side mass inventory

Fig. 40 - Comparison between measured and calculated trends of steadystate mass flowrate in broken loop during the whole transient as a function of primary side mass inventory

6. CONCLUSIONS

The appplication of RELAPS/MOD2 to post-test analysis, of the natural circulation test A2-77A performed in LOBI/MOD2 facility, showed the good quantitative capability of the code in predicting a transient characterized by low fluid velocities. The "tuning" of some input conditions also allowed the achievement of good quantitative agreement.

The analysis led to the identification of deep differences in the interpretation of an assigned experimental data base, from various users.

The evaluation of the relationship between user choice and code predicted test scenario was not the main purpose of the present paper; nevertheless, the study allowed to emphasize the role of user and raised some concern in relation to the validity of large system code assessment activities performed all over the world.

The discrepancies between experimental data and code predictions are due both to the intrinsec code limit and to the user behaviour. There is a worthful need to quantify the percentage of disagreement due to the poor utilization of the code and that due to the code itself.

The quality of the input data base (nodalization) should be consistent with the quality of the prediction: i.e. a good prediction obtained by the use of a unrealistic input deck should be seen as unphysical. A check in this connection is always necessary from the code user; in particular, exhaustive and objective information related to the approach pursued in setting up the nodalization and possibly in its "tuning" should be supplied in any calculation report.

Unfortunately the emphasized uncertainties in the boundary conditions prevent the possibility to definetely evidence code limitations. An example of this is given by the influence of roughness on the results (sect. 4.4).

Taking into account of the above, the following thermalhydraulic models are presumed to be mostly responsible of the disagreement between measured and calculated trends:

- one-dimensional approach in modelling some components (e.g. steam generator plena);
- flooding (especially during the "reflux condensation" period);
- stratification in horizontal pipes: the inadequacy of this model, with main reference to the prediction of piping related phenomena, is the possible cause of the flowrates overestimation shown during the transient by most codes;

- dependence of local pressure drop upon Reynolds number and void fraction. As an additional result of the study, an explanation of the

phenomenology at the basis of the oscillations measured in the LOBI PWR simulator during natural circulation has been proposed. The origin of the instabilities is a sort of siphon condensation controlled by flooding that occurs in the ascending legs of U-tubes. The energy transfer between core and steam generators constitutes the driving force for the instability mechanism.

The application of basis thermalhydraulic models demonstrated that different configurations presumably occur at the same time for the different U-tubes. Some of these are in stalled conditions as far as inlet steam flow is concerned, others are in the flooding process, the CCFL occurs in a third group ot U-tubes (Fig. 26).

The application of RELAPS/MOD2 code made it possible to evaluate the influence on the phenomena of boundary and initial conditions. In particular, it was found that different parameters affect the frequency and amplitude of oscillations; the length/diameter ratios of the hot legs and core power are significant ones. The occurrence of voiding in the horizontal parts of cold legs at the onset of instability also resulted

from calculations and is in agreement with the experimental observation. This conclusion should be obvious if one considers the limitations of the current generation codes not discussed in the present paper.

The phenomenon is typical of U-tube geometry, so oscillations are foreseeable for nuclear plant situations. A code calculation performed with reference to the **DOEL** PWR supported this conclusion. Interaction between neutronics and thermal hydraulics should be evaluated in different plant configurations.

Finally the consideration of two parallel U-tubes in SG PS does not appear convenient due to the basic **I-D** model implemented in the code (case E) and owing to uncertainties in fixing the boundaries, which overshadow the differences between calculations utilizing one or two U-tubes.

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ABBREVIATIONS

SYMBOLS

Greek Symbols

Subscripts

 $\hat{\mathcal{L}}$

 $\sim 10^7$

 \bar{z}

APPENDIX A

Measured and calculated trend for case A (Initial calculation)

UPPER PLENUM PRESSURE

FIGURE A2 MASS FLOW RATE IN DOWNCOMER VS PS HASS INVENTORY

FIGURE A6 FLUID VELOCITIES IN I.L. PUMP INLET

APPENDIX B

Measured and calculated trends

for case B

(Reference calculation)

<u>.</u>

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FIGURE B4 **UPPER PLENUM PRESS RE**

 $\mathbb{R}^{\times}\otimes_{\mathbb{Q}}\mathbb{R}^{\times}\otimes\mathbb{B}6$ B.L. HOT LEG PRESSURE

 \longrightarrow \rightarrow

CUFFERENTIAL PRESSURE PD3017 **FIDURE B8**

 $F32RE$ B10 DIFFERENTIAL PRESSURE PD8227

FIGURE **B12** DIFFERENTIAL PRESSURE PD3DBT

 $F3URE$ B14 DIFFERENTIAL PRESSURE PD3RKA

FIGURE B15 DIFFERENTIAL PRESSURE PD9CBN

FIGURE B16 DIFFERENTIAL PRESSURE PDSCBN

 \sim

 $F16.7E$ B18 DIFFERENTIAL PRESSURE PD163DB3

 ~ 10

FIGURE B20 DIFFERENTIAL PRESSURE PD3R11A4

 $F(3,5)$ DIFFERENTIAL PRESSURE RPV TOTAL IL SIDE

 \sim

FLUID VELOCITIES IN I.L. PUMP INLET F) $\ensuremath{\mathsf{G}}\xspace$, RE **B24**

FIDURE B26 FLUID VELOCITIES IN BLL. PUMP INLET

 $\ddot{}$

FIGURE B28 FLOWRATE DRAINED FROM LP

FIGURE **B30** FLUID DENSITY IN ILL-HILL VESSEL CUTLET

 $F332$ FLUID DENSITY IN ILL-CLL, VESSEL INLET

 $\mathcal{L}^{\text{max}}_{\text{max}}$

 \mathbb{R}^2

FIGURE B33 FLUID DENSITY IN B.L.-H.L. VESSEL OUTLET

FIGURE B34 FLUID DENSITY IN B.L.-PUMP INLET

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 $\hspace{0.1mm}-\hspace{0.1mm}$

FIGURE B38 FLUID TEMPERATURE IN T.L.-S.G. INLET

FIGURE B40 FLUID TEMPERATURE IN I.L.-PUMP INLET

FIGURE B42 FLUID TEMPERATURE IN BILL-HILL VESSEL OUTLET

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FISURE B44 FLUID TEMPERATURE IN B.L.-S.G. OUTLET

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FIGURE B46 FLUID TEMPERATURE IN B.L.-PUMP OUTLET -

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 $\ddot{}$

 $\sigma_{\rm{eff}}$ FIRURE **B48** FLUID TEMPERATURE IN D.C. MIDDLE SECTION

FIGURE B50 FLUID TEMPERATURE IN L.P.

 $\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)$ FIGURE B52 SURFACE ROD TEMPERATURE $\sim 10^{-11}$ χ^2

FIGURE B54 SURFACE ROD TEMPERATURE \mathcal{A}

 \sim \sim \pm $\mathcal{L}_{\mathcal{L}}$ SURFACE ROD TEMPERATURE FIGURE B56

FEUID TEMPERATURES IN UPPER PLENUM FIGURE B58

FIGURE B59 FEUID TEMPERATURES IN ILL. S.G. SECONDARY SIDE TOP

FIGURE B60 FLUID TEMPERATURES IN ILL. S.G. SECONDARY SIDE COWNOOMER T

 \sim \sim

 \mathcal{A}

FIGURE B62 FLUID TEMPERATURES IN E.C. S.O. SECONDARY SIDE RISER

FIGURE B63 FLUID TEMPERATURES IN I.L. S.G. SECONDARY SIDE U-TUBES GUTLE

98.

 $\ddot{}$

FIGURE B66 DIFFERENTIAL PRESSURE IN I.L. UT ASCENDING LEG

FIGURE B67 DIFFERENTIAL PRESSURE IN I.L. UT DESCENDING LEG

FIGURE B68 FLUID VELOCITIES IN TRIPLE LOOP H.L.

FIGURE B70 FLUID VELOCITIES IN TRIPLE LOOP C.L.

FIGURE B71 FLUID VELOCITIES IN SINGLE LOOP C.L.

FIGURE B73 FLUID DENSITY IN THEFFOLL.

FIGURE B74 FLUID DENSITY IN B.L. C.L. it.
D α , β

FIGURE B75 FLUID DENSITY IN B.L-H.L.

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FICUREB76 DIFFERENTIAL PRESSURE IN B.L. dT ASCENDING **LEG**

FIGJRE B77 DIFFERENTIAL PRESSURE IN B.L. jT DESCENDiNG LEG

105.

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 $\label{eq:2.1} \begin{split} \mathcal{L}_{\text{max}}(\mathbf{r}) & = \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \\ & = \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}$

 $\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$

APPENDIX C

 $\mathcal{L}(\mathcal{L})$ and $\mathcal{L}(\mathcal{L})$.

 $\mathcal{L}^{\text{max}}_{\text{max}}$, where $\mathcal{L}^{\text{max}}_{\text{max}}$

Measured and calculated trends

for case C

(Piping roughness changed)

 $\sim 10^{-11}$

107.

 $\mathcal{L}^{\text{max}}_{\text{max}}$

 $\label{eq:2.1} \frac{1}{2} \int_{\mathbb{R}^3} \left| \frac{d\mu}{d\mu} \right|^2 \, d\mu = \frac{1}{2} \int_{\mathbb{R}^3} \left| \frac{d\mu}{d\mu} \right|^2 \,$ $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ and $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ and $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ and $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\$ $\mathcal{L}^{\mathcal{L}}(x)$ and $\mathcal{L}^{\mathcal{L}}(x)$ are the set of the set of

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$

 $\mathcal{L}^{\text{max}}_{\text{max}}$

MASS FLOW RATE IN DOWNCOMER VS PS MASS INVENTORY **FIGURE** $C₂$

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FIGURE C3 FLUID VELOCITIES IN B.L.S.G. INLET

,FIGURE C4 FLUID VELOCITIES IN B.L. PUMP INLET

FIGURE C6 FLUID VELOCITIES IN I.L. PUMP INLET

APPENDIX D

Measured and calculated trends

for case E

(Effect of parallel U-Tubes)

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\text{max}}(\mathcal{L}^{\text{max}}_{\text{max}}(\mathcal{L}^{\text{max}}_{\text{max}}(\mathcal{L}^{\text{max}}_{\text{max}})))$

 $\label{eq:2.1} \mathcal{L}=\left\{ \begin{array}{ll} \mathcal{L}^{\frac{1}{2}}\left(\mathcal{L}^{\frac{1}{2}}\right) & \mathcal{L}^{\frac{1}{2}}\left(\mathcal{L}^{\frac{1}{2}}\right) & \mathcal{L}^{\frac{1}{2}}\left(\mathcal{L}^{\frac{1}{2}}\right) \\ \mathcal{L}^{\frac{1}{2}}\left(\mathcal{L}^{\frac{1}{2}}\right) & \mathcal{L}^{\frac{1}{2}}\left(\mathcal{L}^{\frac{1}{2}}\right) & \mathcal{L}^{\frac{1}{2}}\left(\math$

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$

MASS FLOW RATE IN SINGLE LOOP C.L. VS PS HASS INVENTORY FIGURE D₂

 $\ddot{}$

FIGURE D4 UPPER PLENUM PRESSURE

FIGURE D6 B.L. HOT LEG PRESSURE

 $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$

DIFFERENTIAL PRESSURE PD9217 FIGURE D8

 $D9$

DIFFERENTIAL PRESSURE PD1714

FIGURE D10 DIFFERENTIAL PRESSURE PDE227

 $\frac{1}{k}$

FIGURE D12 DIFFERENTIAL PRESSURE PD3DBT

FIGURE D18 DIFFERENTIAL PRESSURE PD163DB3

÷.

FIGURE **D22** DIFFERENTIAL PRESSURE RPV TOTAL IL SIDE

्री $\mathcal{F}_{\mathcal{A}}$, \mathcal{F} FLUID VELOCITIES IN I.L. PUMP INLET FIGURE D24 $\mathcal{O}(\mathcal{O}_\mathcal{O})$

 $\bar{\star}$

FIGURE $D26$ FLUID VELOCITIES IN B.L. PUMP INLET

 \vec{v}

 F I GURE D28 FLOWRATE DRAINED FROM LP .-

PRIMARY LOOP RESIDUAL MASS

FIGURE D30

 $\mathcal{O}(\frac{1}{2})$

 $\mathcal{F} \subseteq \mathbb{C}$. FIGURE D32 FLUID DENSITY IN I.L.-C.L. VESSEL INLET

 $\frac{1}{2}$

FLUID DENSITY IN B.L.-PUMP INLET FIGURE D34

D38 **FIGURE** FLUID TEMPERATURE IN I.L.-S.G. INLET

 $\alpha = \alpha$. FLUID TEMPERATURE IN I.L.-PUMP INLET: FIGURE D40

FIGURE D44 FLUID TEMPERATURE IN B.L.-S.G. OUTLET

FIGURE D46 FLUID TEMPERATURE IN B.L.-PUMP OUTLE*

FIGURE D48 FLUID TEMPERATURE IN D.C. MIDDLE SECTION AND

FLUID TEMPERATURE IN D.C. BOTTON

FIGURE **D52** SURFACE ROD TEMPERATURE $\sim 10^6$

 $\frac{1}{2}$. 4

FIGURE D53 SURFACE ROD TEMPERATURE

FIGURE D54 SURFACE ROD TEMPERATURE

SURFACE ROD TEMPERATURE

FLUID TEMPERATURES IN I.L. S.G. SECONDARY SIDE DOWNCOMER T. **FIGURE** D60

FIGURE D62 FLUID TEMPERATURES IN I.L. S.G. SECONDARY SIDE RISER

FIGURE]D64 FLUID DOWNCOMER LEVEL IN I.L.S.G.

 \bar{z}

 \mathbb{R}^2

 $\mathcal{F}_{\rm{int}}$

FIGURED66 DIFFERENTIAL PRESSURE IN I.L. UT ASCENDING LEG UPPER TUBE

FIGURE **D67** DIFFERENTIAL PRESSURE IN I.L. UT DESCENDING LEG UPPER TUBE

 $\ddot{}$

FIGURED68 DIFFERENTIAL PRESSURE IN I.L. UT ASCENDING LEG LOWER TUBE

FIGURE D69 DIFFERENTIAL PRESSURE IN I.L. UT DESCENDING LEG LOWER TUBE

 $\frac{1}{2}$

FIGURE D70 DIFFERENTIAL PRESSURE IN B.L. UT ASCENDING LEG UPPER TUBE

FIGURE D71 DIFFERENTIAL PRESSURE IN B.L. UT DESCENDING LEG UPPER TUBE

FIGURE **D72** DIFFERENTIAL PRESSURE IN B.L. UT ASCENDING LEG LOWER TUBE

FIGURE D73 DIFFERENTIAL PRESSURE IN B.L. UT DESCENDING **LEG** LOWER TUBE

FIGURE D74 FLUID VELOCITIES IN TRIPLE LOOP H.L.

PEIGURE D75 FLUID VELOCITIES IN TRIPLE LOOP C.L. $\sim 2.7 \times 10^{-1}$

FIGURE D77 FLUID VELOCITIES IN SINGLE LOOP H.L.

 $\bar{\alpha}$

 $\gamma_{\rm c} \sim 3 \gamma \sim 3$

FIGURE D78 FLUID DENSITY IN I.L.-H.L.

 $\beta = 1$

460

FIGURE D80 FLUID DENSITY IN B.L-H.L.

FIGURE D81 FLUID DENSITY IN B.L.C.L.

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ $\label{eq:2.1} \mathcal{L}_{\mathcal{A}}(\mathcal{A}) = \mathcal{L}_{\mathcal{A}}(\mathcal{A}) = \mathcal{L}_{\mathcal{A}}(\mathcal{A}) = \mathcal{L}_{\mathcal{A}}(\mathcal{A})$

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}))\leq \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}))\leq \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}))$

 $\label{eq:2.1} \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2\left(\frac{1}{\sqrt{2}}\right)^2.$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal$

 $\hat{\boldsymbol{\beta}}$

 $\hat{\boldsymbol{\theta}}$

 $\label{eq:2.1} \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma\delta}(\mathcal{H}^{\alpha\beta\gamma\delta}) = \mathcal{H}^{\alpha\beta\gamma$ $\label{eq:2} \mathcal{L}(\mathbf{z},t) = \mathcal{L}(\mathbf{z},t) + \mathcal{L}(\mathbf{z},t) + \mathcal{L}(\mathbf{z},t) + \mathcal{L}(\mathbf{z},t)$ $\label{eq:2.1} \frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\$ $\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1$ ~ 100 km s $^{-1}$

 $\label{eq:2.1} \frac{d\mathbf{r}}{d\mathbf{r}}\left(\mathbf{r}\right) = \frac{1}{2}\left(\mathbf{r}\right)^{2} \mathbf{r}^{2} \mathbf{r}^{2}$ $\label{eq:2.1} \frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\$

 $\label{eq:2.1} \Psi_{\mu\nu} = -\frac{1}{2} \frac{d\mu}{d\mu} \left[\frac{d\mu}{d\mu} \right] \left[\frac{d\mu}{d\mu$ $\label{eq:2.1} \frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\$ $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ and the contribution of the contribution of the contribution of $\mathcal{L}^{\mathcal{L}}$

 $\label{eq:2.1} \begin{split} \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) = \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{X}) \mathcal{L}_{\text{max}}(\mathbf{X},\mathbf{$ $\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1$

 $\label{eq:2.1} \Phi_{\mathbf{t}}(\mathbf{t},\mathbf{t}) = \mathbb{E}_{\mathbf{t}}\left[\mathbf{t}^{\top}(\mathbf{t},\mathbf{t})\right] = \mathbb{E}_{\mathbf{t}}\left[\mathbf{t}^{\top}(\mathbf{t},\mathbf{t})\right] = \mathbb{E}_{\mathbf{t}}\left[\mathbf{t}^{\top}(\mathbf{t},\mathbf{t})\right] = \mathbb{E}_{\mathbf{t}}\left[\mathbf{t}^{\top}(\mathbf{t},\mathbf{t})\right] = \mathbb{E}_{\mathbf{t}}\left[\mathbf{t}^{\top}(\mathbf{t$ $\mathcal{L}_{\text{max}}(\mathcal{L}_{\text{max}})$ and \mathcal{L}_{max} and \mathcal{L}_{max} $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal$

in the contract of the contrac $\label{eq:2.1} \frac{d\mathbf{r}}{d\mathbf{r}} = \frac{1}{2\sqrt{2}}\sum_{\mathbf{r}}\left(\mathbf{r}^{\mathbf{r}}_{\mathbf{r}}\right)^{2} \mathbf{r}^{\mathbf{r}}_{\mathbf{r}}$

 $\mathcal{L}(\mathcal{A})$ and $\mathcal{L}(\mathcal{A})$ are $\mathcal{L}(\mathcal{A})$. Then $\mathcal{L}(\mathcal{A})$ $\label{eq:2.1} \mathcal{L}(\mathcal{L}) = \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{1}{2$ $\label{eq:2.1} \mathcal{L}(\mathcal{L}) = \mathcal{L}(\mathcal{L}) \mathcal{L}(\mathcal{L}) = \mathcal{L}(\mathcal{L}) \mathcal{L}(\mathcal{L})$ $\frac{1}{2}$

 $\label{eq:2.1} \begin{array}{l} \mathcal{L}^{\mathcal{A}}(\mathbb{R}^d) = \mathbb{R}^d\\ \mathcal{L}^{\mathcal{A}}(\mathbb{R}^d) = \mathbb{R}^{d \times d} \end{array}$

. The second constraints of the second constraint in the second constraint $\mathcal{L}^{\mathcal{L}}$ $\sim 10^{11}$ km s $^{-1}$

 $\label{eq:2.1} \frac{\partial \mathcal{H}}{\partial \mathcal{H}} = -\frac{1}{2} \mathbf{q} + \mathcal{L} \mathbf{\tilde{T}} \mathbf{q} + \frac{1}{2} \mathbf{q} \mathbf{q} + \mathbf{q} \mathbf{q} + \frac{1}{2} \mathbf{q} \mathbf{q} + \frac{$ $\label{eq:2} \mathcal{L}^{(1)}(\mathbf{z}^{\mathcal{A}}) = \mathcal{L}^{(1)}(\mathbf{z}^{\mathcal{A}}) = \mathcal{L}^{(1)}(\mathbf{z}^{\mathcal{A}})$ $\label{eq:Ricci} \begin{split} \mathcal{F} &\rightarrow \mathcal{F}^{\prime} \quad \text{and} \quad \mathcal{F}^{\prime} \rightarrow \mathcal{F}^{\prime} \$

 $\langle \mathfrak{M} \rangle = \langle \mathfrak{M} \rangle$. $\label{eq:1} \frac{\Omega\epsilon}{\Omega^2}=\frac{2\Omega\mathcal{E}}{4\pi^2\omega^2}\,,$ where ω^2 is a particular function

 $\label{eq:2.1} \mathcal{L}^{(1)} = \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \$ THIS DOCUMENT WAS PRINTED USING RECYCLED PAPER

 $\label{eq:2.1} \frac{1}{2}\sum_{i=1}^n\frac{1}{2}\left(\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\left(\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\left(\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\right)\right)\right)^2}{\left(\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\left(\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{i=1}^n\frac{1}{2}\right)\right)^2}.$ $\mathcal{L}(\mathbf{z})$ and $\mathcal{L}(\mathbf{z})$ are the set of \mathbf{z}

 $\mathcal{L}(\mathcal{A})$ and $\mathcal{L}(\mathcal{A})$ are the set of the set of the set of \mathcal{A}

 $\label{eq:2.1} \begin{array}{l} \mathcal{S}_{\mathcal{A}} \left(\mathcal{S}_{\mathcal{A}} \right) \left(\mathcal{S}_{\mathcal{A}} \right) \\ \mathcal{S}_{\mathcal{A}} \left(\mathcal{S}_{\mathcal{A}} \right) \left(\mathcal{S}_{\mathcal{A}} \right$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ \sim \sim

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac$

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 $\label{eq:2.1} \mathcal{L}(\mathcal{A}) = \mathcal{L}(\mathcal{A}) \otimes \mathcal{L}(\mathcal{A}) \otimes \mathcal{L}(\mathcal{A})$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3} \left|\frac{d\mathbf{x}}{d\mathbf{x}}\right|^2 \, d\mathbf{x} \, d\mathbf{x$

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