

# Assessment of RELAP5/MOD 2 Against Marviken Jet Impingement Test **11** Level Swell

Prepared by **0.** Rosdahl, D. Caraher

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Office of Nuclear Regulatory Research U.S. Nuclear Regulatory Commission Washington, D.C. 20555

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Osten Rosdahl David Caraher\*

> ICAP ASSESSMENT OF RELAP5/MOD2 AGAINST MARVIKEN JET IMPINGEMENT TEST **11** LEVEL SWELL

# ABSTRACT

RELAP5/MOD2 simulations of level swell of saturated liquid in a large (5 m diameter, 22 m high) vessel are reported. For certain nodalizations RELAP5 is shown to predict the measured void fraction profile with fair accuracy. RELAP5 results are shown to be dependent upon nodalization with the accuracy of computed results deteriorating significantly when a large number of nodes is employed.

\* Engineering Software Consulting

Approved of Willshand

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APPENDIX

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# EXECUTIVE SUMMARY

RELAP5/MOD2 simulations have been conducted for MARVIKEN Jet Impingement Test **11** (JIT **11).** The purpose of the simulations was to assess the ability of RELAP5/MOD2 to simulate level swell in a large vessel.

The experimental facility consisted of a large vessel 5.2 m in diameter and 22 m high having a total volume of 420  $m^3$ . A standpipe 1 m in diameter and 18 m tall was inserted in the vessel. A discharge pipe containing a valve, nozzle, and rupture disks was attached to the lower end of the standpipe at the bottom of the vessel.

The vessel was filled to the 10.2 m elevation with nearly saturated liquid; the remaining part of the vessel and the standpipe were filled with saturated steam. The initial pressure in the vessel was 5.0 MPa.

The test was initiated by breaking the rupture disks. Because of the standpipe only steam flowed from the vessel. Differential pressures were recorded at various elevations in the vessel, thus affording a history of fluid density versus elevation to be obtained. Discharge mass flow rate was also measured. The experiment was terminated when the pressure in the vessel reached 1.9 MPa.

After the rupture disks broke bulk flashing occurred in the liquid. The level of the resulting two phase mixture rose rapidly and reached a maximum height of about 18 m - the top of the standpipe - within 15 seconds. The mixture level declined slowly thereafter receding to near the

14 m elevation by the time the test ended (80 seconds). For elevations below the 13 m height the differential pressure measurements remained fairly constant over the 15 to 80 second time period; indicating that the void fraction was fairly constant.

RELAP5/MOD2 simulations were conducted using 20, 40 and **100** nodes to model the annular region in the vessel below the top of the standpipe. The experimental mass flow rate was used as a boundary condition. Differential pressures calculated by RELAP5 were compared to measured ones.

The 20 node and the 40 node simulations showed similar results. Both calculations indicated that the RELAP5 underpredicted the void fraction of the swelled two phase mixture for elevations below 13 m and overpredicted the void fraction for higher elevations. The results imply that the interfacial drag force in RELAP5 fell off too rapidly with increasing void fraction. Consequently RELAP5 carried less liquid to the upper elevations than was actually carried there and RELAP5 allowed the liquid to drain from the upper elevations somewhat faster than it actually did.

The **100** node simulation was characterized by very erratic differential pressure histories which were, for some elevations, much different from the **dP** histories of the 20 and 40 node cases. Moreover, the **100** node simulation was found to be sensitive to time step size - changing the step size from **0.1** s to 0.05 s (material Courant limit = 0.12 s) produced large changes in void fraction profiles. The behaviour of the **100** node simulation is believed to be related to the

interphase drag model in RELAP5 - its strong dependence on void fraction in the bubble-toslug flow transition region; its explicit connection to the numerical solution; and its algorithm for damping large changes in computed values.

Time step studies on the 20 node model revealed that the RELAP5 calculation was sensitive to time step size during the time period (0-30 s) when the level swelled to its maximum height.

The 20 and 40 node simulations have demonstrated that RELAP5/MOD2 can give a fairly accurate simulation of the void profile occurring during level swell in large vessels. Improvements in the RELAP5 interphase drag model could lead to better simulations but would probably sacrifice the generality of the present model.

The **100** node simulation has demonstrated that the RELAP5 solution does not converge with increasing nodalization - at least not for time step sizes allowed by RELAP5's automatic time step control algorithm. The results suggest that at some level of nodalization the numerical models for smoothing locally computed variables begin to overshadow the physics.

# **1.** INTRODUCTION

The International Thermal-Hydraulic Code Assessment and Applications Program (ICAP) is being conducted by several countries and coordinated by the USNRC. The goal of ICAP is to make quantitative statements regarding the accuracy of the current state-of-the-art **•.** armal-hydraulic computer programs developed under the auspices of the USNRC.

Sweden's contributions to ICAP relate both to TRAC-PWR **(1)** and RELAP5 (2). The assessment calculations are being conducted by Studsvik Energiteknik AB for the Swedish Nuclear Power Inspectorate. The assessment matrix is shown in Table **1-1.**

This report presents the results of an assessment of RELAP5's ability to calculate level swell. In particular, Jet Impingement Test lI(JIT **11)** has been simulated using RELAP5. JIT **11** was one of a series of full scale jet impingement experiments conducted at the Marviken power station, Marviken, Sweden, during 1981.

This report is organized as follows: section 2 describes the experimental facility and section 3 describes the RELAP5 model used to simulate the experiments. In section 4 results from the simulations are presented and discussed. Computational efficiency of RELAP5 and numerical problems encountered during the simulations are given in section 5. Conclusions are presented in section 6.

# Table 1

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ICAP Assessment Matrix - Sweden.



# 2. FACILITY AND TEST DESCRIPTION

The Marviken Power Plant was built as a boiling heavy water direct cycle nuclear reactor but was never commissioned. The nuclear steam supply system was left intact and an oil fired boiler was built to provide steam for the turbine. During 1981 and 1982 Marviken was the site of the Jet Impingement Test (JIT) program. One of the tests, JIT **11,** examined the flow of steam from the Marviken vessel onto a test plate. Differential pressure measurements made in the vessel during JIT **11** provided data regarding the level swell occurring in the vessel as steam flowed out of the vessel.

Figure 2-1 depicts the Marviken pressure vessel and the location of the differential pressure measurements. For JIT **11** a standpipe was inserted into the vessel to ensure that only steam flowed out of the vessel. The outlet nozzle was located beneath the vessel. The piping leading to the nozzle and the nozzle are depicted in Figure 2-2. Initial and boundary condition for JIT **11** are summarized in Table 2-1. A complete description of the experimental facility for the JIT program is given in Reference 3. A description of JIT **<sup>11</sup>** is presented along with test results in Reference 4. The probable error (one standard deviation) in the measured differential pressure values shown in this report is 0.6 kPa; the 99 **%** confidence error is 1.5 kPa.

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Table 2-1
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Important parameters for Marviken JIT **11.**



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# Figure 2-1

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Marviken test vessel. Differential pressure transducers A through J, internal standpipe and initial (I) and final (F) inventory level are shown for Test T-11.



Figure 2-2

Arrangement of components in the discharge pipe.

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3. CODE AND MODEL DESCRIPTION

The simulations of JIT **11** were made with RELAP5/ MOD2, cycle 36.02.

# 3.1 Input description

The RELAP5 model of the Marviken vessel is shown in Figure 3-1. The standpipe was modelled by a single volume pipe component. At its upper end, this pipe was connected to a branch component which corresponded to the vessel region from the top of the standpipe to 0.5 m above the standpipe. The upper region of the vessel was modelled by a pipe component having 3 volumes.

The annular vessel region lying below the top of the standpipe was modelled as an annulus component. For the base case this component was divided into 20 nearly equally sized volumes. For nodalization studies the number of cells in this component was increased to 40 and then to 100.

A mass flow rate boundary condition was applied to the lower end of the component representing the standpipe. The boundary condition was taken from the experimental data.

The RELAP5 control system was used to calculate differential pressures for comparison with experimentally measured ones. In constructing the control system it was assumed that the pressure, which is a point quantity in RELAP5, varied linearly between the midpoints of adjacent cells. With this assumption pressure could be calculated at the exact location of each of the differential pressure taps. The assumption is a good one as long as the pressure variation with

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elevation is dominated by the gravity head of the fluid rather than its acceleration. This is the case throughout most of JIT 11 which was, for the most part, a quasi-steady flow test.



Figure 3-1

Schematic of RELAP5 model for JIT 11 level swell simulations.

# 4. RESULTS AND DISCUSSION

For the level swell simulations a discharge mass flow rate was imposed as a boundary condition. The boundary condition was obtained by drawing a smooth curve through the measured discharge flow rate. The boundary condition is compared to the measured flow rate in Figure 4-1. Figure 4-2, which compares the computed pressure response to the measured pressure response, demonstrates that the boundary condition served its purpose.

## 4.1 Base case

For the base case simulation the vessel region below the level of the standpipe was represented by 20 nodes  $(\Delta Z = 0.9 \text{ m})$ . Differential pressures were computed for comparison with measured differential pressures.

Figure 4-3 compares computed and measured dPs over the top and bottom halves of the vessel. The center tap for these measurements was initially located in the liquid but was uncovered at the end of the experiment. The fairly good agreement between the calculated and measured dPs indicates that RELAP5 calculated the gross behaviour of level swell correctly; that is, it correctly calculated the mass of liquid which was lifted above the 9.28 m elevation. It remains to be seen if the distribution of mass above and below the 9.28 m elevation was computed correctly.

The fact that the RELAP5 calculated dP203 lies below the measurement and the calculated dP205 lies slightly above the measurement during the quasi-steady part (20-70 s) of the experiment

indicates that RELAP5 lifted slightly more liquid above the 9.28 m elevation than was actually transported there.

Besides dP203 one other valid measurement of differential pressure below the 9.28 m elevation was obtained during JIT **11.** It is shown together with the corresponding RELAP5 value in Figure 4-4. The fact that the computed **dP** lies above the measured dP means that the average void fraction in the region being spanned by the **dP** taps was higher in RELAP5 than in the experiment (approximately 0.05 higher at 50 s).

The "overshoot" in the calculated dP in the 15-20 s time period is related to time step size. The base calculation generally proceeded at the maximum specified time step size (0.5 s). The differential pressure histories at all locations changed significantly in the 0-25 s range when time step size was reduced (see section 4.3).

Differential pressure measurement 213 spanned the initial liquid level surface (Figure 4-5). The experimental dP response over this span (9.23 to 11.43 m) showed that the average void fraction initially decreased (0-3 s) as the liquid level swelled then increased (3-20 s) to a value near the initial value as the two phase mixture rose to its maximum elevation. The average void fraction remained fairly constant until 60 s when it began decreasing as liquid drained into the region. When the outlet valve was completely closed (85 s) the average void fraction increased rapidly to unity as all liquid drained out of the region.

The RELAP5 simulation of dP213 behaved qualitatively the same as the measurement. However, RELAP5 calculated the average void fraction in the 9.23 m to 11.43 m region to be lower than it was in the experiment. The 2 kPa difference in dP between the measurement and the calculation corresponds to a void'fraction difference of about 0.13 between the experiment and the calculation  $($   $\sim$  0.60 in the experiment versus **-** 0.47 in the calculation).

When the outlet valve began closing at 75 s the RELAPS calculation responded immediately - the dP decreased, indicating that liquid was draining into the region from above. The experiment showed no such behaviour. The calculated dP began increasing at 82 s, before the valve was fully closed (84 s), and before the experimental dP began increasing. By 90 s the experimental dP had risen nearly to its maximum value (no liquid in the region) whereas the calculation showed a significant amount of liquid remained in the region. The inability of RELAP5 to mimic the experiment during the 5-20 s time period and the 85-90 s time period is attributed to the nature of the interfacial friction model in RELAP5. For the geometry and mass fluxes under consideration RELAP5 assumes that bubbly flow exists for void fractions less than 0.25 and slug flow exists for void fractions between 0.25 to 0.95. The interfacial friction associated with bubbly flow is much larger than that associated with slug flow. In the early (t **<** 20 s) part of the transient RELAP5 was unable to lift enough liquid out of the 9.23 - 11.43 region because the drag between the vapor and liquid diminished too quickly as the void fraction increased. Near the end of the transient RELAP5 was unable to let liquid drain out of the region as fast as it really did because

of an increasing drag between the rising vapor and falling liquid as the local void fraction dropped below 0.25. In the experiment it is quite likely that, once the valve was closed, the flow regime in the 9.23 - 11.43 region switched from slug flow to annular flow (the liquid draining downward along the vessel and sta lpipe surfaces) - a flow regime change which cannot be simulated in RELAP5 with the current flow regime maps.

The measured dP response over the 11.47 m to 13.46 m range (Figure 4-6) shows the same qualitative response as seen at the 9.23 to 11.43 m range. A rapid decrease in differential pressure occurred as liquid entered the region followed by an increase as the local void fraction increased. The dP, and hence the average void fraction, in the region remained relatively constant between 20 and 75 seconds. The dP increased sharply as the valve was closed and liquid drained out of the region.

RELAP5 calculated the dP response in the 9.23 - 11.43 range qualitatively correctly. However, it calculated too low a void fraction in the region during most of the transient. Unlike the calculated dP213 (Figure 4-5), the calculated dP212 showed the same response as the measured dP212 when the valve was closed. This is because the local void fraction in the computational cells which spanned the region never dropped below a value of 0.25 as they drained.

Differential pressure **dP211** (Figure 4-7) spanned the 13.46 m to 15.45 m elevation. The accuracy of this measurement is questionable because it did

not return to the correct final value at the end of the experiment - it was low by 4 kPa. The instruments response does appear to be correct at **<sup>10</sup>**s, because it indicates a dP value slightly less than that of dP210, which lies above it. This is reasonable, assuming a linear increase in void fraction with elevation. It is also likely that the instrument's response in the **<sup>10</sup>**- 50 s range is reasonably good because it remains fairly constant, exhibiting a behaviour similar to dP212 and dP213. If one assumes that the two phase level in the experiment declined steadily then **dP211** should have remained constant until 50 s and then should have increased at the same rate as **dP210** above it. In fact, the instrument shows this response which suggests that the instruments response is correct over the 0-80 s temperature range.

The uppermost dP measurement, dP210, spanned the 15.45 m to 17.45 m elevation. The response of **dP210** (Figure 4-8) shows that a two phase mixture passed the 15.45 elevation at about **10** seconds, reached its maximum level at 15 s, and slowly subsided over the next 30 seconds. Measurements of the discharge flow rate indicated that a small amount of liquid entered the standpipe during the 17 to 25 s time period. From this one concludes that the maximum elevation reached by the twophase level was about 18 m.

RELAPS calculated the **dP210** response to decline less than it actually did and to rise more rapidly. RELAP5 carried somewhat less liquid into the region spanned by **dP210** than did the experiment. RELAP5 also calculated that the liquid drained from the region much faster than it actually did. This faster drainage manifests itself in the different slopes of the measured and calculated dPs after 20 s.

Time step studies conducted subsequent to the base calculation have shown that when smaller time steps were used the RELAP5 calcr ated history of **dP210** dipped briefly to a value of about **-1.0** kPa then recovered to zero. Thus the calculated history shown in Figure 4-8 is an unconverged one.

The RELAP5 simulation of the level swell, which occurred in experiment JIT **11,** can be summarized as follows:

RELAP5 calculated that somewhat more liquid was lifted above the initial liquid level than actually was (Figure 4-3, 4-4). Even though RELAP5 carried more liquid above the initial liquid level, it did not calculate the two-phase level to rise and remain as high as it did in the experiments (Figure 4-8, 4-7).

For the quasi-steady part of the experiment (30 - 50 s) RELAP5 calculated the liquid fraction profile in the swelled region (the region above the original water level) to be more concentrated in the lower elevations than it actually was (Figure 4-5 and 4-6). The latter observation is a recognition of the fact that the RELAP5 overprediction of void fraction in the 4.97 to 9.23 m range (Figure 4-4) is more than offset by the underprediction of the void fraction in the 9.23 to 13.43 m range (Figures 4-5, 4-6.

# 4.2 Nodalization studies

In order to determine the sensitivity of computed results to nodalization the RELAP5 calculation of JIT **11** was repeated using 40 and **100** nodes to represent the 18 m long annular region from the vessel bottom to the top of the standpipe. Some results from the 20, 40, and **100** node cases are given in Figures 4-9 through 4-12.

The most significant feature in Figures 4-9 through 4-12 is the erratic nature of the **<sup>100</sup>** node calculation. Other notable features are the reduction, with increasing nodes, in the dP "overshoot" in the lower elevations (Figure 4-9,  $4-10$ ;  $15 < t < 25 s$ ), and the decreasing calculated height of the two-phase level with increasing number of nodes (Figure 4-12). Figures 4-9 to 4-12 show that the RELAP5 calculation is sensitive to nodalization and does not converge with increasing nodalization.

The reason for the erratic calculated dPs in the **<sup>100</sup>**node calculation is revealed in Figure 4-13, which shows the void fraction profile in the vessel at 42.5 and 43.5 seconds. For the region spanned by dP212 (11.47 to 13.46 m) the average void fraction changes from a value of 0.49 at 42.5 s to a value of 0.64 at 43.5 s. The calculated change in **dP** use to this change in void fraction is  $(\Delta dP = \Delta \alpha \rho_{fq} \Delta z)$  2.3 kPa which is about the change in dP shown in Figure 4-10. The erratic nature of the results from the **100** node calculation are due to RELAP5 calculating a fluctuating axial void profile (for elevations between 9 and 15'm) rather than the quasi-steady void profile indicated by the experimental measurements and by the 40 and 20 node RELAP5 calculations.

# 4.3 Time step studies

The **100** node calculation was examined further by restarting it at 40 s and reducing the maximum time step from **0.1** s to 0.05 s (the material Courant limit was 0.12 s). The axial void profile from these two calculations are shown in Figure 4-14. Local differences in void fraction are large even though the average step size utilized for the two calculations (0.075 s and 0.05 s for  $DTMAX = 0.1$  and  $DTMAX = 0.05$  respectively) was much less than the Courant limit.

The sensitivity of the **100** node calculation to time step size demonstrates the unreliability of the mass error time step control algorithm - the algorithm does not necessarily produce time step sizes which lead to a convergent solution.

It is believed that the erratic solution and the time step sensitivity exhibited to the **100** node calculated are related to the strong coupling between void fraction and interfacial drag in the 20 - 40 percent void fraction range. Because of the large flow areas and short lengths of the cells in the **100** node simulation void fractions were observed to sometimes change dramatically from one time step to the next. As a consequence the interfacial drag changed by as much as a factor of **10** from one time step to the next.

Having seen the sensitivity of the **100** node calculations to time step size we decided to investigate the sensitivity of the base (20 node) calculation to time step size.

The original calculation proceeded at the specified maximum time step which was 0.5 s. For the

two sensitivity studies the maximum step size was set to **0.1** s and 0.05 s. The results of the sensitivity calculations are depicted, along with the base case results in Figure 4-15 to 4-19.

The figures show that the RELAP5 base case solution was not converged during the first 30 seconds of the transient. Using large time steps, although entirely consistent with the mass error time step control algorithm, gave an artificial increase in the amount of liquid being carried upward.

It is believed that the sensitivity to time step is caused by the relaxation algorithm being used in the interphase drag model. The model calculates an interphase drag coefficient based upon new time level hydraulic parameters and then, to get a interphase drag coefficient for use in the momentum equations, weights the new time level value with the old time level value, giving favor to the old value. The result of the scheme is that the code "remembers" previous values of the interphase drag for several time steps regardless of time step size. During a period of rapidly changing interphase drag (bubble-to-slug flow) the computed results therefore become sensitive to time step size.



# Figure 4-1

Mass flow rate boundary condition and experimental discharge flow rate.



# Figure 4-2

Measured and calculated pressure response for **JIT 11.** 



 $\frac{Figure 4-3}{}$ Comparison of computed and measured diverse



# Figure  $4-4$ <br> $\frac{1}{2}$  and calculated dependence  $4.97 \text{ m}$ Measured

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# Figure 4-5





# Figure 4-6

Measured and calculated dP over the 11.47 to 13.43 m range.



# Figure 4-8

Measured and calculated dP over the 15.45 to 17.45 m span.

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# Figure 4-12

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Measured and calculated dPs over the 15.45 to 17.45 m elevation.

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Axial void profile for the 100 node simulation. DTMAX =  $0.1$  s and DTMAX =  $0.05$  s.



Sensitivity of dP213 to time step size.

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Figure 4-18 Sensitivity of dP211 to time step size.

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Sensitivity of dP210 to time step size.

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### 5. COMPUTATIONAL EFFICIENCY

The computational efficiency of the RELAP5 simulations is summarized in Table 5-1. The CPU times are for a CYBER 180-810 computer.

In general, the simulations used the maximum time step size specified in the input. This time step size was usually slightly smaller than the material Courant limit.

# Table 5-1

Run time statistics for JIT **11** simulations. (ninety second transient)



# 6. CONCLUSIONS

Assessment of RELAP5 against the level swell data obtained from JIT **11** led to the following conclusions:

- **1.** For the 9 to 13 m range RELAP5 calculated a quasi-steady void fraction which was 0.10 to 0.15 lower than measurements indicated.
- 2. For the 13 to 18 m range RELAP5 calculated a quasi-steady void fraction which was too high. In the base case RELAP calculated the maximum level of the two phase front well but allowed liquid to drain out of the upper region of the vessel faster than it actually did.
- 3. Nodalization studies demonstrated that the RELAP5 solution did not converge with increasing number of nodes. The 20 node case and the 40 node case gave similar results. The **100** node case, however, exhibited erratic fluctuations in the axial void profile and was sensitive to time step sizes. The behaviour of the **<sup>100</sup>**node case is believed to be caused by the interphase drag model.
- 4. The computed results were sensitive to time step size. Time step size had to be reduced considerably below the value allowed by the RELAP5 time step control algorithm in order to obtain converged results during the time when the level was rising. The sensitivity is believed to be caused by the relaxation algorithm in the interphase drag model.

The JIT **11** simulations have shown that the RELAP5 interphase drag model allows the interphase drag force to diminish too rapidly with increasing void fraction for the case of level swell in large vessels. It would, however, be premature to recommend changes in the interphase drag model based upon the results reported here so long as

the model is intended to be used for a variety of geometries (tubes, rod bundles, vessels). Refinements of the interphase drag model should consider separating the treatment of level swell in large vessels from the treatment of flow in ducts.

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7. REFERENCES

 $\mathcal{L}_{\text{max}}$ 

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### APPENDIX A

# RELAPS INPUT DECK

=RELAPS INPUT FOR JIT-11 LEVEL SWELL (CASE **I. RUN** 2) **4 \*** MARVIKEN JET IMPINGEMENT TEST (JIT) NO. **H1 <sup>4</sup>**IS USED FOR **A** LEVEL SWELL CALCULATION WITH RELAPS/MOD2/36.02 **AS A** PART OF THE ICAAP AGREEMENT. **\*** STEADY STATE RUN TO GET INITIAL PRESSURES **<sup>4</sup>**COMPONENIS 902,952,903,gs3,AND TABLE 210 WILL BE DELETED ON THE RESTART FOR TRANSIENT RUN.  $\bullet$ **100** NEW STOY-ST  $\sim 100$ 101 RUN 105 **10.** 20.  $\bullet$ **4** END DTMIN DTREQ FLAGS MINED MAJED RESTART **2'0.0** 1.0-6 **0.50** 00003 **I** 20 **100** 201 **4 <sup>4</sup>**MINOR EDITS **-------------- 301** CNTRLVAR 313 **302** CNTRLVAR 101 303 **CITT** RL **VAR 210** *304* MFLOUJ 200030000 305 MFLOWJ 9S1000000 306 TEMPF 100010000 TEMPF 307 100050000 308 TEMPF 100100000 1001500c0 309 TEMPF 311 CNTRLVAR 403 312 CNTRLVAR 40S 313 CNTRLVAR 410 314 CNTRLVAR 41 **315** CNTRLVAR 412 **316** CNTRLVAR 41 317 CNTRLVAR 414 318 CNTRLVAR 415 319 CNTRLVAR 41'  $\bullet$ TRIPS \*------------ $\bullet$ 501 TIME 0 GT NULL 100. L

601 501 AND 501 L **4** COMPONENT 100 \* REACTOR VESSEL BELOW TOP OF STANDPIPE • LEVELS: 0. - **18.33** 1000000 VESSLOW ANNULUS • NV 1000001 20 \* VOLAREA VOLNO 1000101 0. 20 **4 <sup>4</sup>**LENGTH VOLNO 1000301 1.23 **1** 1000302 0.90  $\bullet$ **<sup>4</sup>**VOLUME VOLNO 1000401 9.22 **1** 1000402 15.20 2 1000403 17.78 1000404 18.194 20 VANGLE VOLNO 1000601 90. 20  $\bullet$ • WROUGH HYD-DIA VOLNO 1000801 0. 4. 20 **4** \* VOLFLGS VOLNO 1001001 00 20  $\bullet$ • JUNFLGS JUNNO 1001101 01000 19  $\bullet$  $\bullet$ \* COMPONENT 200 **<sup>4</sup>**REACTOR VESSEL ABOVE TOP OF STANOPIPE, BELOW SPRAY SHIELD • LEVELS: 18.33 - 18.83 2000000 VESSMID BRANCH  $\bullet$ \* NJ 2000001 3 \* VOLAREA LENGTH VOLUME AZI INCL ELEV ROUGH DIA **FLGS** 2000101 0. 0.50 10.79 0. 90. 0.50 0. 0. **c0**  $\bullet$ \* FROM TO JUNAREA LOSSF LOSER FLGS 2001101 100010000 200000000 0. 0. 0. 01100 2002101 200010000 300000000 19.6 0. 0. 01100 2003101 200000000 400010000 0. 0. **0.** 01100



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**\*** UJOLFLGS VOLNO 4001001 **00 I**  $\bullet$ **\*** COMPONENT 901 **4** ATMOSPHERE **9** 9010000 ATMOSPH Ti1DPVOL VOLAREA LENGTH VOLUME AZI INCL ELEV 90. 0.50 ROUGH DIA FLGS 10. **0.** 9010101 10. 1. 0. 0. 00 CW **9** 9010200 2  $\bullet$  ,  $\bullet$  ,  $\bullet$  ,  $\bullet$ \* TIME PRESS QUAL 9.10201 **0. O.IE6 1. 4** • CI OMPONENT **951 4** N( OZZLE FLOW **AS** FUNCTION OF TIME 9510000 NOZZLE TMDPJUN **4** FROM TO JUNAREA  $\bullet$  . 9510101 400000000 901000000 0  $\bullet$  $\bullet$ cW 9510200 **1** 601  $\bullet$ TIME LIOFLOW VAPFLOW ZEROL 9510201 **0. 0. 0.** 0. **ss 10202 0.** 624. **0.** 0.04 0.06 9510203 **0.** 632. **0.** 9510204 208. **0. 0.0B** 0. 9510205 0.10 **0.** 647. **0.** 0.16 9510206 **0.** 125. **0.** 9510207 0.20 139. **0. 0.** 0.24 9510208 349. **0. 0.** 9510209 0.28 **0.** 432. 0. 0.32 9510210 470. 432. 0. **0.** 0.40 9510211 **0.** 0. 0. 0.46 9510212 389. 0. 0.  $\bar{\mathbf{a}}$  . 0.52 393. 9510213 .<br>M 0.60 9510214 411. O0.  $\overline{a}$ 0.70 9510215 415. 0.  $\tilde{\mathbf{e}}$ 0.80 403. 9510216 0.  $\tilde{a}$ 9510217 392. **I.** 0.  $\frac{1}{2}$ 9510218 1.5 392. 0.  $\bar{\mathbf{a}}$ . 9510219 401. 2. 0. 0. 9510220 5. 394. 0. ດົ 9510221 7. **390.** 0.  $\bar{a}$ 9510222 375<br>336 9. 0.  $\bar{o}$ 9610223 16. **32.** 0. 289. 228. .<br>M 9510224 0.  $\bar{a}$ 9510225 50. 0.  $\overline{a}$ 9510226 75. 175.

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20549901 0. 1. CNTRLVAR 414
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* STEADY STATE CARDS
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\bullet\bulletCOMPONENT 902
• STEADY STATE PRESSURIZER
9020000 PRZ
                    TMDPVOL
\bullet .
• VOLAREA
LENGTH VOLUME AZI INCL ELEV ROUGH DIA FLGS
9020101 10.
1. 10. 0. 90. 1. 0. 0. 00
        cW
9020200
2
\bullet4 TIME
                      PRESS
                              QUAL
9020201 0.
                      5.OEE
                               I.
\bullet4
   COMPONENT 952
     STEADY STATE CONNECTION TO THE PRESSURIZER
9
9520000 PRZLIN
SNGLJUN
\bullet* FROM<br>9520101 3000100
                    TO JUNAREA
LOSF
0.
LOSSR
                                                       FL G
01100J
9520101 300010000
902000000 0.
                                               0.
9520201 0
                   0. 0. 0.
\bullet\bullet\bullet• COMPONENT 903
4STEADY STATE LIOUID FILL VOLUME
9030000 FILL
TMOPVOL
\bullet4 VOLAREA
LENGTH VOLUME AZI INCL ELEV ROUGH DIA FLGS
9030101 10.
1. 10. 0. 90. 1. 0. 0. 00
\bulletCW
4
9030200
3
4
         TIME
PRESS QUAL
9030201 0.
                     5.OEB 533.
\bulletCOMPONENT 953
* STEADY STATE FILL
JUNCTION TO OBTAIN THE CORRECT WATER LEVEL
9530000 FILLJUN
                   TMOPJUN
\bullet• FROM
                   TO JUNAREA
9530101 903000000
100000000 0.
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# 1986-07-11



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 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}})) \leq \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}}))$ 

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ 

 $\label{eq:2.1} \begin{split} \mathcal{L}_{\text{max}}(\mathbf{r}) & = \frac{1}{2} \sum_{i=1}^{N} \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \\ & = \frac{1}{2} \sum_{i=1}^{N} \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf$ 

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$  $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$ 



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 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}})) \leq \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}}))$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$ 



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OFFICIAL BUSINESS<br>PENALTY FOR PRIVATE USE, \$300

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