DRAGON Theory for 3-D CANDU Problems

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Summary 1

1. Introduction to the Collision Probability Method
2. 3–D Collision Probability Calculations
3. Solving the Collision Probability Equations
4. Condensation and Homogenization Techniques
5. Discussion and Conclusion

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Introduction to the CP Method 1

Contents

- The Transport Equation
- The Collision Probability Technique
- Boundary Conditions
- Cross Sections Considerations

The Transport Equation 1

The transport equation is a neutron balance equation

$$\mathcal{L}(\vec{r}, E, \vec{\Omega}) = \mathcal{Q}(\vec{r}, E, \vec{\Omega})$$

- $\mathcal{L}$ represents neutron lost from the system:

$$\mathcal{L} = \vec{\Omega} \cdot \vec{\nabla}\Phi(\vec{r}, E, \vec{\Omega}) + \Sigma(\vec{r}, E)\Phi(\vec{r}, E, \vec{\Omega})$$

- $\mathcal{Q}$ represents neutron created in the system

$$\mathcal{Q}_s = \int dE' d^2\Omega \Sigma_s(\vec{r}, E' \rightarrow E, \vec{\Omega}' \rightarrow \vec{\Omega})\Phi(\vec{r}, E', \vec{\Omega}')$$

$$\mathcal{Q}_f = \chi(\vec{r}, E) \int dE' d^2\Omega \nu\Sigma_f(\vec{r}, E')\Phi(\vec{r}, E', \vec{\Omega}')$$

The Transport Equation 2

Multigroup transport equation

$$\left[\vec{\Omega}\cdot\vec{\nabla}+\Sigma^g(\vec{r})+D^g(\vec{r})B^2\right]\Phi^g(\vec{r},\vec{\Omega})=Q_s^g(\vec{r},\vec{\Omega})+\frac{1}{k}Q_f^g(\vec{r},\vec{\Omega})$$

Scattering source $Q_s^g(\vec{r},\vec{\Omega})$

$$Q_s^g(\vec{r},\vec{\Omega})=\sum_{h=1}^{G}\int d^2\Omega'\Sigma_s^{hg}(\vec{r},\vec{\Omega}'\rightarrow\vec{\Omega})\Phi^h(\vec{r},\vec{\Omega}')$$

Fission source $Q_f^g(\vec{r},\vec{\Omega})$

$$Q_f^g(\vec{r})=\chi^g\sum_{h=1}^{G}\nu\Sigma_f^h(\vec{r})\int d^2\Omega'\Phi^h(\vec{r},\vec{\Omega}')$$

The Transport Equation 3

Transport equation in the absence of external sources is an eigenvalue problem:

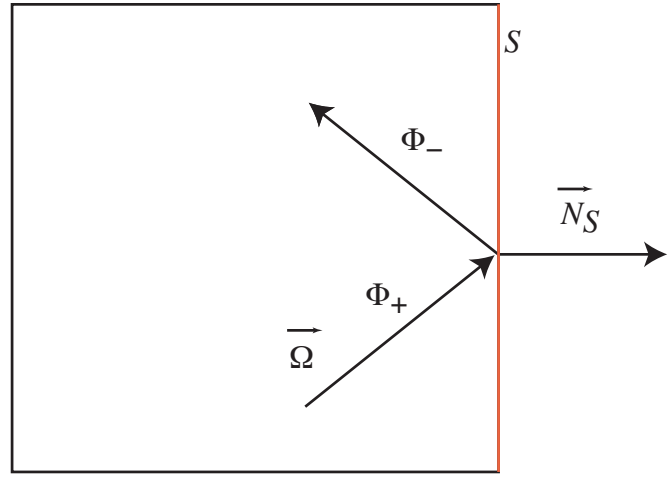
1. k eigenvalue with imposed leakage ($D^g(\vec{r})B^2$ fixed):
   - k indicates how the fission rate should be modified to make the system critical (reach a non-trivial solution to the transport equation)
2. Buckling eigenvalue with imposed k:
   - $D^g(\vec{r})B^2$ represent the amount of leakage required to make the system critical

The Transport Equation 4

Boundary conditions are required to close the system

- Albedo conditions

$$\phi_-(\vec{r}_S, \vec{\Omega} - 2(\vec{N}_S \cdot \vec{\Omega})) = \beta(\vec{r}_S, \vec{\Omega})\phi_+(\vec{r}_S, \vec{\Omega})$$



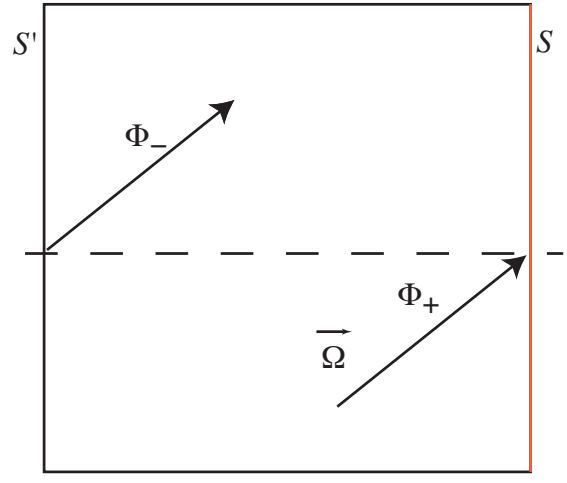
$\beta(\vec{r}_S, \vec{\Omega}) = 0$ for void BC

$\beta(\vec{r}_S, \vec{\Omega}) = 1$ for reflective BC

The Transport Equation 5

- Periodic conditions

$$\phi_-(\vec{r}_S, \vec{\Omega}) = \phi_+(\vec{r}_{S'}, \vec{\Omega})$$



The Transport Equation 6

Integral transport equation (case without leakage)

- Flux at a point $\vec{r}$ due to neutrons created at any point $\vec{r}' = \vec{r} - R\vec{\Omega}$ surrounding it

$$\left[-\frac{d}{dR} + \Sigma^g(\vec{r} - R\vec{\Omega})\right] \Phi^g(\vec{r} - R\vec{\Omega}, \vec{\Omega}) = Q^g(\vec{r} - R\vec{\Omega}, \vec{\Omega})$$

- Integrate transport equation over R and $\vec{\Omega}$

$$\phi^g(\vec{r}) = \int d^2\Omega \, e^{-\tau^g(R_S)} (\vec{\Omega} \cdot \vec{N}_-) \Phi^g_-(\vec{r}'_S, \vec{\Omega}) \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega})$$
$$+ \int d^2\Omega \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \vec{\Omega}) \Theta(\vec{r}, \vec{r}', \vec{\Omega}) dR'$$

The Transport Equation 7

Definitions

$$\phi^g(\vec{r}) = \int d^2\Omega \Phi^g(\vec{r}, \vec{\Omega})$$

$$\tau^g(R) = \int_O^R \Sigma^g(\vec{r} - R'\vec{\Omega}) dR'$$

$$\Theta(\vec{r}, \vec{r}', \vec{\Omega}) = \begin{cases} 1 & \text{if } \vec{r} = \vec{r}' + R'\vec{\Omega} \\ 0 & \text{otherwise} \end{cases}$$

and $\Phi_-^g(\vec{r}'_S, \vec{\Omega})$ is the incoming angular flux on surface S

The Transport Equation 8

Equation for the outgoing flux $\phi_+(\vec{r}_S)$ at S

$$\phi_+^g(\vec{r}_S) = \int d^2\Omega(\vec{\Omega}\cdot\vec{N}_+)\int_0^R e^{-\tau^g(R')}Q^g(\vec{r}',\vec{\Omega})\Theta(\vec{r}_S,\vec{r}',\vec{\Omega})dR'$$

$$+\int d^2\Omega(\vec{\Omega}\cdot\vec{N}_+)(\vec{\Omega}\cdot\vec{N}_-)\, e^{-\tau^g(R_S)}\Phi_-^g(\vec{r}'_S,\vec{\Omega})\Theta(\vec{r}_S,\vec{r}'_S,\vec{\Omega})$$

where

$$\phi_+^g(\vec{r}_S) = \int d^2\Omega\ \Phi_+^g(\vec{r}_S,\vec{\Omega})(\vec{\Omega}\cdot\vec{N}_+)$$

The Transport Equation 9

CP approximations

- Divide domain into N_V regions of volume V_i where the cross sections sources are independent of $\vec{r}$ and $\vec{\Omega}$

$$\Sigma^g(\vec{r}) = \Sigma^g_j \qquad \text{for } \vec{r} \in V_j$$

$$Q^g(\vec{r},\vec{\Omega}) = \frac{1}{4\pi} q^g(\vec{r}) = q^g_j \qquad \text{for } \vec{r} \in V_j$$

(the source is assumed isotropic)

- Divide the external boundary S into N_S surfaces of area S_α and assume angular flux constant on these surfaces

$$\Phi^g_-(\vec{r}_S,\vec{\Omega}) = \frac{1}{4\pi}\phi^g_{\alpha,-}$$

The Transport Equation

$$\int d^2\Omega \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \vec{\Omega}) \Theta(\vec{r}, \vec{r}', \vec{\Omega}) dR' =$$

$$q_j^g \int_{V_i} e^{-\tau^g(R')} \Theta(\vec{r}, \vec{r}', \vec{\Omega}) \frac{d^3 r}{4\pi |\vec{r} - \vec{r}'|^2}$$

$$\int d^2\Omega \; e^{-\tau^g(R_S)} \Phi_-^g(\vec{r}'_S, \vec{\Omega}) \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega}) =$$

$$\phi_{\alpha,-}^g \int_{S_\alpha} e^{-\tau^g(R_S)} \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega}) \frac{d^2 r}{4\pi |\vec{r} - \vec{r}'|^2}$$

The Transport Equation

Transport equations in CP form

$$\phi_i^g = \sum_{\alpha=1}^{N_S} p_{i\alpha}^g \phi_{-,\alpha}^g + \sum_{j=1}^{N_V} p_{ij}^g q_j^g$$

$$\phi_{+,\alpha}^g = \sum_{\beta=1}^{N_S} p_{\alpha\beta}^g \phi_{-,\beta}^g + \sum_{j=1}^{N_V} p_{\alpha j}^g q_j^g$$

where

$$\phi_i^g = \frac{1}{V_i} \int_{V_i} d^3 r \phi^g(\vec{r})$$

$$\phi_{+,\alpha}^g = \frac{1}{S_\alpha} \int_{S_\alpha} d^2 r \phi^g(\vec{r}_S)$$

The Transport Equation

Four types of probabilities

$$\tilde{p}^g_{ij} = V_i p^g_{ij} = \int_{V_i}\int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2}\Theta_i\Theta_j d^3r' d^3r$$

$$\tilde{p}^g_{i\alpha} = V_i p^g_{i\alpha} = \int_{V_i}\int_{S_\alpha} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2}(\vec{\Omega}\cdot\vec{N}_-)\Theta_i\Theta_\alpha d^3r' d^2r$$

$$\tilde{p}^g_{\alpha i} = \frac{S_\alpha}{4} p^g_{\alpha i} = \int_{S_\alpha}\int_{V_i} \frac{e^{-\tau^g(R)}}{4\pi R^2}(\vec{\Omega}\cdot\vec{N}_+)\Theta_\alpha\Theta_i d^2r' d^3r$$

$$\tilde{p}^g_{\alpha\beta} = \frac{S_\alpha}{4} p^g_{\alpha\beta} = \int_{S_\alpha}\int_{S_\beta} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2}(\vec{\Omega}\cdot\vec{N}_-)(\vec{\Omega}\cdot\vec{N}_+)$$
$$\times\Theta_\alpha\Theta_\beta d^2r d^2r'$$

The Transport Equation

Symmetry relations

$$
\begin{aligned}
V_i p^g_{ij} &= V_j p^g_{ji} \\
4 V_i p^g_{i\alpha} &= S_\alpha p^g_{\alpha i} \\
S_\alpha p^g_{\alpha\beta} &= S_\beta p^g_{\beta\alpha}
\end{aligned}
$$

Conservation properties

$$
\sum_{\alpha=1}^{N_\alpha} p^g_{i\alpha} + \sum_{j=1}^{N_j} p^g_{ij} \Sigma^g_j = 1
$$

$$
\sum_{\beta=1}^{N_\beta} p^g_{\alpha\beta} + \sum_{i=1}^{N_i} p^g_{\alpha i} \Sigma_i = 1
$$

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Boundary Conditions 1

Surface flux approximation:

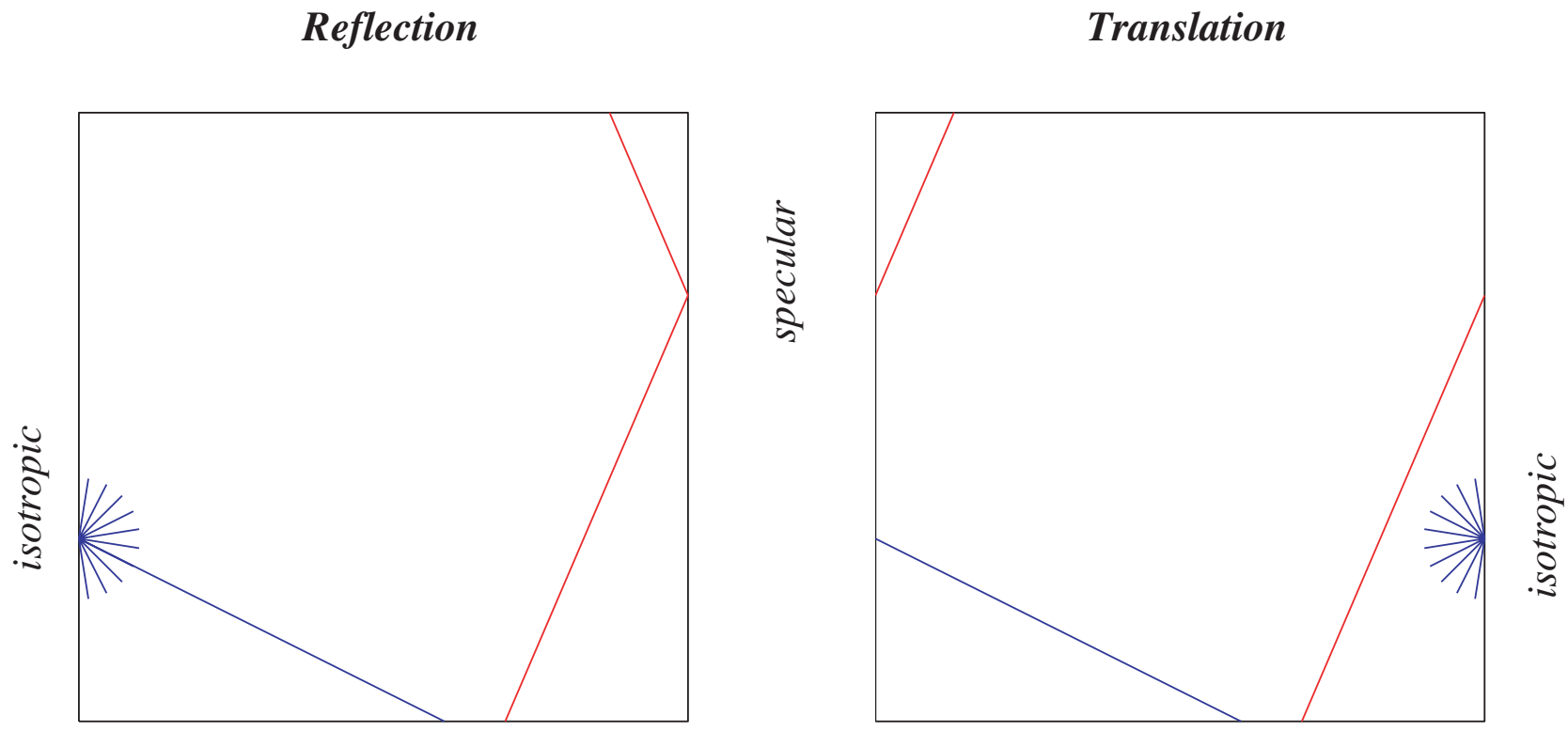
- Incoming angular flux on outer surfaces assumed to be independent of $\vec{\Omega}$

Comments

- Outgoing angular flux on outer surfaces integrated over $\vec{\Omega}$
- Angular flux not used at region interfaces
- Approximation for incoming angular flux exact for vacuum BC
- Approximation for incoming angular flux leads to large errors in surface flux when flux is not isotropic

Boundary Conditions 2

Illustration of approximate BC



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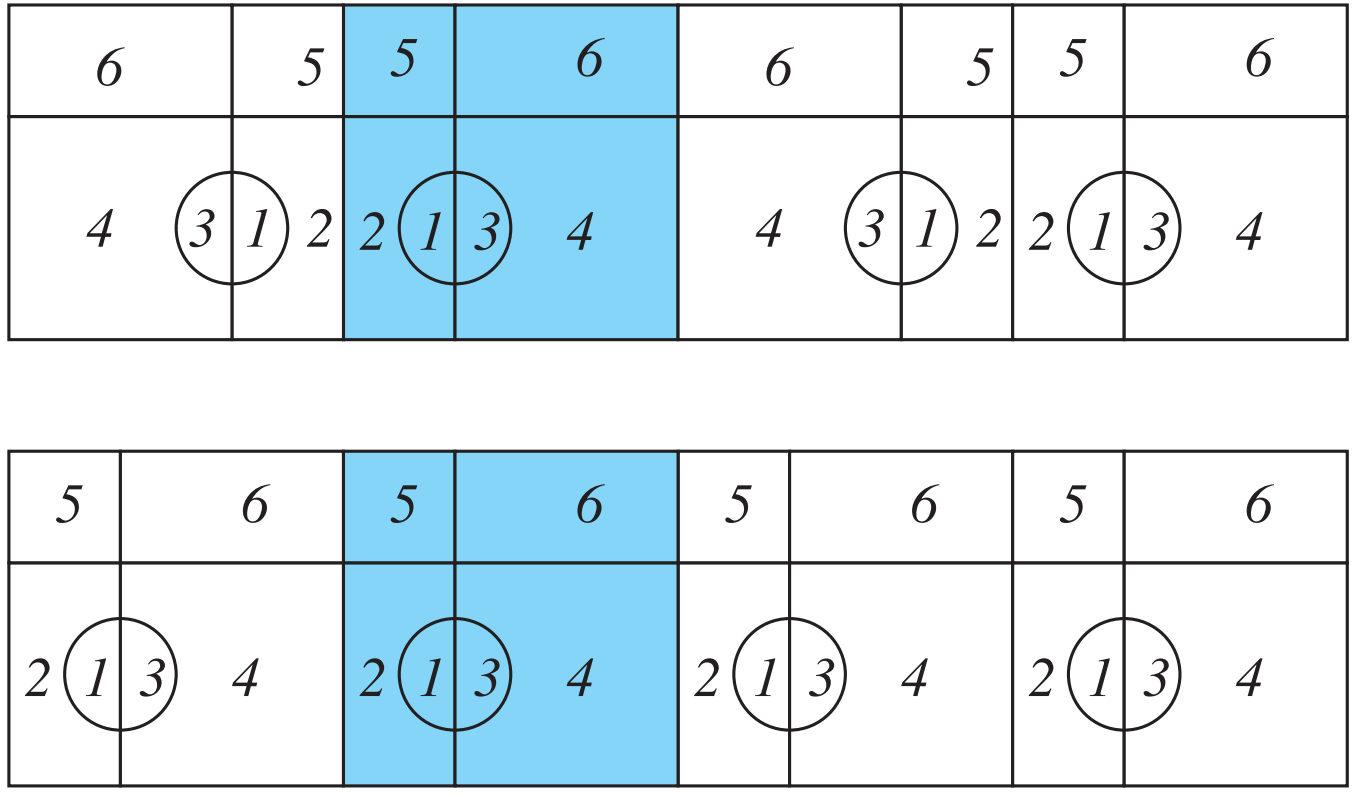
Boundary Conditions 3

Recommendations to reduce errors due to approximate use of BC for a fixed direction

- No special treatment for 2 vacuum BC
- Unfold cell once for 1 vacuum and 1 reflection BC
- Multiply unfold cell for 2 reflection or periodic cell, apply approximate BC on final surfaces and consider results in cell located far from these surfaces.

Boundary Conditions 4

Example of cell unfolding in direction X



Boundary Conditions

Simplifying CP equations using approximate BC

- Assume

$$\vec{J}^g_- = \mathbf{A}^g \vec{J}^g_+$$

- Final transport equation

$$\vec{\phi}^g = \mathbf{P}^g_{c,vv} \vec{q}^g$$

with the complete collision probability matrix $\mathbf{P}^g_{c,vv}$:

$$\mathbf{P}^g_{c,vv} = \left(\mathbf{P}^g_{vv} + \mathbf{P}^g_{vs}((\mathbf{A}^g)^{-1} - \mathbf{P}^g_{ss})^{-1} \mathbf{P}^g_{sv} \right)$$

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Cross Sections Consideration 1

Two types of multigroup cross-section database can be read by DRAGON

- Mixture macroscopic cross-section
- Isotope microscopic cross-section that contains itself a macroscopic cross-section database

Cross Sections Consideration 2

Minimum cross-section requirements for each mixture m

- The multigroup total cross section Σ_m^g
- The isotropic component of the multigroup scattering cross section $\Sigma_{m,s,0}^{h\to g}$ defined as

$$\Sigma_{m,s,0}^{h\to g} = \int_{4\pi} d^2\Omega^2 \Sigma_{m,s}^{h\to g}(\vec{\Omega}' \to \vec{\Omega}) P_0(\vec{\Omega}' \cdot \vec{\Omega})$$

- The product of the average neutron emitted per fission with the multigroup fission cross section $\nu\Sigma_{m,f}^g$
- The multigroup fission spectrum χ_m^g

Cross Sections Consideration 3

The linearly isotropic component of the multigroup scattering cross section $\Sigma_{m,s,1}^{h\to g}$

$$\Sigma_{m,s,1}^{h\to g} = \int_{4\pi} d^2\Omega^2 \Sigma_{m,s}^{h\to g}(\vec{\Omega}' \to \vec{\Omega}) P_1(\vec{\Omega}' \cdot \vec{\Omega})$$

- Required only if B_1 leakage method is used

Cross Sections Consideration 4

The transport correction $\Sigma^g_{m,\text{tc}}$

- The transport calculations are performed using transport corrected total ($\tilde{\Sigma}^g_m$) and scattering ($\tilde{\Sigma}^{h\to g}_{m,s,0}$) cross sections

$$\tilde{\Sigma}^g_m = \Sigma^g_m - \Sigma^g_{m,\text{tc}}$$
$$\tilde{\Sigma}^{h\to g}_{m,s,0} = \Sigma^{h\to g}_{m,s,0} - \delta^{gh}\Sigma^g_{m,\text{tc}}$$

- Takes partially into account the linearly anisotropic scattering contributions

Cross Sections Consideration 5

For contribution of multi-neutron production reactions such as

$$\mathrm{x}_A^Z + \mathrm{n}_1^0 \rightarrow \mathrm{x}_{A-1}^Z + 2\mathrm{n}_1^0 + \gamma$$

The scattering cross section must be corrected to take into account this effect

$$\tilde{\Sigma}_{m,s,0}^{h\rightarrow g} = \Sigma_{m,s,0}^{h\rightarrow g} + 2\delta^{gh}\Sigma_{m,(\mathrm{n},2\mathrm{n})}^{g}$$

where $\Sigma_{(\mathrm{n},2\mathrm{n})}^{g}$ is the macroscopic cross section associated with the reaction

Cross Sections Consideration 6

Macroscopic cross section data base can be created using

- from the input file using the `MAC:` module
- from a GOXS file using the `MAC:` module
- from a microscopic library using the `LIB:` module
- from the homogenization and condensation module `EDI:`
- from a WIMS-AECL execution using the information available on `TAPE16` (side-step method)
- from a HELIOS execution

Microscopic libraries 1

Many formats can be processed by DRAGON including

- WIMS–AECL format
- MATXS format
- WIMD-D4 format

In DRAGON resonance self-shielding calculations are performed using the Stamm'ler method

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3–D CP Calculations 1

Contents

- Collision Probabilities in 3-D
- Numerical Quadrature and Tracking
- Collision Probability Integration
- Neutron Conservation and CP Normalization

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3–D CP Calculations 2

Recall CP approximations

- Divide domain into N_V regions of volume V_i
- Assume total cross sections constant inside each region
- Assume sources constant inside each region
  - This has an impact on the selection of the spatial mesh
- Assume sources isotropic inside each region
  - This may lead to problem when scattering is highly anisotropic.

3–D CP Calculations 3

Recall BC approximations

- Divide the external boundary S into N_S surfaces of area S_α
- Assume flux constant on external surfaces
  - This has an impact on the selection of the spatial mesh
- Assume flux isotropic on external surfaces
  - This may lead to problems when flux is highly anisotropic near external boundaries

Collision Probabilities in 3-D 1

Collision probability definition

$$\tilde{p}_{ij}^{g} = V_i p_{ij}^{g} = \int_{V_i} \int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2} \Theta_i \Theta_j d^3r' d^3r$$

- Spherical coordinates for d^3r integral

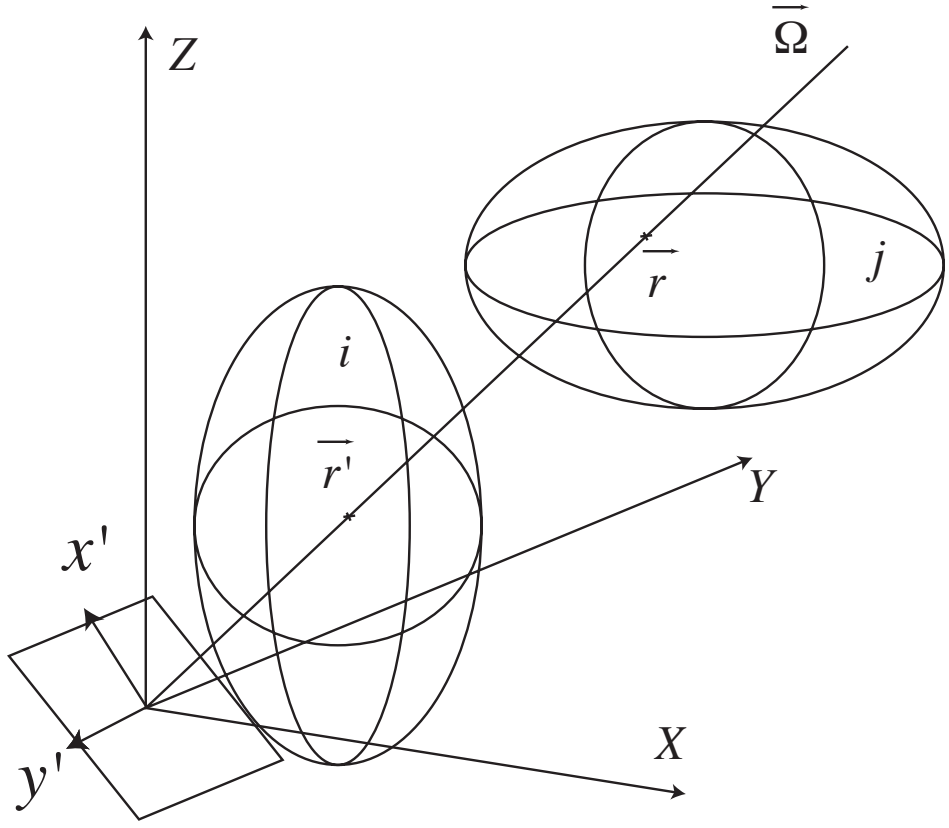
$$\int_{V_j} d^3r \Theta_j = \int_{4\pi} d^2\Omega \int_{R_{i-\frac{1}{2}}}^{R_{i+\frac{1}{2}}} R^2 dR \Theta_j$$

- Cartesian coordinates for d^3r' integral

$$\int_{V_i} d^3r' \Theta_i = \int dx' \int dy' \int dR' \Theta_i$$

Collision Probabilities in 3-D 2

General 3–D geometry for collision probability integration



Collision Probabilities in 3-D 3

Final form for collision probability integration

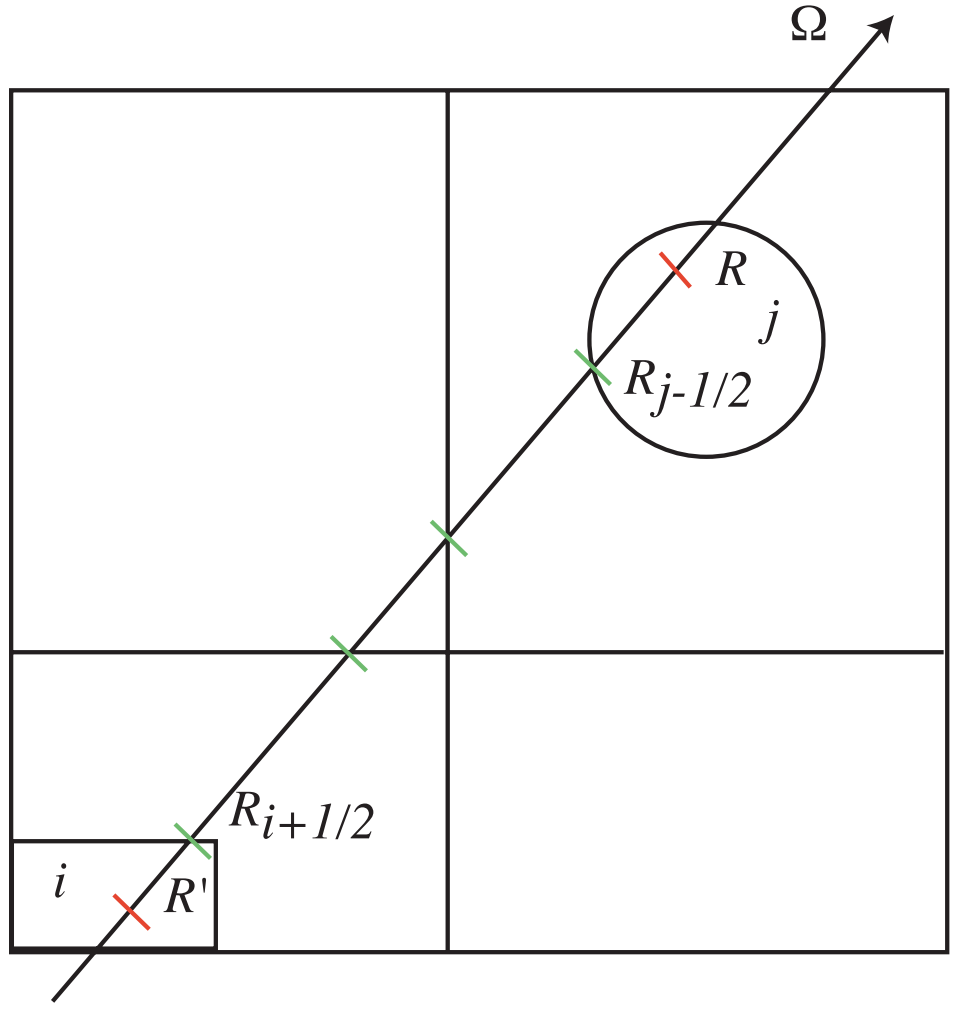
$$\tilde{p}_{ij}^{g} = \int_{4\pi} \frac{d^2\Omega}{4\pi} \int dx' \int dy' \int_{R_{i-\frac{1}{2}}}^{R_{i+\frac{1}{2}}} dR' \int_{R_{j-\frac{1}{2}}}^{R_{j+\frac{1}{2}}} dR \; e^{-\tau^g(R',R)} \Theta_i \Theta_j$$

$$\tau^g = (R_{i+\frac{1}{2}} - R')\Sigma_i^g + \sum_{k=i+1}^{j-1} \Delta R_k \Sigma_k^g + (R - R_{j-\frac{1}{2}})\Sigma_j^g$$

with $\Delta R_k = R_{k+\frac{1}{2}} - R_{k-\frac{1}{2}}$.

Collision Probabilities in 3-D 4

Notation for optical path



Collision Probabilities in 3-D

After integration over R' and R, one obtains

$$\begin{aligned}\tilde{p}_{ij}^{g} =& \frac{1}{4\pi\Sigma_i^g\Sigma_j^g}\int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i\Theta_j \\ & \times \left[1-\exp\left(-\tau^g_{i-\frac{1}{2},i+\frac{1}{2}}\right)\right]\exp\left(-\tau^g_{i+\frac{1}{2},j-\frac{1}{2}}\right) \\ & \times \left[1-\exp\left(-\tau^g_{j-\frac{1}{2},j+\frac{1}{2}}\right)\right]\end{aligned}$$

with

$$\tau^g_{i\pm\frac{1}{2},j\pm\frac{1}{2}} = \Sigma_i^g(R_{i+\frac{1}{2}} - R_{i\pm\frac{1}{2}}) + \tau_{i+\frac{1}{2},j-\frac{1}{2}} + \Sigma_j^g(R_{j\pm\frac{1}{2}} - R_{j-\frac{1}{2}})$$

Collision Probabilities in 3-D

Case where $\Sigma_i^g = 0$

$$\tilde{p}_{ij}^g = \frac{1}{4\pi\Sigma_j^g} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_j \times \Delta R_i \exp(-\tau^g_{i+\frac{1}{2},j-\frac{1}{2}}) \left[1 - \exp(-\tau^g_{j-\frac{1}{2},j+\frac{1}{2}})\right]$$

Case where $\Sigma_i^g = \Sigma_j^g = 0$

$$\tilde{p}_{ij}^g = \frac{1}{4\pi} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_j \Delta R_i \Delta R_j \exp(-\tau^g_{i+\frac{1}{2},j-\frac{1}{2}})$$

Collision Probabilities in 3-D

For $\tilde{p}^g_{ii}$, one obtains after integration over R' and R

$$\tilde{p}^g_{ii} = \frac{1}{2\pi \left(\Sigma^g_i\right)^2} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_i \times \left[\tau^g_{i-\frac{1}{2},i+\frac{1}{2}} - \left(1 - \exp(-\tau^g_{i-\frac{1}{2},i+\frac{1}{2}})\right)\right]$$

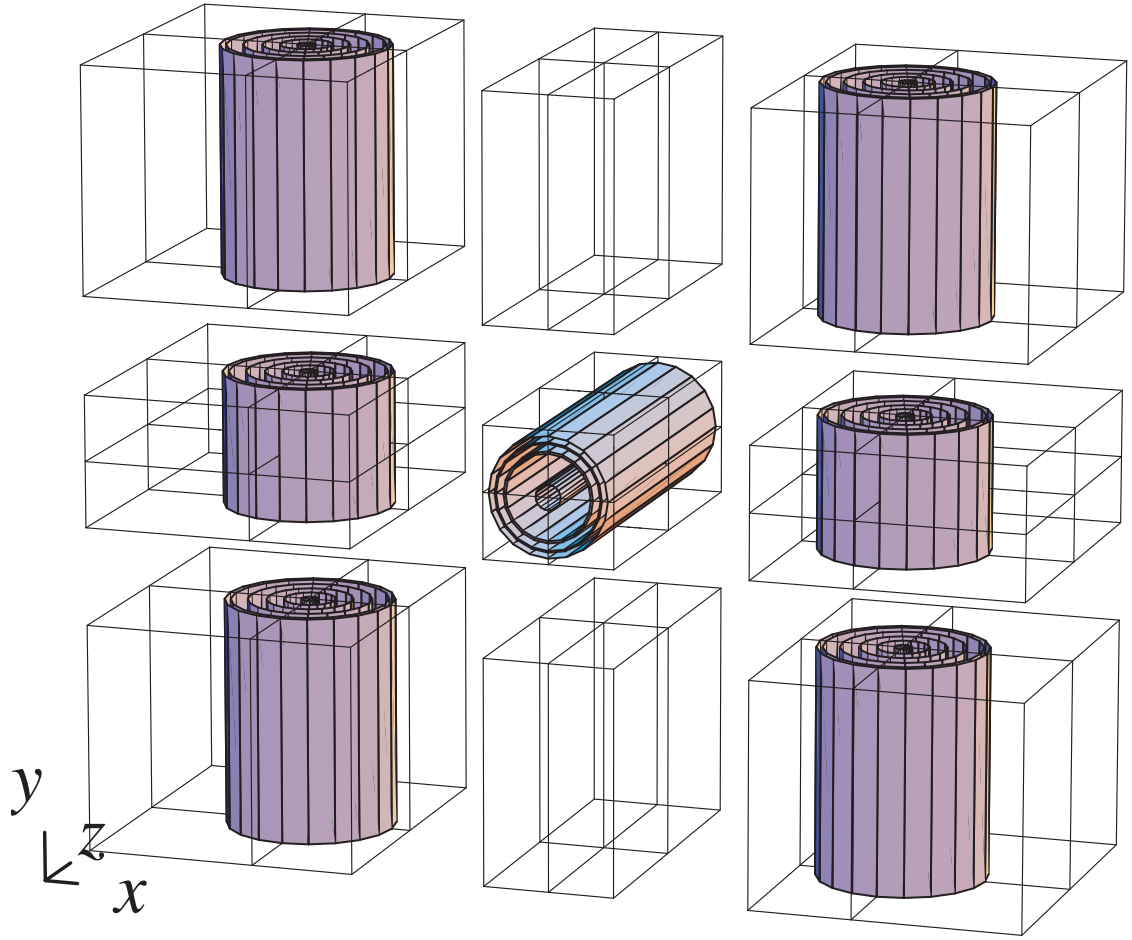
For $\Sigma^g_i = 0$ this is simplified to

$$\tilde{p}^g_{ii} = \frac{1}{4\pi} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_i \left(\Delta R_i\right)^2$$

Similar relations are obtained for $\tilde{p}^g_{i\alpha}$ and $\tilde{p}^g_{\alpha\beta}$

Quadrature and Tracking 1

3-D Adjuster model in DRAGON



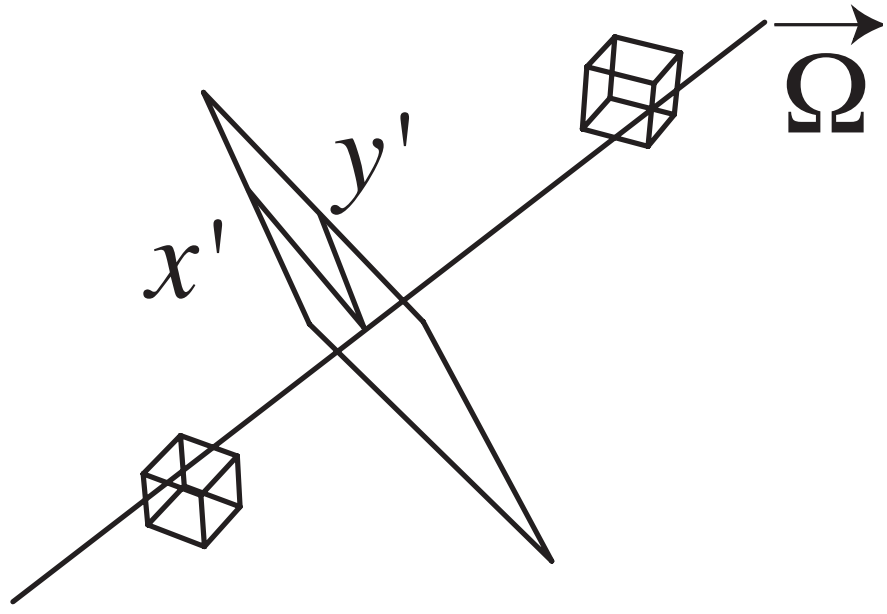
Quadrature and Tracking 2

Angular quadrature

- S_n type EQ$_N$ quadrature with $4N_\Omega(N_\Omega+2)/8$ angular directions $\vec{\Omega}_{1,i}$, $\vec{\Omega}_{2,i}$, $\vec{\Omega}_{3,i}$ and $\vec{\Omega}_{4,i}$
- Global quadrature weight $W_\Omega = 2/(N_\Omega(N_\Omega+2))$
- Tracking in lower half sphere only
- Number of tracking quadrant automatically reduced for symmetric cell
- Select as many angles as possible (Neutrons travel on a straight line)

Quadrature and Tracking 3

Cartesian surface quadrature



Quadrature and Tracking 4

Cartesian quadrature

- Identify the radius h_+ of the smallest sphere including the geometry
- Select a tracking density ρ_p and define the line spacing δ

$$\delta = \frac{2h_+}{N_p} \qquad N_p = (2\sqrt{\rho_p}h_+) + 1$$

- Integration line l_{mn} passes through

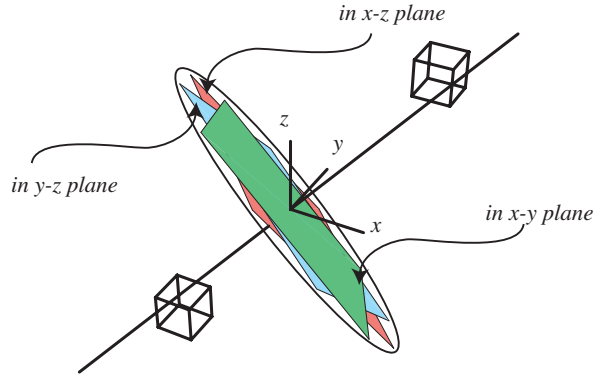
$$u_m^x = \left(\frac{2m-1}{2}\right)\delta \qquad u_n^y = \left(\frac{2n-1}{2}\right)\delta$$

- Integration weight $W_p = \delta^2$

Quadrature and Tracking 5

Comments on Cartesian surface quadrature

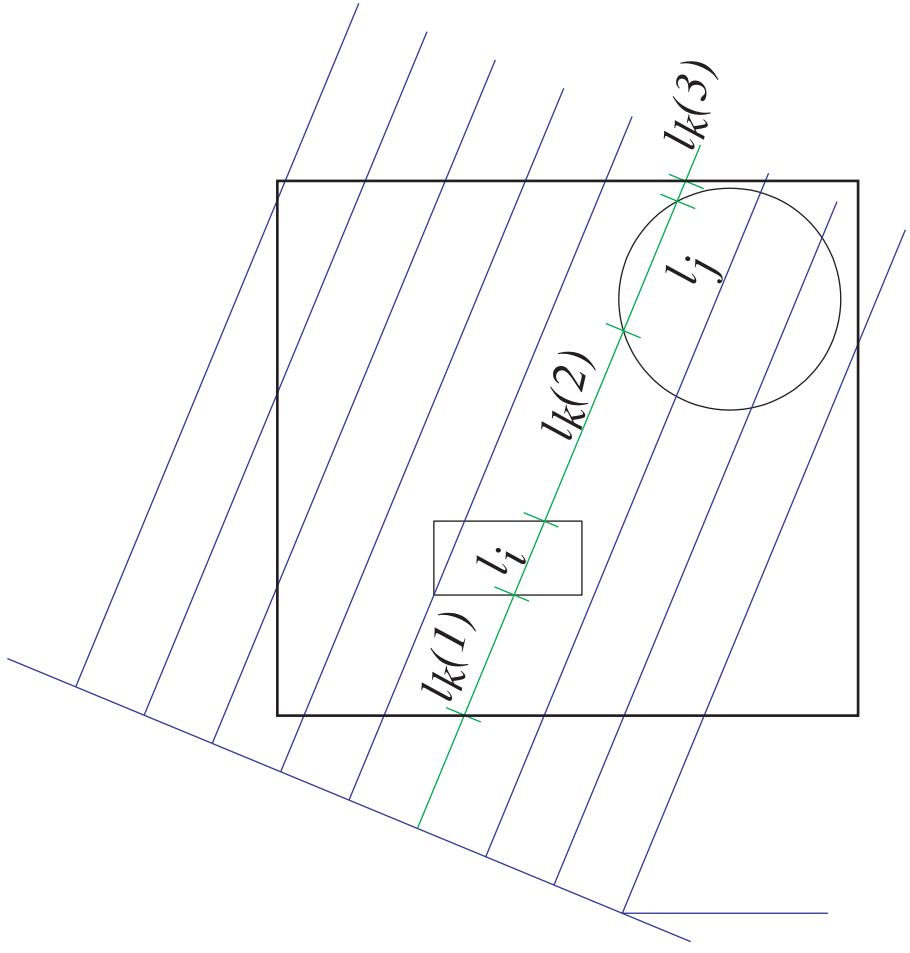
- In Dragon 3 different planes are selected for each given spatial direction



- Select tracking density as dense as possible
  - Each region must be touched by a maximum number of lines

Quadrature and Tracking 6

DRAGON tracking example



Quadrature and Tracking 7

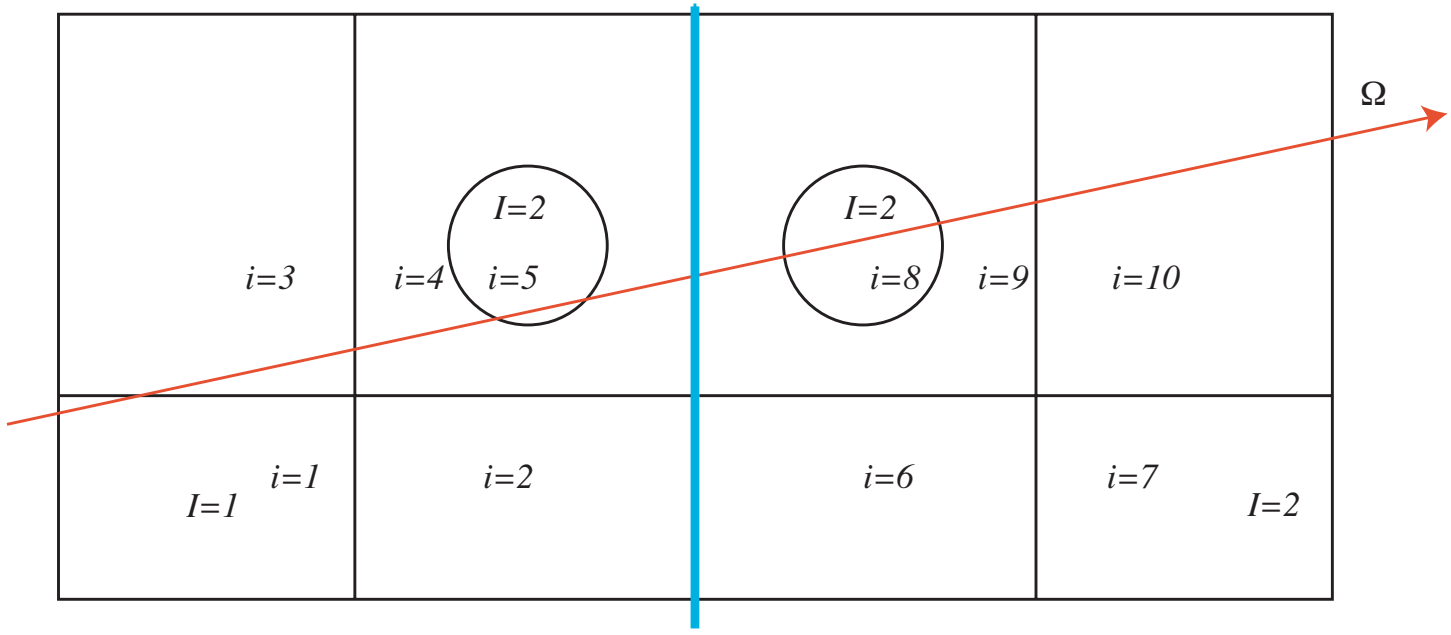
DRAGON tracking of a line $l_{m,n}$

1. Follow the tracking line as it travels through the cell
2. Identify each region i and surface s uniquely
3. Identify final region number I associated with set of regions i (a flux region)
4. Identify final surface number S associated with sets of surfaces s
5. Identify external surfaces and regions i crossed by line
6. Compute distance $\tilde{l}_{i,m,n}$ the neutron travels in each region
7. Store information on temporary tracking file

Quadrature and Tracking 8

Case of symmetric cells (mirror reflection on one side of the cell)

- Unfold the cell according to symmetry



- Regions in unfolded cells are originally assigned new region numbers

Quadrature and Tracking 9

Angle selection for symmetric cells

- For symmetry with respect to a (y, z) plane, track only in quadrants corresponding to directions $\vec{\Omega}_{1,i}$ and $\vec{\Omega}_{3,i}$
- For symmetry with respect to a (z, x) plane, track only in quadrants corresponding to directions $\vec{\Omega}_{1,i}$ and $\vec{\Omega}_{2,i}$
- For symmetry with respect to a (x, y) plane, track only in all quadrants

Quadrature and Tracking 10

Post treatment of tracking file

- Assign to all spatial region i its final flux region number combining track segments as required
- For each direction, normalize tracks using

$$l_{i,m,n} = \left(\frac{V_i}{\tilde{V}_i} \right) \tilde{l}_{i,m,n}$$

where

$$\tilde{V}_i = \frac{W_p}{3} \sum_{m=1}^{N_p} \sum_{n=1}^{N_p} \tilde{l}_{i,m,n}$$

Quadrature and Tracking 11

Comments on storage requirements for tracking

- Maximum number of line segments tracks d_t

$$d_t < 6N(N_\Omega(N_\Omega + 2))\rho_p h_+^2$$

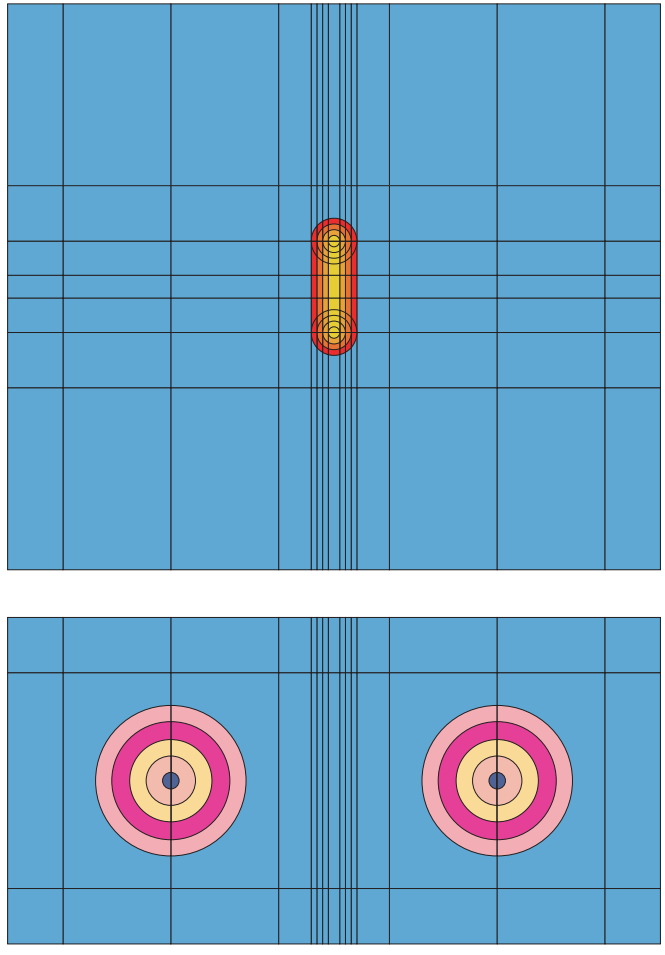
- For $N = 1000$, $h_+ = 50$, $\rho_p = 20$ t/cm^2 and $N_\Omega = 8$

$$d_t < 20 \times 10^9$$

tracks segments.

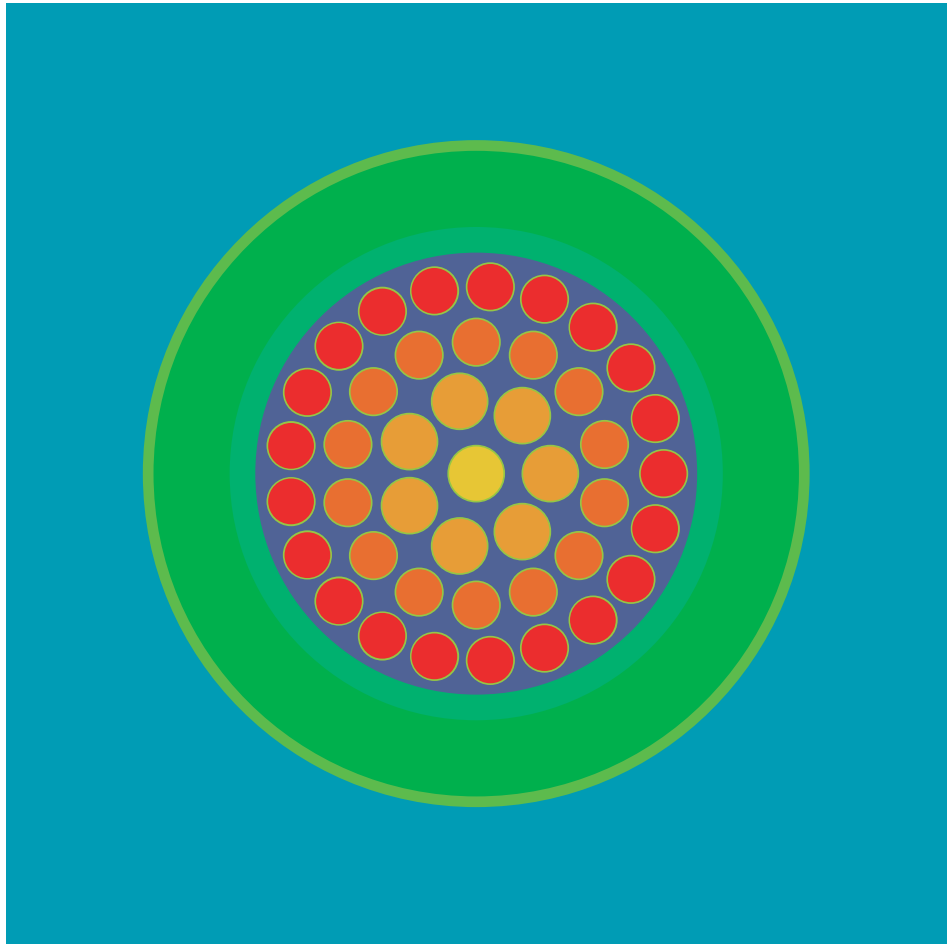
Quadrature and Tracking 12

ACR control rods model in DRAGON



Quadrature and Tracking 13

3-D cluster analysis not currently allowed in DRAGON



CP Integration 1

For each energy group

- Read a line from tracking file
- Scan this line and add contribution to $\tilde{p}_{ii}^{g}$

$$\frac{1}{2}(\Sigma_i^g)^2\tilde{p}_{ii}^{g} = \frac{1}{2}(\Sigma_i^g)^2\tilde{p}_{ii}^{g} + \sum_n W_n \sum_{m\in i}\left(\tau_{i,n,m}^{g} - \kappa_{i,n,m}^{g}\right)$$

where $W_n = W_\Omega W_p/3$ and

$$\kappa_{i,n,m}^{g} = \left(1 - \exp\left[-\tau_{i,n,m}^{g}\right]\right)$$

$$\tau_{i,n,m}^{g} = \Sigma_i^g l_{i,n,m} \qquad \text{with } m \in i$$

CP Integration 2

- Scan this line a second time and add contributions to $\tilde{p}^g_{ij}$

$$\Sigma^g_i \Sigma^g_j \tilde{p}^g_{ij} = \Sigma^g_i \Sigma^g_j \tilde{p}^g_{ij} + \sum_n W_n \sum_{m \in i} \sum_{m' \in j} \kappa^g_{i,n,m} \kappa^g_{n,m+1,m'-1} \kappa^g_{j,n,m'}$$

using

$$\kappa^g_{n,m,m'} = \prod_{l=m}^{m'} \exp\left[-\tau^g_{i,n,l}\right]$$

CP Integration 3

Finish CP calculations

- Only the contributions with $m < m'$ has been considered
- Symmetrize p_{ij}^g using

$$\tilde{p}_{ij}^g = \tilde{p}_{ij}^g + \tilde{p}_{ji}^g$$

CP Normalization 1

Compute errors on CP conservation rules

$$R_j^g = \Sigma_j^g V_j - \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} \Sigma_i^g p_{\alpha j}^g - \sum_{i=1}^{N_i} \Sigma_j^g \Sigma_i^g V_i p_{ij}^g$$

$$R_\beta^g = \frac{S_\beta}{4} - \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} p_{\alpha\beta}^g - \sum_{i=1}^{N_i} \Sigma_i^g V_i p_{i\beta}^g$$

CP Normalization 2

Diagonal Normalization

$$p^g_{D,ii} = p^g_{ii} - \frac{R^g_i}{(\Sigma^g_i)^2 V_i}$$

$$p^g_{D,\alpha\alpha} = p^g_{\alpha\alpha} - \frac{4R^g_\alpha}{S_\alpha}$$

- May result in non-physical negative probabilities
- Cannot be applied to problems involving voided zones

CP Normalization 3

HELIOS Type Normalization

$$p^g_{H,ij} = (w^g_i + w^g_j) p^g_{ij} \qquad p^g_{H,\alpha\alpha} = (w^g_\alpha + w^g_\beta) p^g_{\alpha\alpha}$$

- Apply conservation laws to above relation
- Solve resulting system for w^g using an iterative process
- Does not lead to negative probabilities and works for void regions
- Default option in DRAGON

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CP Normalization 4

Comments on storage requirements for CP matrices

- Number of elements per groups is N^2
- Memory space required for execution is about $5N^2$
- Total disk space required for storage of G group CP is GN^2

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Solving the CP Equations 1

Contents

- The Power Iteration
- The Multigroup Iteration
- Leakage Models

Solving the CP Equations 2

The multigroup transport equation has the form

$$\vec{\phi} = \mathbf{P}_{c,vv}(\vec{q}_s + \frac{1}{k}\vec{q}_f)$$

$$\vec{q}_s = \mathbf{\Sigma}_s\vec{\phi}$$

$$\vec{q}_f = \chi\nu\mathbf{\Sigma}_f\vec{\phi}$$

- $\vec{\phi}$ is a $N \times G$ dimensional vector
- $\mathbf{P}_{c,vv}$ is the multigroup CP matrix
  - Diagonal in energy, full in space.
- $\chi\nu\Sigma_f$ is a matrix for neutron production by fission and Σ_s is the scattering matrix
  - Diagonal in space, full in energy

Solving the CP Equations 3

We can decompose the scattering matrix as

$$\Sigma_s = \Sigma_{d,s} + \Sigma_{u,s} + \Sigma_{w,s}$$

with

- $\Sigma_{u,s}$ the up-scattering matrix (lower triangular in energy)
- $\Sigma_{d,s}$ the down-scattering matrix (upper triangular in energy)
- $\Sigma_{w,s}$ the within-group scattering matrix (diagonal in energy)

Solving the CP Equations 4

Defining $\mathbf{W}$ the scattering modified CP matrix

$$\mathbf{W} = (\mathbf{I} - \mathbf{P}_{c,vv}\mathbf{\Sigma}_{w,s})^{-1}\mathbf{P}_{c,vv}$$

the transport equation now becomes

$$\vec{\phi} = \mathbf{W}(\mathbf{\Sigma}_{d,s}\vec{\phi} + \mathbf{\Sigma}_{u,s}\vec{\phi} + \frac{1}{k}\vec{q}_f)$$

Assuming

- $\vec{q}_f/k$ is fixed and $\mathbf{\Sigma}_{w,s} = 0$

The above equation can be solved directly from group $g = 1$ to $g = G$

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Solving the CP Equations 5

The general equation can be solve using two iteration processes

- The Power iteration
  illustrated by solving the transport problem with $\Sigma_{w,s} = 0$
- The multigroup iteration
  illustrated by solving the transport problem with $\vec{q}_f/k$ fixed

The Power Iteration 1

- Assume $\Sigma_{u,s} = 0$
- Write an iterative group by broup solution to the transport problem as

$$\vec{\phi}(l) = \mathbf{W}(\mathbf{\Sigma}_{d,s}\vec{\phi}(l) + \frac{\chi}{k(l-1)}\nu\mathbf{\Sigma}_f\vec{\phi}(l-1))$$

$$k(l) = \sum_{g=1}^{G}\sum_{i=1}^{N} V_i\chi_i^g \sum_{h=1}^{G}\nu\Sigma_{f,i}^h\phi_i^h(l)$$

with $\vec{\phi}(0)$ a known arbitrary flux distribution and

$$k(0) = \sum_{g=1}^{G}\sum_{i=1}^{N} V_i\chi_i^g \sum_{h=1}^{G}\nu\Sigma_{f,i}^h\phi_i^h(0)$$

The Power Iteration 1

The iteration process is repeated until

$$k(l) - k(l-1) < \epsilon_1$$

$$\left| \frac{\vec{\phi}(l)}{k(l)} - \frac{\vec{\phi}(l-1)}{k(l-1)} \right| < \epsilon_2 \left| \frac{\vec{\phi}(l)}{k(l)} \right|$$

are both satisfied

- The parameters ϵ_1 and ϵ_2 can be defined independently in DRAGON

The Multigroup Iteration 1

- Assume $\vec{q}_f/k = \vec{q}$ is fixed
- Solve group-by-group this fixed source problem using a Gauss-Seidel strategy

$$\vec{\phi}(l) = \mathbf{W}\left(\mathbf{\Sigma}_{d,s}\vec{\phi}(l) + \mathbf{\Sigma}_{u,s}\vec{\phi}(l-1) + \vec{q}\right)$$

- Iterate until

$$\left|\frac{\vec{\phi}^g(l) - \vec{\phi}^g(l-1)}{\vec{\phi}^g(l)}\right| < \epsilon_3$$

The Multigroup Iteration 2

Multigroup rebalancing technique

- Neutron conservation states that for a converged solution $\phi_i^g(l)$

$$\sum_{i=1}^{N} \Sigma_i^g V_i \phi_i^g(l) = \sum_{i=1}^{N} R_i^g V_i q_i^g(l)$$

$$+ \sum_{i=1}^{N} R_i^g V_i \left(\sum_{h=1}^{g} \Sigma_s^{h \to g} \phi_i^h(l) + \sum_{h=g+1}^{G} \Sigma_s^{h \to g} \phi_i^h(l) \right)$$

with $R_i^g = 1 - \sum_{i=1}^{N} \Sigma_j^g p_{c,ij}^g$

The Multigroup Iteration 3

- The multigroup problem we effectively solve yields

$$\sum_{i=1}^{N} \Sigma_i^g V_i \phi_i^g(l) = \sum_{i=1}^{N} R_i^g V_i q_i^g(l)$$

$$+ \sum_{i=1}^{N} R_i^g V_i \left(\sum_{h=1}^{g} \Sigma_s^{h \to g} \phi_i^h(l) + \sum_{h=g+1}^{G} \Sigma_s^{h \to g} \phi_i^h(l-1) \right)$$

The Multigroup Iteration 4

To restore conservation at each iteration

- Use $\tilde{\phi}_i^g = \alpha^g \phi_i^g$ and assume $\tilde{\phi}_i^g$ satisfies conservation relations, then α^g must satisfy

$$\sum_{h=1}^{G} M^{h\to g}\alpha^h = q^g$$

$$M^{h\to g} = \sum_{i=1}^{N} R_i^g V_i \left(\Sigma_i^h \delta_{gh} - \Sigma_s^{h\to g}\right) \phi_i^h$$

$$q^g = \sum_{i=1}^{N} R_i^g V_i q_i^g$$

- Solve for $\alpha^g(k)$ and rebalance flux

The Multigroup Iteration 5

Add relaxation parameter to the Gauss–Seidel iteration scheme

- For appromimate solution $\vec{\Gamma}(l)$

$$\vec{\Gamma}(l) = \mathbf{W}\left(\mathbf{\Sigma}_{d,s}\vec{\Gamma}(l) + \mathbf{\Sigma}_{u,s}\vec{\Gamma}(l-1) + \vec{q}\right)$$

- Define an improved flux distribution for the next iteration using

$$\vec{\phi}(l) = \vec{\Gamma}(l) + \omega(l)\vec{\Delta}(l)$$
$$\vec{\Delta}(l) = \vec{\Gamma}(l) - \vec{\phi}(l-1)$$

and $\omega(l)$ will be computed using a variational procedure

The Multigroup Iteration 6

Select $\omega(l)$ in such a way that $\vec{\phi}(l)$ minimizes the tranport functional

$$\mathcal{F}[\vec{\phi}] = \frac{1}{2}\vec{\phi}^T\mathbf{Z}^T\mathbf{Z}\vec{\phi} - \vec{\phi}^T\mathbf{Z}^T\mathbf{W}\vec{q}$$

$$\mathbf{Z} = [\mathbf{I} - \mathbf{W}\left(\mathbf{\Sigma}_{d,s} + \mathbf{\Sigma}_{u,s}\right)]$$

This yield

$$\omega(l) = -\frac{[\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]^T[\vec{\Gamma} - \mathbf{W}\vec{S}_2(l)]}{[\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]^T[\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]}$$

$$\vec{S}_1(l) = (\mathbf{\Sigma}_{d,s} + \mathbf{\Sigma}_{u,s})\vec{\Delta}(l)$$

$$\vec{S}_2(l) = (\vec{q} + \mathbf{\Sigma}_{d,s}\vec{\Gamma}(l) + \mathbf{\Sigma}_{u,s}\vec{\Gamma}(l))$$

Leakage Models

In 3–D, the transport equation can be solved in DRAGON using the B_0 and B_1 leakage models

- Both of these models are based on the following factorization of the flux

$$\Phi^g(\vec{r},\vec{\Omega}) \approx \Psi^g(\vec{r},\vec{\Omega})\exp(i\vec{B}\cdot\vec{r})$$

- Transport equation with leakage

$$\vec{\Omega}\cdot\vec{\nabla}\Psi^g(\vec{r},\vec{\Omega}) + [\Sigma^g(\vec{r}) + i\vec{B}\cdot\vec{\Omega}]\Psi^g(\vec{r},\vec{\Omega}) =$$

$$Q_s^g(\vec{r},\vec{\Omega}) + \frac{1}{k_{eff}}Q_f^g(\vec{r})$$

- In general we will assume that k_{eff}=1

Leakage Models 2

- For an infinite homogeneous media the scalar flux and vector current are related to each other according to

$$\vec{\Omega}\Psi^g(\vec{\Omega}) = \vec{J}^g(\vec{\Omega}) = -iD^g\vec{B}\Psi^g(\vec{\Omega})$$

  with D^g is an homogeneous diffusion coefficient

- Apply to heterogeneous systems

$$\vec{\Omega}\cdot\vec{\nabla}\Psi^g(\vec{r},\vec{\Omega}) + [\Sigma^g(\vec{r}) + D^g B^2]\Psi^g(\vec{r},\vec{\Omega}) = Q_s^g(\vec{r},\vec{\Omega}) + Q_f^g(\vec{r})$$

- Find the homogeneous diffusion coefficient compatible with this heterogeneous problem

Leakage Models 3

- Assume an heterogeneous solution is known for $B^2 = 0$
- Use this solution to define an equivalent infinite homogeneous problem

$$\Sigma^g \Psi^g(\vec{\Omega}) + i\vec{B} \cdot \vec{J}^g(\vec{\Omega}) = Q_s^g(\vec{\Omega}) + Q_f^g$$

  where the cross sections and sources are homogenized using the heterogeneous flux
- Solve the homogeneous problem for D^g and B
- Insert in heterogeneous transport equation and obtain an improved solution.
- Repeat until the iterative procedure is converged

Leakage Models

Solving the homogeneous problem (B_1 model)

- Use a 2 terms expansion for the scattering cross section in Legendre polynomials

$$\Sigma_s^{h\to g}(\vec{\Omega}' \to \vec{\Omega}) = \Sigma_{s,0}^{h\to g} + 3\Sigma_{s,1}^{h\to g}\vec{\Omega}\cdot\vec{\Omega}'$$

- Define

$$\psi^g = \int d^2\Omega \Psi^g(\vec{\Omega})$$

$$\vec{j}^g = \int d^2\Omega \vec{\Omega}' \Psi^g(\vec{\Omega})$$

Leakage Models

- Insert into the homogeneous transport equation, and integrate to obtain

$$\psi^g = \alpha^g \sum_h (\Sigma_{s,0}^{h\to g} + \chi^g \nu\Sigma_f^h)\psi^h + 3\beta^g \sum_h \Sigma_{s,1}^{h\to g} \frac{\vec{B}\cdot\vec{j}^h}{iB^2}$$

$$\vec{j}^g = \beta^g \sum_h \left[(\Sigma_{s,0}^{h\to g} + \chi^g \nu\Sigma_f^h) \frac{\vec{B}\psi^h}{iB^2} + 3\Sigma^g \Sigma_{s,1}^{h\to g} \frac{\vec{j}^h}{B^2} \right]$$

$$\alpha^g = \frac{1}{B} \arctan\left(\frac{B}{\Sigma^g}\right)$$

$$\beta^g = 1 - \Sigma^g \alpha^g$$

Solve for B, ψ^g and $\vec{j}^g$ and compute $D^g = i\vec{B}\cdot\vec{j}^g / B^2\psi^g$

Leakage Models

The B_0 homogeneous problem

- Assume $\Sigma_{s,1}^{h\to g} = 0$ and obtain

$$\psi^g = \alpha^g \sum_h (\Sigma_{s,0}^{h\to g} + \chi^g \nu\Sigma_f^h)\psi^h$$

$$\vec{j}^g = \beta^g \frac{\vec{B}}{iB^2} \sum_h (\Sigma_{s,0}^{h\to g} + \chi^g \nu\Sigma_f^h)\psi^h$$

The homogeneous diffusion coefficient is then given by

$$D^g = \frac{\beta^g}{\alpha^g}$$

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Condensation and Homogenization 1

Contents

- Condensation Technique
- Full Cell Homogenization
- Partial Cell Homogenization and SPH Factors
- Microscopic Cross Section Homogenization

Condensation and Homogenization 2

Condensation and homogenization techniques in DRAGON are based on the following assumptions

- Reaction rates are physically meaningfull and should be preserved by the condensation/homogenization procedure

$$R_i = \sum_g V_i \phi_i^g \Sigma_i^g = V_i \phi_i \Sigma_i$$

$$R^g = \sum_i V_i \phi_i^g \Sigma_i^g = V \phi^g \Sigma^g$$

- The eigenvalue is physically meaningfull and should be preserved by the condensation/homogenization procedure

Condensation Technique 1

Condensed transport equation (macrogroup K that includes $g \in G_K$)

$$\vec{\Omega} \cdot \vec{\nabla} \sum_{g \in G_K} \Phi^g(\vec{r}, \vec{\Omega}) + \sum_{g \in G_K} \Sigma^g(\vec{r}) \Phi^g(\vec{r}, \vec{\Omega}) =$$

$$\sum_{g \in G_K} [Q_s^g(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q_f^g(\vec{r})]$$

Few group version of the same equation is

$$\vec{\Omega} \cdot \vec{\nabla} \Phi^K(\vec{r}, \vec{\Omega}) + \Sigma^K(\vec{r}) \Phi^K(\vec{r}, \vec{\Omega}) = [Q_s^K(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q_f^K(\vec{r}, \vec{\Omega})]$$

It should reproduce condensed multigroup results

Condensation Technique 2

The condensation procedure that satisfies our requirements

$$\phi_i^K = \sum_{g \in G_K} \phi_i^g$$

$$\Sigma_i^K = \frac{1}{\phi_i^K} \sum_{g \in G_K} \Sigma_i^g \phi_i^g$$

$$\Sigma_{s,i}^{L \to K} = \frac{1}{\phi_i^L} \sum_{h \in G_L} \sum_{g \in G_K} \Sigma_{s,i}^{h \to g} \phi_i^h$$

$$\chi_i^K = \sum_{g \in G_K} \chi_i^g$$

$$\nu\Sigma_{f,i}^K = \frac{1}{\phi_i^K} \sum_{g \in G_K} \nu\Sigma_{f,i}^h \phi_i^g$$

Condensation Technique 3

Multiplying CP transport equation by $\Sigma_i^g V_i$ and summing over all regions i yields

$$\sum_i \Sigma_i^g V_i \phi_i^g = \sum_i V_i [Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]$$

The equivalent transport equation in a homogeneous infinite cell is

$$\hat{\Sigma}^g V \hat{\phi}^g = V [\hat{Q}_{s,i}^g + \frac{1}{k} \hat{Q}_{f,i}^g]$$

The homogenized and homogeneous transport equations are identical if one selects a flux-volume homogenization technique

Condensation Technique 4

Flux-volume homogenization technique

$$\hat{\phi}^g = \frac{1}{V}\sum_i V_i \phi_i^g$$

$$\hat{\Sigma}^g = \frac{1}{V\hat{\phi}^g}\sum_i V_i \Sigma_i^g \phi_i^g$$

$$\hat{\Sigma}_s^{h\to g} = \frac{1}{V\hat{\phi}^h}\sum_i V_i \Sigma_{s,i}^{h\to g} \phi_i^h$$

$$\hat{\nu}\hat{\Sigma}_f^g = \frac{1}{V\hat{\phi}^g}\sum_i V_i \nu\Sigma_{f,i}^g \phi_i^g$$

$$\hat{\chi}^g = \frac{1}{V\sum_h \hat{\nu}\hat{\Sigma}_f^g \hat{\phi}^g}\sum_i \chi_i^g V_i \sum_h \nu\Sigma_{f,i}^g \phi_i^h$$

Condensation Technique 5

Flux-volume homogenization fails if

- The cell is finite (a cell with leakage) and

$$\sum_{j=1}^{N_j} p_{ij}^g \Sigma_j^g \neq 1$$

- Partial cell homogenization cell is considered

Partial Cell Homogenization 1

The heterogeneous N region transport equation homogenized over M regions takes the form

$$\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = \sum_{i \in M_I} \sum_J \sum_{j \in M_J} p_{ji}^g(\Sigma^g)[Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]$$

The M region heterogeneous transport equation takes the from

$$V_I \hat{\Sigma}_I^g \hat{\phi}_I^g = \sum_J \hat{p}_{JI}^g(\hat{\Sigma}^g)[Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g]$$

where $\hat{P}_{JI}^g(\hat{\Sigma}^g)$ indicates that the CP are computed using homogenized cross sections

Partial Cell Homogenization 2

We need

$$\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = V_I \Sigma_I^g \phi_I^g$$

and

$$\sum_J \hat{p}_{JI}^g(\hat{\Sigma}^g)[Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g] =$$

$$\sum_{i \in M_I} \sum_J \sum_{j \in M_J} p_{ji}^g(\Sigma^g)[Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]$$

to be simultaneously true

Partial Cell Homogenization 3

The flux-volume homogenization method is not longer adequate because

- There is no simple relation between $\hat{p}^g_{JI}(\hat{\Sigma}^g)$ and $p^g_{ji}(\Sigma^g)$

The alternative here is to use a non-linear process

- Consider a flux-volume homogenization for ϕ^g_I and Σ^g_I
- Redefine the homogeneous flux $\hat{\phi}^g_I$ and cross sections $\hat{\Sigma}^g_I$ as follows

$$\hat{\phi}^g_I = \frac{1}{\mu^g_I}\phi^g_I \qquad \hat{\Sigma}^g_I = \mu^g_I \Sigma^g_I$$

Partial Cell Homogenization 4

- Determine the SPH factors μ_I^g numerically in such a way that

$$\sum_J \hat{p}_{JI}^g(\hat{\Sigma}^g)[Q_{s,J}^g + \frac{1}{k}Q_{f,J}^g] =$$

$$\sum_{i\in M_I}\sum_J\sum_{j\in M_J} p_{ji}^g(\Sigma^g)[Q_{s,i}^g + \frac{1}{k}Q_{f,i}^g]$$

  is true

- The definition of the SPH factors automatically ensures

$$\sum_{i\in M_I} V_i\Sigma_i^g\phi_i^g = V_I\Sigma_I^g\phi_I^g = V_I\tilde{\Sigma}_I^g\tilde{\phi}_I^g$$

Microscopic Cross Section 1

The macroscopic cross section associated with a material is simply the sum over all isotopes of the isotopic macroscopic cross section Σ_I namely

$$\Sigma_i^g = \sum_I \Sigma_{I,i}^g$$

where

$$\Sigma_{I,i}^g = N_{I,i}\sigma_I^g$$

with $N_{I,i}$, the concentration of isotope I in region i

- The homogenization and condensation procedure described above remain valid for $\Sigma_{I,i}^g$

Microscopic Cross Section 2

Since the final concentration of isotope I in the cell is given by:

$$N_I = \frac{1}{V} \sum_i N_{I,i} V_i$$

we can define the equivalent homogenized microscopic cross section as:

$$\hat{\sigma}_I^K = \frac{\mu_I^K}{N_I V \phi^K} \sum_{i \in M_I} \sum_{g \in G_K} N_{I,i} V_i \sigma_I^g \phi_i^g$$

where the microscopic cross sections now become dependent on the spatial position

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Conclusions 1

Some comments and warning on the CP method

- The sources are assumed constant inside each region
  - Select an adequate spatial discretization
    This may lead to a large number of region (CP is proportionnal to N^2)
    Some regions may be very small causing problem with tracking
- Select a problem that is not too heterogeneous

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Conclusions 2

- The angular flux on each external surface are assumed constant and isotropic
  - Try to get rid of external surfaces with re-entrant angular flux
  - Select a model where the region of interest is far from the external surfaces