

## **DRAGON Theory for 3-D CANDU Problems**

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- 1. Introduction to the Collision Probability Method
- 2. 3–D Collision Probability Calculations
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- 4. Solving the Collision Probability Equations
- 5. Condensation and Homogenization Techniques
- 6. Managing <sup>a</sup> DRAGON Execution
- 7. Discussion and Conclusion

#### **Introduction to the CP Method 1** ÉCOLE

## **Contents**

- **C** The Transport Equation
- **C** The Collision Probability Technique
- **C** Boundary Conditions
- Cross Sections Considerations

### **The Transport Equation 1** ③}<br>◎ ÉCOLE<br><mark>DLYTECHNIQUE</mark>

## The transport equation is <sup>a</sup> neutron balance equation

$$
\mathcal{L}(\vec{r},E,\vec{\Omega})=\!\mathcal{Q}(\vec{r},E,\vec{\Omega})
$$

 $\mathcal L$  represents neutron lost from the system:

$$
\mathcal{L} = \vec{\Omega} \cdot \vec{\nabla} \Phi(\vec{r}, E, \vec{\Omega}) + \Sigma(\vec{r}, E) \Phi(\vec{r}, E, \vec{\Omega})
$$

#### Q represents neutron created in the system  $\bullet$

$$
Q_s = \int dE' d^2\Omega \Sigma_s(\vec{r}, E' \to E, \vec{\Omega}' \to \vec{\Omega}) \Phi(\vec{r}, E', \vec{\Omega}')
$$
  

$$
Q_f = \chi(\vec{r}, E) \int dE' d^2\Omega \nu \Sigma_f(\vec{r}, E') \Phi(\vec{r}, E', \vec{\Omega}')
$$

#### **The Transport Equation 2 O**<br>SE ÉCOLE<br>POLYTECHNIQUE

Multigroup transport equation

$$
\left[\vec{\Omega}\cdot\vec{\nabla}+\Sigma^g(\vec{r})+D^g(\vec{r})B^2\right]\Phi^g(\vec{r},\vec{\Omega})=Q_s^g(\vec{r},\vec{\Omega})+\frac{1}{k}Q_f^g(\vec{r},\vec{\Omega})
$$

Scattering source  $Q_s^g$  $_{s}^{g}(\vec{r}% )\sim\vec{r}^{2}\sqrt{\vec{r}^{2}-\vec{r}^{2}^{2}}$  $\vec{r}, \vec{\Omega})$ 

$$
Q_s^g(\vec{r}, \vec{\Omega}) = \sum_{h=1}^G \int d^2\Omega' \Sigma_s^{hg}(\vec{r}, \vec{\Omega}' \to \vec{\Omega}) \Phi^h(\vec{r}, \vec{\Omega}')
$$

Fission source  $Q^g_{\,\,\mu}$  $\frac{g}{f}(\vec{r})$  $\vec{r}, \vec{\Omega})$ 

$$
Q_f^g(\vec{r}) = \chi^g \sum_{h=1}^G \nu \Sigma_f^h(\vec{r}) \int d^2\Omega' \Phi^h(\vec{r}, \vec{\Omega}')
$$

### **The Transport Equation 3** (<mark>⑥</mark>)<br>POLYTECHNIQUE<br>M. Q. N. T. P. É. A. L

Transport equation in the absence of external sources is an eigenvalue problem:

- 1.  $k$  eigenvalue with imposed leakage  $(D^g(\vec{r})B)$ 2 fixed):
	- $k$  indicates how the fission rate should be modified to make the system critical (reach <sup>a</sup> non-trivial solution to the transport equation)
- 2. Buckling eigenvalue with imposed  $k$ :
	- $D^g(\vec{r})B$  $^2$  represent the amount of leakage required to make the system critical

#### **The Transport Equation 4** ୍ଞ୍ୟୁ<br>ାଧ≺ ÉCOLE<br><mark>OLYTECHNIQUE</mark>

## Boundary conditions are required to close the system Albedo conditions

$$
\phi_-(\vec{r}_S,\vec{\Omega}-2(\vec{N}_S\cdot\vec{\Omega}))=\beta(\vec{r}_S,\vec{\Omega})\phi_+(\vec{r}_S,\vec{\Omega})
$$



 $\beta(\vec{r}_S,\vec{\Omega})=0$  $\Omega)=0$  for void BC  $\beta(\vec{r}_S,\vec{\Omega})=1$  $\Omega) = 1$  for reflective BC



### Periodic conditions

$$
\phi_-(\vec{r}_S,\vec{\Omega})=\phi_+(\vec{r}_{S'},\vec{\Omega})
$$



### **The Transport Equation 6** <u>®</u><br>◎ ÉCOLE<br><mark>OLYTECHNIQUE</mark>



Integral transport equation (case without leakage)

Flux at a point  $\vec{r}$  due to neutrons created at any point  $\vec{r}^{\, \prime} = \vec{r} - R \vec{\Omega}$  surrounding it

$$
\left[-\frac{d}{dR} + \Sigma^g(\vec{r} - R\vec{\Omega})\right] \Phi^g(\vec{r} - R\vec{\Omega}, \vec{\Omega}) = Q^g(\vec{r} - R\vec{\Omega}, \vec{\Omega})
$$

Integrate transport equation over  $R$  and  $\vec{\Omega}$ 

$$
\phi^g(\vec{r}) = \int d^2\Omega \ e^{-\tau^g(R_S)} (\vec{\Omega} \cdot \vec{N}_-) \Phi_-^g(\vec{r}'_S, \vec{\Omega}) \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega}) + \int d^2\Omega \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \vec{\Omega}) \Theta(\vec{r}, \vec{r}', \vec{\Omega}) dR'
$$

#### **The Transport Equation 7 O**<br>POLYTECHNIQUE<br>MONTRÉAL

### De finitions

$$
\phi^g(\vec{r}) = \int d^2\Omega \Phi^g(\vec{r}, \vec{\Omega})
$$

$$
\tau^g(R) = \int_O^R \Sigma^g(\vec{r} - R'\vec{\Omega})dR'
$$

$$
\Theta(\vec{r}, \vec{r}', \vec{\Omega}) = \begin{cases} 1 & \text{if } \vec{r} = \vec{r}' + R'\vec{\Omega} \\ 0 & \text{otherwise} \end{cases}
$$

and Φ g  $\frac{g}{\Gamma}(\vec{r}^{\, \prime}_S,\vec{\Omega})$ is the incoming angular flux on surface  $S$ 

### **The Transport Equation O**<br>
ECOLE<br>
POLYTECHNIQUE



Equation for the outgoing flux  $\phi_+(\vec{r}_S)$  at  $S$ 

$$
\phi^g_+(\vec{r}_S) = \int d^2\Omega(\vec{\Omega} \cdot \vec{N}_+) \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \vec{\Omega}) \Theta(\vec{r}_S, \vec{r}', \vec{\Omega}) dR'
$$

$$
+ \int d^2\Omega(\vec{\Omega} \cdot \vec{N}_+)(\vec{\Omega} \cdot \vec{N}_-) e^{-\tau^g(R_S)} \Phi^g_-(\vec{r}'_S, \vec{\Omega}) \Theta(\vec{r}_S, \vec{r}'_S, \vec{\Omega})
$$

where

$$
\phi^g_+(\vec{r}_S) = \int d^2\Omega \ \Phi^g_+(\vec{r}_S, \vec{\Omega}) (\vec{\Omega} \cdot \vec{N}_+)
$$

# **The Transport Equation 9**

- CP approximations
	- Divide domain into  $N_V$  regions of volume  $V_i$  where the cross sections sources are independent of  $\vec{r}$  and  $\vec{\Omega}$

$$
\Sigma^{g}(\vec{r}) = \Sigma^{g}_{j} \qquad \text{for } \vec{r} \in V_{j}
$$
\n
$$
Q^{g}(\vec{r}, \vec{\Omega}) = \frac{1}{4\pi} q^{g}(\vec{r}) = q^{g}_{j} \qquad \text{for } \vec{r} \in V_{j}
$$

(the source is assumed isotropic)

Divide the external boundary  $S$  into  $N_S$  surfaces of area  $S_\alpha$  and assume angular flux constant on these surfaces

$$
\Phi^g_-(\vec{r}_S,\vec{\Omega})=\frac{1}{4\pi}\phi^g_{\alpha,-}
$$

#### **110 The Transport Equation O**<br>
ECOLE<br>
POLYTECHNIQUE MONTRÉAL

 $\alpha,$ —

$$
\int d^2\Omega \int_0^R e^{-\tau^g(R')} Q^g(\vec{r}', \vec{\Omega}) \Theta(\vec{r}, \vec{r}', \vec{\Omega}) dR' =
$$
  

$$
q_j^g \int_{V_i} e^{-\tau^g(R')} \Theta(\vec{r}, \vec{r}', \vec{\Omega}) \frac{d^3r}{4\pi |\vec{r} - \vec{r}'|^2}
$$
  

$$
\int d^2\Omega \ e^{-\tau^g(R_S)} \Phi_-^g(\vec{r}'_S, \vec{\Omega}) \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega}) =
$$
  

$$
\phi_{\alpha,-}^g \int_{S_{\alpha}} e^{-\tau^g(R_S)} \Theta(\vec{r}, \vec{r}'_S, \vec{\Omega}) \frac{d^2r}{4\pi |\vec{r} - \vec{r}'|^2}
$$

 $-|\vec{r}'|^2$ |



### Transport equations in CP form



### where

$$
\phi_i^g = \frac{1}{V_i} \int_{V_i} d^3r \phi^g(\vec{r})
$$

$$
\phi_{+,\alpha}^g = \frac{1}{S_{\alpha}} \int_{S_{\alpha}} d^2r \phi^g(\vec{r}_S)
$$



Four types of probabilities

$$
\tilde{p}_{ij}^g = V_i p_{ij}^g = \int_{V_i} \int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2} \Theta_i \Theta_j d^3 r' d^3 r
$$
\n
$$
\tilde{p}_{i\alpha}^g = V_i p_{i\alpha}^g = \int_{V_i} \int_{S_\alpha} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2} (\vec{\Omega} \cdot \vec{N}_-) \Theta_i \Theta_\alpha d^3 r' d^2 r
$$
\n
$$
\tilde{p}_{\alpha i}^g = \frac{S_\alpha}{4} p_{\alpha i}^g = \int_{S_\alpha} \int_{V_i} \frac{e^{-\tau^g(R)}}{4\pi R^2} (\vec{\Omega} \cdot \vec{N}_+) \Theta_\alpha \Theta_i d^2 r' d^3 r
$$
\n
$$
\tilde{p}_{\alpha \beta}^g = \frac{S_\alpha}{4} p_{\alpha \beta}^g = \int_{S_\alpha} \int_{S_\beta} \frac{e^{-\tau^g(R_S)}}{4\pi R_S^2} (\vec{\Omega} \cdot \vec{N}_-) (\vec{\Omega} \cdot \vec{N}_+)
$$
\n
$$
\times \Theta_\alpha \Theta_\beta d^2 r d^2 r'
$$



### Symmetry relations

$$
V_i p_{ij}^g = V_j p_{ji}^g
$$

$$
4V_i p_{i\alpha}^g = S_{\alpha} p_{\alpha i}^g
$$

$$
S_{\alpha} p_{\alpha \beta}^g = S_{\beta} p_{\beta \alpha}^g
$$

Conservation properties

$$
\sum_{\alpha=1}^{N_{\alpha}} p_{i\alpha}^{g} + \sum_{j=1}^{N_{j}} p_{ij}^{g} \Sigma_{j}^{g} = 1
$$

$$
\sum_{\beta=1}^{N_{\beta}} p_{\alpha\beta}^{g} + \sum_{i=1}^{N_{i}} p_{\alpha i}^{g} \Sigma_{i} = 1
$$

## **Boundary Conditions**

Surface flux approximation:

Incoming angular flux on outer surfaces assumed to be independent of  $\vec{\Omega}$ 

**Comments** 

- Outgoing angular flux on outer surfaces integrated over  $\vec{\Omega}$
- Angular flux not used at region interfaces
- Approximation for incoming angular flux exact for vacuum BC
- Approximation for incoming angular flux leads to large errors in surface flux when flux is not isotropic



## Illustration of approximate BC

*specular isotropic*

*Reflection Translation*

*isotropic*

#### **Boundary Conditions 3** <sup>≠</sup><br>® ÉCOLE<br>⁄**TECHNIQUE**

Recommendations to reduce errors due to approximate use of BC for a fixed direction

- No special treatment for 2 vacuum BC
- Unfold cell once for 1 vacuum and 1 reflection BC
- Multiply unfold cell for 2 re flection or periodic cell, apply approximate BC on final surfaces and consider results in cell located far from these surfaces.



### Example of cell unfolding in direction  $X$





#### **Boundary Conditions 5 O**<br>
ECOLE<br>
POLYTECHNIQUE

## Simplifying CP equations using approximate BC **C** Assume

$$
\vec{J}^g_- = \mathbf{A}^g \vec{J}^g_+
$$

**Final transport equation** 

$$
\vec{\phi}^g = \mathbf{P}^g_{c, vv} \vec{q}^g
$$

with the complete collision probability matrix **P**  $_{c, vv}^g$  .

$$
\mathbf{P}_{c, vv}^g = \left( \mathbf{P}_{vv}^g + \mathbf{P}_{vs}^g ((\mathbf{A}^g)^{-1} - \mathbf{P}_{ss}^g)^{-1} \mathbf{P}_{sv}^g \right)
$$

Two types of multigroup cross-section database can be read by DRAGON

- Mixture macroscopic cross-section
- Isotope microscopic cross-section that contains itself <sup>a</sup> macroscopic cross-section database

Minimum cross-section requirements for each mixture  $m$ 

- The multigroup total cross section  $\Sigma$ g  $\,m$
- The isotropic component of the multigroup scattering cross section  $\Sigma$  $h{\rightarrow}g$  $_{m,s,0}^{\prime\prime\rightarrow g}$  defined as

$$
\Sigma_{m,s,0}^{h\to g} = \int_{4\pi} d^2\Omega^2 \Sigma_{m,s}^{h\to g} (\vec{\Omega}' \to \vec{\Omega}) P_0 (\vec{\Omega}' \cdot \vec{\Omega})
$$

- The product of the average neutron emitted per fission with the multigroup fission cross section  $\nu\Sigma$ g  $m,f$
- The multigroup fission spectrum  $\chi$ g  $\,m$

#### **Cross Sections Consideration 3** ⊬<br>®ÉCOLE

The linearly isotropic component of the multigroup scattering cross section  $\Sigma$  $h{\rightarrow}g$  $m,\!s,\!1$ 

$$
\Sigma_{m,s,1}^{h\to g} = \int_{4\pi} d^2\Omega^2 \Sigma_{m,s}^{h\to g} (\vec{\Omega}' \to \vec{\Omega}) P_1 (\vec{\Omega}' \cdot \vec{\Omega})
$$

Required only if  $B_1$  leakage method is used

The transport correction  $\Sigma$ g  $m,$ tc

**C** The transport calculations are performed using transport corrected total ( $\tilde{\Sigma}_m^g$ ) and scattering ( $\tilde{\Sigma}_{m.s.0}^{h\rightarrow g}$  $\binom{n \rightarrow y}{m,s,0}$ cross sections

$$
\begin{array}{l} \tilde{\Sigma}_m^g=\!\!\!\!\!\!\Sigma_m^g-\Sigma_{m,\text{tc}}^g\\ \tilde{\Sigma}_{m,s,0}^{h\to g}=\!\!\!\!\!\Sigma_{m,s,0}^{h\to g}-\delta^{gh}\Sigma_{m,\text{tc}}^g \end{array}
$$

Takes partially into account the linearly anisotropic scattering contributions

For contribution of multi-neutron production reactions such as

$$
\mathbf{X}_A^Z+\mathbf{n}_1^0\rightarrow\mathbf{X}_{A-1}^Z+2\mathbf{n}_1^0+\gamma
$$

The scattering cross section must be corrected to take into account this effect

$$
\tilde{\Sigma}_{m,s,0}^{h\to g} = \Sigma_{m,s,0}^{h\to g} + 2\delta^{gh}\Sigma_{m,(\text{n},2\text{n})}^g
$$

where  $\Sigma$ g  $\mathcal{C}_{(n,2n)}^{(n)}$  is the macroscopic cross section associated with the reaction

Macroscopic cross section data base can be created using

- from the input file using the <code>MAC</code>: module
- from a GOXS file using the MAC: module
- from <sup>a</sup> microscopic library using the LIB: module
- **•** from the homogenization and condensation module EDI:
- from <sup>a</sup> WIMS-AECL execution using the information available on TAPE16 (side-step method)
- **•** from a HELIOS execution

#### **Microscopic libraries 1** ÉCOLE<br>**TECHNIQUE**

Many formats can be processed by DRAGON including

- WIMS–AECL format
- **C** MATXS format
- WIMD-D4 format

In DRAGON resonance self-shielding calculations are preformed using the Stamm'ler method



## **Contents**

- Collision Probabilities in 3-D
- **C** Numerical Quadrature and Tracking
- **Collision Probability Integration**
- **C** Neutron Conservation and CP Normalization



Recall CP approximations

- Divide domain into  $N_V$  regions of volume  $V_i$
- Assume total cross sections constant inside eachregion
- Assume sources constant inside each region
	- This has an impact on the selection of the spatial  $\bullet$ mesh
- **Assume sources isotropic inside each region This may lead to problem when scattering is highly** anisotropic.



Recall BC approximations

- Divide the external boundary  $S$  into  $N_S$  surfaces of area  $S_{\alpha}$
- Assume flux constant on external surfaces
	- **This has an impact on the selection of the spatial** mesh
- Assume flux isotropic on external surfaces
	- This may lead to problems when flux is highly anisotropic near external boundaries



## **Collision Probabilities in 3-D 1**

Collision probability de finition

$$
\tilde{p}_{ij}^g = V_i p_{ij}^g = \int_{V_i} \int_{V_j} \frac{e^{-\tau^g(R)}}{4\pi R^2} \Theta_i \Theta_j d^3r' d^3r
$$

 $^3r$  integral Spherical coordinates for  $d$  $\bullet$ 

$$
\int_{V_j} d^3r \Theta_j = \int_{4\pi} d^2\Omega \int_{R_{i-\frac{1}{2}}}^{R_{i+\frac{1}{2}}} R^2 dR\Theta_j
$$

Cartesian coordinates for  $d$ 3  $r^\prime$  integral

$$
\int_{V_i} d^3r' \Theta_i = \int dx' \int dy' \int dR' \Theta_i
$$

#### **Collision Probabilities in 3-D 2 MANA**<br>POLYTECHNIQUE

General 3–D geometry for collision probability integration







## Final form for collision probability integration

$$
\tilde{p}_{ij}^g = \int_{4\pi} \frac{d^2\Omega}{4\pi} \int dx' \int dy' \int_{R_{i-\frac{1}{2}}}^{R_{i+\frac{1}{2}}} dR' \int_{R_{j-\frac{1}{2}}}^{R_{j+\frac{1}{2}}} dR \ e^{-\tau^g(R',R)} \Theta_i \Theta_j
$$
\n
$$
\tau^g = (R_{i+\frac{1}{2}} - R')\Sigma_i^g + \sum_{k=i+1}^{j-1} \Delta R_k \Sigma_k^g + (R - R_{j-\frac{1}{2}})\Sigma_j^g
$$

with  $\Delta R_k$  $=R_{k+\frac{1}{2}}$  $-R_{k-}$  $\frac{1}{2}$  .



## **Collision Probabilities in 3-D 4**

## Notation for optical path







After integration over  $R'$  and  $R$ , one obtains

$$
\tilde{p}_{ij}^g = \frac{1}{4\pi \Sigma_i^g \Sigma_j^g} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_j
$$
\n
$$
\times \left[1 - \exp\left(-\tau_{i-\frac{1}{2},i+\frac{1}{2}}^g\right)\right] \exp\left(-\tau_{i+\frac{1}{2},j-\frac{1}{2}}^g\right)
$$
\n
$$
\times \left[1 - \exp\left(-\tau_{j-\frac{1}{2},j+\frac{1}{2}}^g\right)\right]
$$

### with

$$
\tau_{i\pm\frac{1}{2},j\pm\frac{1}{2}}^g = \Sigma_i^g (R_{i+\frac{1}{2}} - R_{i\pm\frac{1}{2}}) + \tau_{i+\frac{1}{2},j-\frac{1}{2}} + \Sigma_j^g (R_{j\pm\frac{1}{2}} - R_{j-\frac{1}{2}})
$$


## **Collision Probabilities in 3-D 6**

Case where  $\Sigma$ g  $\frac{9}{i} = 0$ 

$$
\tilde{p}_{ij}^{g} = \frac{1}{4\pi \Sigma_{j}^{g}} \int_{0}^{4\pi} d^{2} \Omega \int dx' \int dy' \Theta_{i} \Theta_{j}
$$
  
 
$$
\times \Delta R_{i} \exp(-\tau_{i+\frac{1}{2},j-\frac{1}{2}}^{g}) \left[1 - \exp(-\tau_{j-\frac{1}{2},j+\frac{1}{2}}^{g})\right]
$$

Case where  $\Sigma$ g  $\frac{g}{i}=\Sigma_j^g=0$ 

$$
\tilde{p}_{ij}^g = \frac{1}{4\pi} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_j \Delta R_i \Delta R_j \exp(-\tau_{i+\frac{1}{2},j-\frac{1}{2}}^g)
$$

#### **Collision Probabilities in 3-D 7** ÉCOLE

For  $\tilde{p}^g_{ii}$ , one obtains after integration over  $R'$  and  $R$ 

$$
\tilde{p}_{ii}^{g} = \frac{1}{2\pi (\Sigma_{i}^{g})^{2}} \int_{0}^{4\pi} d^{2}\Omega \int dx' \int dy' \Theta_{i} \Theta_{i}
$$
\n
$$
\times \left[ \tau_{i-\frac{1}{2},i+\frac{1}{2}}^{g} - \left( 1 - \exp(-\tau_{i-\frac{1}{2},i+\frac{1}{2}}^{g}) \right) \right]
$$

For Σ  $g\,$  $\frac{g}{i}=0$  this is simplified to

$$
\tilde{p}_{ii}^g = \frac{1}{4\pi} \int_0^{4\pi} d^2\Omega \int dx' \int dy' \Theta_i \Theta_i \left(\Delta R_i\right)^2
$$

Similar relations are obtained for  $\tilde{p}^g_{i\ell}$  $_{i\alpha}^g$  and  $\tilde{p}_{\alpha}^g$  $\alpha\beta$ 



### 3-D Adjuster model in DRAGON



# **Quadrature and Tracking 2**

### Angular quadrature

- $S_n$  type  $\mathsf{EQ}_N$  quadrature with  $4N_{\Omega}(N_{\Omega}+2)/8$  angular directions  $\vec{\Omega}_{1,i}$ ,  $\vec{\Omega}_{2,i}$ ,  $\vec{\Omega}_{3,i}$  and  $\vec{\Omega}_{4,i}$
- Global quadrature weight  $W_{\Omega}=2/(N_{\Omega}(N_{\Omega}+2))$
- Tracking in lower half sphere only
- Number of tracking quadrant automatically reduced for symmetric cell
- Select as many angles as possible (Neutrons travel on <sup>a</sup> straight line)



#### Cartesian surface quadrature



# **Quadrature and Tracking 4**

### Cartesian quadrature

- Identify the radius  $h_{\pm}$  of the smallest sphere including the geometry
- Select a tracking density  $\rho_p$  and define the line spacing  $\delta$

$$
\delta = \frac{2h_+}{N_p} \qquad N_p = (2\sqrt{\rho_p}h_+) + 1
$$

Integration line  $\mathit{l}_{mn}$  passes through

$$
u_m^x = \left(\frac{2m-1}{2}\right)\delta \qquad u_n^y = \left(\frac{2n-1}{2}\right)\delta
$$

Integration weight  $W_p=\delta$ 2



### Comments on Cartesian surface quadrature

In Dragon 3 different planes are selected for each given spatial direction



- Select tracking density as dense as possible
	- **Each region must be touched by a maximum number** of lines



### DRAGON tracking example





DRAGON tracking of a line  $\mathit{l}_{m,n}$ 

- 1. Follow the tracking line as it travels through the cell
- 2. Identify each region  $i$  and surface  $s$  uniquely
- 3. Identify final region number  $I$  associated with set of regions  $i$  (a flux region)
- 4. Identify final surface number  $S$  associated with sets of surfaces  $s$
- 5. Identify external surfaces and regions  $i$  crossed by line
- 6. Compute distance  $l_{i,m,n}$  the neutron travels in each ˜region
- 7. Store information on temporary tracking file

#### **Quadrature and Tracking 8** ÉCOLE

Case of symmetric cells (mirror re flection on one side of the cell)

**C** Unfold the cell according to symmetry



Regions in unfolded cells are originally assigned new region numbers

# **Quadrature and Tracking 9**



Angle selection for symmetric cells

- For symmetry with respect to a  $(y,z)$  plane, track only in quadrants corresponding to directions  $\vec{\Omega}_{1,i}$  and  $\vec{\Omega}_{3,i}$
- For symmetry with respect to a  $(z,x)$  plane, track only in quadrants corresponding to directions  $\vec{\Omega}_{1,i}$  and  $\vec{\Omega}_{2,i}$
- For symmetry with respect to a  $(x,y)$  plane, track only in all quadrants



Post treatment of tracking file

- Assign to all spatial region  $i$  its final flux region number  $\bullet$ combining track segments as required
- **For each direction, normalize tracks using**

$$
l_{i,m,n} = \left(\frac{V_i}{\tilde{V}_i}\right) \tilde{l}_{i,m,n}
$$

where

$$
\tilde{V}_i = \frac{W_p}{3} \sum_{m=1}^{N_p} \sum_{n=1}^{N_p} \tilde{l}_{i,m,n}
$$

#### **Quadrature** and Tracking ÉCOLE<br>**TECHNIQUE**

Comments on storage requirements for tracking Maximum number of line segments tracks  $d_t$  $\bullet$ 

$$
d_t < 6N(N_\Omega(N_\Omega+2))\rho_p h_+^2
$$

• For 
$$
N = 1000
$$
,  $h_+ = 50$ ,  $\rho_p = 20$  t/cm<sup>2</sup> and  $N_{\Omega} = 8$   
 $d_t < 20 \times 10^9$ 

tracks segments.



#### ACR control rod model in DRAGON







#### 3-D cluster analysis not currently allowed in DRAGON





For each energy group

- Read <sup>a</sup> line from tracking file  $\bullet$
- Scan this line and add contribution to  $\tilde{p}^g_i$  $ii \$

$$
\frac{1}{2}(\Sigma_i^g)^2 \tilde{p}_{ii}^g = \frac{1}{2}(\Sigma_i^g)^2 \tilde{p}_{ii}^g
$$

$$
+ \sum_n W_n \sum_{m \in i} \left(\tau_{i,n,m}^g - \kappa_{i,n,m}^g\right)
$$

where  $W_n$  $= W_\Omega W_p/3$  and

$$
\kappa_{i,n,m}^g = \left(1 - \exp\left[-\tau_{i,n,m}^g\right]\right)
$$
  

$$
\tau_{i,n,m}^g = \Sigma_i^g l_{i,n,m} \qquad \text{with } m \in i
$$



#### Scan this line a second time and add contributions to  $\tilde{p}^g_i$  $ij$

$$
\Sigma_i^g \Sigma_j^g \tilde{p}_{ij}^g = \Sigma_i^g \Sigma_j^g \tilde{p}_{ij}^g
$$
  
+ 
$$
\sum_n W_n \sum_{m \in i} \sum_{m' \in j} \kappa_{i,n,m}^g \kappa_{n,m+1,m'-1}^g \kappa_{j,n,m'}^g
$$

using

$$
\kappa_{n,m,m'}^g = \prod_{l=m}^{m'} \exp\left[-\tau_{i,n,l}^g\right]
$$



### Finish CP calculations

- Only the contributions with  $m < m'$  has been considered
- Symmetrize  $p$  $_{ij}^g$  using

$$
\tilde{p}_{ij}^g=\tilde{p}_{ij}^g+\tilde{p}_{ji}^g
$$



### Compute errors on CP conservation rules

$$
R_j^g = \Sigma_j^g V_j - \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} \Sigma_i^g p_{\alpha j}^g - \sum_{i=1}^{N_i} \Sigma_j^g \Sigma_i^g V_i p_{ij}^g
$$

$$
R_\beta^g = \frac{S_\beta}{4} - \sum_{\alpha=1}^{N_\alpha} \frac{S_\alpha}{4} p_{\alpha \beta}^g - \sum_{i=1}^{N_i} \Sigma_i^g V_i p_{i \beta}^g
$$







Diagonal Normalization

$$
p_{D,ii}^g = p_{ii}^g - \frac{R_i^g}{(\Sigma_i^g)^2 V_i}
$$

$$
p_{D,\alpha\alpha}^g = p_{\alpha\alpha}^g - \frac{4R_\alpha^g}{S_\alpha}
$$

- May result in non-physical negative probabilities
- Cannot be applied to problems involving voided zones



HELIOS Type Normalization

$$
p^g_{H,ij} = (w^g_i + w^g_j)p^g_{ij} \qquad p^g_{H,\alpha\alpha} = (w^g_\alpha + w^g_\beta)p^g_{\alpha\alpha}
$$

- Apply conservation laws to above relation
- Solve resulting system for  $w^g$  using an iterative process
- Does not lead to negative probabilities and works for void regions
- Default option in DRAGON



Comments on storage requirements for CP matrices

- Number of elements per groups is  $N$ 2
- Memory space required for execution is about  $5N$ 2
- Total disk space required for storage of  $G$  group CP is  $GN^2$



Contents:

- Geometry.
- Collision Probability Integration and Tracking.  $\bullet$
- Region Merging.  $\bullet$



# **3–D DRAGON Examples 2**

## 3–D Geometry restrictions in DRAGON:

- Cartesian mesh in each direction must extend to thewhole geometry.
- **Cluster option not permitted.**
- A single cylinder per cell.
- Cylinders cannot intersect other than axially.
- Cylinders must extend to the whole geometry.
- Cylinders are by default centered in the cell. They can be displaced using the OFFCENTER option.
- Mixtures are specified radially, then in  $x,\,y$  and  $z.$
- Mixtures are speci fied even in location that do not exists.





#### A Simple 3–D Cell:









#### Mixture speci fication for simple 3–D cell:

```
TMPGEO := GEO: :: CAR3D 1 1 1CELL FC1B
X- REFL X+ REFL Y- REFL Y+ REFL Z- REFL Z+ REFL::: FC1B := GEO: CARCELZ 2321
    MESHX <<<MXLP>> <<MXL>> <<MXL>>> <<<<<<br/>MXLD>>> <<<<br/>MXYD>>
    MESHY <<MYLP>> <<CALLCN>> <<PYLP>>MESHZ 0.0 49.5OFFCENTER <<FC1XD>> <<FC1YD>>RADIUS 0.0 <<RF2>> <<RCT>>
    MIX 1 2 311 12 1321 22 2331 32 3341 42 4351 52 53 ;
```
;

Note: mixtures 1, 2, 31 and 32 not used.



CANDU adjuster rod simulation:

- 5 region annular fuel (including PT and CT).
- 6 region annular adjuster rod.





### Coarse mesh geometry de finition for CANDU adjuster:



#### **3–D DRAGON Examples 7** ÉCOLE **OLYTECHNIQUE**

#### DRAGON geometry for CANDU adjuster rod:

```
TMPGEO := GEO: :: CAR3D 2 1 1CELL FC1B AD1T
X- REFL X+ SYME Y- REFL Y+ SYME Z- REFL Z+ SYME::: FC1B := GEO: CARCELZ 5113
    MESHX 0.0000 21.5750 MESHY 0.0000 28.5750MESHZ 0.0000 17.7650 31.7650 49.5300OFFCENTER 3.5 0.0RADIUS 0.0 0.7222 2.1603 3.6007 5.1689 6.5875MIX 1 2 3 4 5 6 1 2 3 4 5 6 1 2 3 4 5 6 1 2 3 4 5 6 ;
 ::: AD1T := GEO: CARCELY 6113
    MESHX 21.5750 35.5750 MESHY 0.0000 28.5750MESHZ 0.0000 17.7650 31.7650 49.5300RADIUS 0.0 0.5770 3.6781 3.8100 4.4450 4.7520 6.3776MIX 7 8 9 10 11 12 1 7 8 9 10 11 12 17 8 9 10 11 12 1 7 8 9 10 11 12 1 ;
```
;



#### Coarse mesh CANDU adjuster rod after unfolding:





#### Fine mesh CANDU adjuster rod after unfolding:



#### **3–D DRAGON Examples 10** ÉCOLE



### DRAGON geometry for CANDU adjuster rod:

```
TMPGEO := GEO: :: CAR3D 2 1 1CELL FC1B AD1T
X- REFL X+ SYME Y- REFL Y+ SYME Z- REFL Z+ SYME::: FC1B := GEO: CARCELZ 5113
    MESHX 0.0000 21.5750 SPLITX 3 MESHY 0.0000 28.5750 SPLITY 3MESHZ 0.0000 17.7650 31.7650 49.5300OFFCENTER 3.5 0.0RADIUS 0.0 0.7222 2.1603 3.6007 5.1689 6.5875MIX 1 2 3 4 5 6 1 2 3 4 5 6 1 2 3 4 5 6 1 2 3 4 5 6 ;
 ::: AD1T := GEO: CARCELY 6113
    MESHX 21.5750 35.5750 SPLITX 2 MESHY 0.0000 28.5750 SPLITY 3MESHZ 0.0000 17.7650 31.7650 49.5300 SPLITZ 2RADIUS 0.0 0.5770 3.6781 3.8100 4.4450 4.7520 6.3776MIX 7 8 9 10 11 12 1 7 8 9 10 11 12 17 8 9 10 11 12 1 7 8 9 10 11 12 1 ;
```
;



Exact boundary conditions:

- VOID: applied at the explicit boundary of the cell or assembly.
- SYME: applied at the center of the cells closest to the explicit assembly boundary speci fied.
- DIAG: applied at the center of the cells closest to the explicit assembly boundary speci fied.
- SSYM: applied at the explicit boundary of the cell or assembly.





Approximate boundary conditions:

- REFL: applied at the explicit boundary of the cell or assembly. Exact specular option not available in 3–D.
- TRAN: applied at the explicit boundary of the cell or assembly. Exact specular option not available in 3–D.
- ALBE: applied at the explicit boundary of the cell or assembly. Exact specular option not available.



Region identi fication for single cell:

- radially outward in <sup>a</sup> cell.
- from lower to upper  $x$  location in a cell.
- from lower to upper  $y$  location in a cell.
- from lower to upper  $z$  location in a cell.

Region identi fication for assembly of cells:

- Inside each cell as above.
- from lower to upper  $x$  cell location in the assembly.
- from lower to upper  $y$  cell location in the assembly.
- from lower to upper  $z$  cell location in the assembly.



#### Region identi fication for cells and assemblies:







*Front*





*one cell three cells assembly*


# Region identi fication for one cell problem:



--

--



# Region identi fication for three cells assembly:

--





Quadrature selection:

- As many angles as possible:  $\rightarrow$  neutron travels on a straight line.
- **•** Tracking density must be as dense as possible:  $\rightarrow$  to touch as often as possible each region and surface.
- For CANDU reactivity devices TRAK TISO <sup>8</sup> <sup>25</sup>:  $\rightarrow$  10 angles per quadrant.

 $\rightarrow 3 \times 25$  tracks per cm $^2$ .

**Integration lines are renormalized using ratio of** approximate to exact volumes.





Comments on storage requirements:

Size of tracking file linear in  $N$ :

$$
d_t \propto \rho h_+^2 N
$$

for  $h_+=50$  cm,  $\rho=20$  t/cm $^2$  and  $N=1000$  regions:  $\rightarrow d_t = 600$  Mb.

Size of CP matrix quadratic in  $N$ :

$$
d_a \propto N^2 G
$$

for  $G=89$  groups and  $N=1000$  regions:  $\rightarrow d_a = 356$  Mb.

Use XSM\_FILE for ASMPIJ data structure.



Example of storage requirements for <sup>a</sup> simple 3–D problem.

- Total volume 1 liter ( $V=10^3~{\sf cm}^3$ ).
- Central fissile region is red ( $V=27~{\rm cm}^3$ ).
- Strong absorber is green ( $V=1~{\rm cm}^3$ ).



#### **Region Merging 2** <sup>≠</sup><br>® ÉCOLE<br>⁄**TECHNIQUE**

- Region with strong absorber:
	- Try to avoid using approximate boundary conditions.
	- **Fine mesh discretization is required.**
- Region with fission:
	- **Try to avoid using approximate boundary conditions.**
	- Medium to fine mesh discretization is required.
- **C** For moderator region
	- **Fine to coarse mesh discretization is required.**



#### Uniform mesh for simple 3–D problem:



Note:  $x-y$  and  $x-z$  planes are identical.



## Non-uniform meshes in DRAGON:



- $N=19\times 16\times 16=4864$  regions,  $d_a=100$  Mb/groups for uniform mesh.
- $N\approx 1500$  regions,  $d_a=9$  Mb/groups for non-uniform  $\bullet$ mesh.



### Using the MRG: module:



...





# Region numbering for first (out of 16)  $z$ -plane:







#### **Solving the CP Equations 1**  $\left(\begin{matrix} \bigcirc \\ \bigcirc \end{matrix}\right)$   $\epsilon$  cole

# Contents

- **C** The Power Iteration
- **C** The Multigroup Iteration
- Leakage Models  $\bullet$

#### **Solving the CP Equations 2** *》*<br>☞ ÉCOLE<br><mark>YTECHNIOUE</mark>

# The multigroup transport equation has the form

$$
\vec{\phi} = \mathbf{P}_{c, vv} (\vec{q}_s + \frac{1}{k} \vec{q}_f)
$$

$$
\vec{q}_s = \Sigma_s \vec{\phi}
$$

$$
\vec{q}_f = \chi \nu \Sigma_f \vec{\phi}
$$

- $\vec{\phi}$  is a  $N \times G$  dimensional vector
- $\bullet$   $\mathbf{P}_{c,vv}$  is the multigroup CP matrix
	- **Diagonal in energy, full in space.**
- $\chi\nu\boldsymbol{\Sigma}_f$  is a matrix for neutron production by fission and  $\Sigma_s$  is the scattering matrix
	- **Diagonal in space, full in energy**





$$
\boldsymbol{\Sigma}_{s}=\!\boldsymbol{\Sigma}_{d,s}+\boldsymbol{\Sigma}_{u,s}+\boldsymbol{\Sigma}_{w,s}
$$

#### with

- $\sum_{u,s}$  the up-scattering matrix (lower triangular in energy)
- $\sum_{d,s}$  the down-scattering matrix (upper triangular in energy)
- Σ<sub>w,s</sub> the within-group scattering matrix (diagonal in energy)



# De fining **W** the scattering modi fied CP matrix

$$
\mathbf{W} = (\mathbf{I}-\mathbf{P}_{c,vv}\mathbf{\Sigma}_{w,s})^{-1}\mathbf{P}_{c,vv}
$$

the transport equation now becomes

$$
\vec{\phi} = \mathbf{W}(\Sigma_{d,s}\vec{\phi} + \Sigma_{u,s}\vec{\phi} + \frac{1}{k}\vec{q}_f)
$$

Assuming

• 
$$
\vec{q}_f/k
$$
 is fixed and  $\Sigma_{w,s} = 0$ 

The above equation can be solved directly from group  $g=1$ to  $g=G$ 

#### **Solving the CP Equations 5**

The general equation can be solve using two iteration processes

- **C** The Power iteration illustrated by solving the transport problem with  $\mathbf{\Sigma}_{w,s}=0$
- **C** The multigroup iteration illustrated by solving the transport problem with  $\vec{q}_f/k$ fixed

#### **The Power Iteration 1** ,<br>⊕ÉCOLE<br>**'TECHNIQUE**

- $\mathsf{Assume}\ \mathbf{\Sigma}_{u,s}=0$
- Write an iterative group by broup solution to the transport problem as

$$
\vec{\phi}(l) = \mathbf{W}(\Sigma_{d,s}\vec{\phi}(l) + \frac{\chi}{k(l-1)}\nu\Sigma_f\vec{\phi}(l-1))
$$

$$
k(l) = \sum_{g=1}^G \sum_{i=1}^N V_i \chi_i^g \sum_{h=1}^G \nu \Sigma_{f,i}^h \phi_i^h(l)
$$

with  $\vec{\phi}(0)$  a known arbitrary flux distribution and

$$
k(0) = \sum_{g=1}^{G} \sum_{i=1}^{N} V_i \chi_i^g \sum_{h=1}^{G} \nu \Sigma_{f,i}^h \phi_i^h(0)
$$



# **The Power Iteration 1**

 $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$ 

# The iteration process is repeated until

$$
k(l) - k(l-1) < \epsilon_1
$$
\n
$$
\left| \frac{\vec{\phi}(l)}{k(l)} - \frac{\vec{\phi}(l-1)}{k(l-1)} \right| < \epsilon_2 \left| \frac{\vec{\phi}(l)}{k(l)} \right|
$$

### are both satisfied

The parameters  $\epsilon_1$  and  $\epsilon_2$  can be defined independently in DRAGON

#### **The Multigroup Iteration 1 LOZ**<br>POLYTECHNIQUE<br>POLYTECHNIQUE

 $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$ 

• Assume 
$$
\vec{q}_f/k = \vec{q}
$$
 is fixed

Solve group-by-group this fixed source problem using <sup>a</sup>  $\bullet$ Gauss-Seidel strategy

$$
\vec{\phi}(l) = \mathbf{W}\left(\mathbf{\Sigma}_{d,s}\vec{\phi}(l) + \mathbf{\Sigma}_{u,s}\vec{\phi}(l-1) + \vec{q}\right)
$$

Iterate until

$$
\left| \frac{\vec{\phi}^g(l) - \vec{\phi}^g(l-1)}{\vec{\phi}^g(l)} \right| < \epsilon_3
$$

#### **The Multigroup Iteration 2 O**<br>Se ÉCOLE<br>POLYTECHNIQUE

Multigroup rebalancing technique

Neutron conservation states that for <sup>a</sup> converged solution  $\phi_i^g$  $_{i}^{g}(l)$ 

$$
\sum_{i=1}^{N} \sum_{i=1}^{g} V_i \phi_i^g(l) = \sum_{i=1}^{N} R_i^g V_i q_i^g(l)
$$
  
+ 
$$
\sum_{i=1}^{N} R_i^g V_i \left( \sum_{h=1}^{g} \sum_{s}^{h \to g} \phi_i^h(l) + \sum_{h=g+1}^{G} \sum_{s}^{h \to g} \phi_i^h(l) \right)
$$

with  $R$ g  $\frac{9}{i} = 1$  $-\sum_{i=1}^N$ Σ g  $\H{j}{p}$ g  $c, ij$ 

#### **The Multigroup Iteration 3**  $\circled{(*)}$   $\circled{})$   $\circled$

**•** The multigroup problem we effectively solve yields

$$
\sum_{i=1}^{N} \sum_{i=1}^{g} V_i \phi_i^g(l) = \sum_{i=1}^{N} R_i^g V_i q_i^g(l)
$$
  
+ 
$$
\sum_{i=1}^{N} R_i^g V_i \left( \sum_{h=1}^{g} \sum_{s=1}^{h \to g} \phi_i^h(l) + \sum_{h=g+1}^{G} \sum_{s=1}^{h \to g} \phi_i^h(l-1) \right)
$$

#### **The Multigroup Iteration 4** (<mark>⑥</mark>)<br>POLYTECHNIQUE<br>POLYTECHNIQUE

To restore conservation at each iteration

Use  $\tilde{\phi}_i^g$  $\frac{g}{i}=\alpha^g\phi^g_i$  $\begin{array}{c} g \end{array}$  and assume  $\tilde{\phi}_i^g$  $\frac{g}{i}$  satisfies conservation relations, then  $\alpha^g$  must satisfy

$$
\sum_{h=1}^{G} M^{h \to g} \alpha^h = q^g
$$
\n
$$
M^{h \to g} = \sum_{i=1}^{N} R_i^g V_i \left( \sum_{i=1}^{h} \delta_{gh} - \sum_{s=1}^{h \to g} \right) \phi_i^h
$$
\n
$$
q^g = \sum_{i=1}^{N} R_i^g V_i q_i^g
$$

Solve for  $\alpha^g(k)$  and rebalance flux

#### **The Multigroup Iteration 5** ÉCOLE



For appromimate solution  $\vec{\Gamma}$  $\Gamma(l)$ 

$$
\vec{\Gamma}(l) = \mathbf{W}\left(\boldsymbol{\Sigma}_{d,s}\vec{\Gamma}(l)+\boldsymbol{\Sigma}_{u,s}\vec{\Gamma}(l-1)+\vec{q}\right)
$$

De fine an improved flux distribution for the next iteration using

$$
\vec{\phi}(l) = \vec{\Gamma}(l) + \omega(l)\vec{\Delta}(l)
$$

$$
\vec{\Delta}(l) = \vec{\Gamma}(l) - \vec{\phi}(l-1)
$$

and  $\omega(l)$  will be computed using a variational procedure

#### **The Multigroup Iteration 6**  $\bigotimes_{\epsilon \in \text{COLE}}$ <br>POLYTECHNIQUE

Select  $\omega(l)$  in such a way that  $\vec{\phi}(l)$  minimizes the tranport functional

$$
\mathcal{F}[\vec{\phi}] = \frac{1}{2} \vec{\phi}^T \mathbf{Z}^T \mathbf{Z} \vec{\phi} - \vec{\phi}^T \mathbf{Z}^T \mathbf{W} \vec{q}
$$

$$
\mathbf{Z} = [\mathbf{I} - \mathbf{W} \left( \mathbf{\Sigma}_{d,s} + \mathbf{\Sigma}_{u,s} \right)]
$$

### This yield

$$
\omega(l) = -\frac{[\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]^T [\vec{\Gamma} - \mathbf{W}\vec{S}_2(l)]}{[\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]^T [\vec{\Delta} - \mathbf{W}\vec{S}_1(l)]}
$$

$$
\vec{S}_1(l) = (\mathbf{\Sigma}_{d,s} + \mathbf{\Sigma}_{u,s})\vec{\Delta}(l)
$$

$$
\vec{S}_2(l) = (\vec{q} + \mathbf{\Sigma}_{d,s}\vec{\Gamma}(l) + \mathbf{\Sigma}_{u,s}\vec{\Gamma}(l))
$$

#### **Leakage Models 1**  $\bigotimes_{k=0}^{\infty}$ école<br>Polytechnique<br>MONTRÉAI

In 3–D, the transport equation can be solved in DRAGON using the  $B_0$  and  $B_1$  leakage models

■ Both of these models are based on the following factorization of the flux

$$
\Phi^g(\vec{r}, \vec{\Omega}) \approx \Psi^g(\vec{r}, \vec{\Omega}) \exp(i\vec{B} \cdot \vec{r})
$$

Transport equation with leakage

$$
\vec{\Omega} \cdot \vec{\nabla} \Psi^g(\vec{r}, \vec{\Omega}) + [\Sigma^g(\vec{r}) + i \vec{B} \cdot \vec{\Omega}] \Psi^g(\vec{r}, \vec{\Omega}) =
$$
  

$$
Q_s^g(\vec{r}, \vec{\Omega}) + \frac{1}{k_{\text{eff}}} Q_f^g(\vec{r})
$$

In general we will assume that  $k_{\mathit{eff}}$ =1



For an in finite homogeneous media the scalar flux and vector current are related to each other according to

$$
\vec{\Omega}\Psi^g(\vec{\Omega}) = \vec{J}^g(\vec{\Omega}) = -iD^g\vec{B}\Psi^g(\vec{\Omega})
$$

with  $D^g$  is an homogeneous diffusion coefficient

**Apply to heterogeneous systems** 

 $\vec{\Omega} \cdot \vec{\nabla}\Psi$  $^{g}(\vec{r}% )^{g}(\vec{r})$  $(\vec r, \vec \Omega) + [\Sigma^g(\vec r) + D]$  $^{g}B^{2}]\Psi^{g}(\vec{r})$  $(\vec{r},\vec{\Omega})=Q_s^g$  $\frac{g}{s}(\vec{r})$  $(\vec{r},\vec{\Omega})+Q^g_{~f}$  $\frac{g}{f}(\vec{r})$ 

Find the homogeneous diffusion coef ficient compatible with this heterogeneous problem

#### **Leakage Models 3** <sup>≠</sup><br>® ÉCOLE<br>⁄**TECHNIQUE**

- Assume an heterogeneous solution is known for  $B$  $^2=0$
- Use this solution to de fine an equivalent in finite homogeneous problem

$$
\Sigma^g \Psi^g(\vec{\Omega}) + i \vec{B} \cdot \vec{J}^g(\vec{\Omega}) = Q_s^g(\vec{\Omega}) + Q_f^g
$$

where the cross sections and sources are homogenized using the heterogeneous flux

- Solve the homogeneous problem for  $D^g$  and  $B$
- Insert in heterogeneous transport equation and obtain an improved solution.
- Repeat until the iterative procedure is converged



Solving the homogeneous problem  $(B_{1}% ,B_{2})\equiv\mathcal{A}(B_{1})$  model)

Use <sup>a</sup> 2 terms expansion for the scattering cross section in Legendre polynomials

$$
\Sigma_s^{h\rightarrow g}(\vec{\Omega}' \rightarrow \vec{\Omega}) = \Sigma_{s,0}^{h\rightarrow g} + 3\Sigma_{s,1}^{h\rightarrow g}\vec{\Omega}\cdot\vec{\Omega}'
$$

Define

$$
\psi^g = \int d^2\Omega \Psi^g(\vec{\Omega})
$$

$$
\vec{j}^g = \int d^2\Omega \vec{\Omega}' \Psi^g(\vec{\Omega})
$$



- 
- Insert into the homogeneous transport equation, and  $\bullet$ integrate to obtain

$$
\psi^g = \alpha^g \sum_h (\Sigma_{s,0}^{h \to g} + \chi^g \nu \Sigma_f^h) \psi^h + 3\beta^g \sum_h \Sigma_{s,1}^{h \to g} \frac{\vec{B} \cdot \vec{j}^h}{iB^2}
$$
  

$$
\vec{j}^g = \beta^g \sum_h \left[ (\Sigma_{s,0}^{h \to g} + \chi^g \nu \Sigma_f^h) \frac{\vec{B} \psi^h}{iB^2} + 3\Sigma^g \Sigma_{s,1}^{h \to g} \frac{\vec{j}^h}{B^2} \right]
$$
  

$$
\alpha^g = \frac{1}{B} \arctan\left(\frac{B}{\Sigma^g}\right)
$$
  

$$
\beta^g = 1 - \Sigma^g \alpha^g
$$

Solve for  $B$ ,  $\psi^g$  and  $\vec{j}^g$  and compute  $D^g=i\vec{B}\cdot\vec{j}^g/B^2$  $\psi^g$ 





# The  $B_{\rm 0}$  homogeneous problem

Assume  $\Sigma$  $h{\rightarrow}g$  $\mathcal{S}_{s,1}^{n \rightarrow g} = 0$  and obtain

$$
\psi^g = \alpha^g \sum_h (\Sigma_{s,0}^{h \to g} + \chi^g \nu \Sigma_f^h) \psi^h
$$

$$
\vec{j}^g = \beta^g \frac{\vec{B}}{iB^2} \sum_h (\Sigma_{s,0}^{h \to g} + \chi^g \nu \Sigma_f^h) \psi^h
$$

The homogeneous diffusion coef ficient is then given by

$$
D^g=\!\frac{\beta^g}{\alpha^g}
$$

# **Condensation and Homogenization 1**

# **Contents**

- **Condensation Technique**
- Full Cell Homogenization  $\bullet$
- Partial Cell Homogenization and SPH Factors
- Microscopic Cross Section Homogenization

# **Condensation and Homogenization 2**

Condensation and homogenization techniques in DRAGON are based on the following assumptions

Reaction rates are physically meaningfull and should be  $\bullet$ preserved by the condensation/homogenization procedure

$$
R_i = \sum_g V_i \phi_i^g \Sigma_i^g = V_i \phi_i \Sigma_i
$$

$$
R^g = \sum_i V_i \phi_i^g \Sigma_i^g = V \phi^g \Sigma^g
$$

The eigenvalue is physically meaningfull and should be preserved by the condensation/homogenization procedure

# **Condensation Technique**

Condensed transport equation (macrogroup  $K$  that includes  $g\in G_K$ )

$$
\vec{\Omega} \cdot \vec{\nabla} \sum_{g \in G_K} \Phi^g(\vec{r}, \vec{\Omega}) + \sum_{g \in G_K} \Sigma^g(\vec{r}) \Phi^g(\vec{r}, \vec{\Omega}) =
$$
  

$$
\sum_{g \in G_K} [Q_s^g(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q_f^g(\vec{r})]
$$

Few group version of the same equation is

$$
\vec{\Omega} \cdot \vec{\nabla} \Phi^K(\vec{r}, \vec{\Omega}) + \Sigma^K(\vec{r}) \Phi^K(\vec{r}, \vec{\Omega}) = [Q^K_s(\vec{r}, \vec{\Omega}) + \frac{1}{k} Q^K_f(\vec{r}, \vec{\Omega})]
$$

It should reproduce condensed multigroup results

#### **Condensation Technique 2 WAY**<br>POLYTECHNIQUE<br>POLYTECHNIQUE

## The condensation procedure that satis fies our requirements

$$
\phi_i^K = \sum_{g \in G_K} \phi_i^g
$$
  
\n
$$
\Sigma_i^K = \frac{1}{\phi_i^K} \sum_{g \in G_K} \Sigma_i^g \phi_i^g
$$
  
\n
$$
\Sigma_{s,i}^{L \to K} = \frac{1}{\phi_i^L} \sum_{h \in G_L} \sum_{g \in G_K} \Sigma_{s,i}^{h \to g} \phi_i^h
$$
  
\n
$$
\chi_i^K = \sum_{g \in G_K} \chi_i^g
$$
  
\n
$$
\nu \Sigma_{f,i}^K = \frac{1}{\phi_i^K} \sum_{g \in G_K} \nu \Sigma_{f,i}^h \phi_i^g
$$

# **Condensation Technique 3**



Multiplying CP transport equation by  $\Sigma$ g  $\frac{g}{i}V_i$  and summing over all regions  $i$  yields

$$
\sum_{i} \Sigma_i^g V_i \phi_i^g = \sum_{i} V_i [Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]
$$

The equivalent transport equation in <sup>a</sup> homogeneous infinite cell is

$$
\hat{\Sigma}^g V \hat{\phi}^g = V [\hat{Q}_{s,i}^g + \frac{1}{k} \hat{Q}_{f,i}^g]
$$

The homogenized and homogeneous transport equations are identical if one selects a flux-volume homogenization technique

#### **Condensation Technique O**<br>Se école<br>POLYTECHNIQUE

# Flux-volume homogenization technique





Flux-volume homogenization fails if

The cell is finite (a cell with leakage) and

$$
\sum_{j=1}^{N_j} p_{ij}^g \Sigma_j^g \neq 1
$$

Partial cell homogenization cell is considered  $\bullet$
## **Partial Cell Homogenization 1**

The heterogeneous  $N$  region transport equation homogenized over  $M$  regions takes the form

$$
\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = \sum_{i \in M_I} \sum_{J} \sum_{j \in M_J} p_{ji}^g (\Sigma^g) [Q_{s,i}^g + \frac{1}{k} Q_{f,i}^g]
$$

The  $M$  region heterogeneous transport equation takes the from

$$
V_I \hat{\Sigma}_I^g \hat{\phi}_I^g = \sum_J \hat{p}_{JI}^g (\hat{\Sigma}^g) [Q_{s,J}^g + \frac{1}{k} Q_{f,J}^g]
$$

where  $\hat{P}^{g}_{\tau}$  $\hat{L}_{JI}^{g}(\hat{\Sigma}^{g})$  indicates that the CP are computed using homogenized cross sections

#### **Partial Cell Homogenization 2**  $\left(\begin{matrix} \bigcirc \\ \bigcirc \\ \bigcirc \end{matrix}\right)$   $\epsilon$  cole

#### We need

$$
\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = V_I \Sigma_I^g \phi_I^g
$$

and

$$
\sum_{J} \hat{p}_{JI}^{g} (\hat{\Sigma}^{g}) [Q_{s,J}^{g} + \frac{1}{k} Q_{f,J}^{g}] =
$$
  

$$
\sum_{i \in M_{I}} \sum_{J} \sum_{j \in M_{J}} p_{ji}^{g} (\Sigma^{g}) [Q_{s,i}^{g} + \frac{1}{k} Q_{f,i}^{g}]
$$

to be simultaneously true

#### **Partial Cell Homogenization 3** 。<br>ÉCOLE

The flux-volume homogenization method is not longer adequate because

- There is no simple relation between  $\hat{p}_{J}^{g}$  $g_{JI}(\hat{\Sigma}^g)$  and  $p$  $\frac{g}{ji}(\Sigma^g)$ The alternative here is to use <sup>a</sup> non-linear process
	- Consider a flux-volume homogenization for  $\phi^g_I$  $^g_I$  and  $\Sigma$ g I
	- Redefine the homogeneous flux  $\hat{\phi}^g_I$  $_I^g$  and cross sections  $\hat{\Sigma}^g$  $_J^g$  as follows

$$
\hat{\phi}_I^g = \frac{1}{\mu_I^g} \phi_I^g \qquad \hat{\Sigma}_I^g = \mu_I^g \Sigma_I^g
$$

#### **Partial Cell Homogenization 4** ÉCOLE

Determine the SPH factors  $\mu$ g  $_J^g$  numerically in such a way that

$$
\sum_{J} \hat{p}_{JI}^{g} (\hat{\Sigma}^{g}) [Q_{s,J}^{g} + \frac{1}{k} Q_{f,J}^{g}] =
$$
  

$$
\sum_{i \in M_{I}} \sum_{J} \sum_{j \in M_{J}} p_{ji}^{g} (\Sigma^{g}) [Q_{s,i}^{g} + \frac{1}{k} Q_{f,i}^{g}]
$$

is true

The de finition of the SPH factors automatically ensures

$$
\sum_{i \in M_I} V_i \Sigma_i^g \phi_i^g = V_I \Sigma_I^g \phi_I^g = V_I \tilde{\Sigma}_I^g \tilde{\phi}_I^g
$$

## **Microscopic Cross Section 1**

The macroscopic cross section associated with <sup>a</sup> material is simply the sum over all isotopes of the isotopic macroscopic cross section  $\Sigma_I$  namely

$$
\Sigma_i^g=\sum_I\Sigma_{I,i}^g
$$

#### where

$$
\Sigma_{I,i}^g = N_{I,i} \sigma_I^g
$$

with  $N_{I,i}$ , the concentration of isotope  $I$  in region  $i$ 

**C** The homogenization and condensation procedure described above remain valid for  $\Sigma$ g  $I, i$ 

#### **Microscopic Cross Section 2** ÉCOLE

Since the final concentration of isotope  $I$  in the cell is given by:

$$
N_I = \frac{1}{V} \sum_i N_{I,i} V_i
$$

we can de fine the equivalent homogenized microscopic cross section as:

$$
\hat{\sigma}_I^K = \frac{\mu_I^K}{N_I V \phi^K} \sum_{i \in M_I} \sum_{g \in G_K} N_{I,i} V_i \sigma_I^g \phi_i^g
$$

where the microscopic cross sections now become dependent on the spatial position

#### **Managing <sup>a</sup> DRAGON Execution 1** ÉCOLE

### **Contents**

- Input file formats.
- Data structure formats.
- Working with variables.
- Conditional execution and loops.
- Working with procedures.
- **C** Flow chart in DRAGON input decks

# **Managing <sup>a</sup> DRAGON Execution 2**

Input file format

- 72 columns, free format instruction ends by ;
- Comments \* or !
- MODULE and objects declarations
- **Sequence of calls to modules**

```
(list of output objects) := GEO: (list of input objects) ::
```

```
(data input)
```
- END: ; statement
- QUIT\_"LIST" \_. end compilation

#### **Managing <sup>a</sup> DRAGON Execution 3** ÉCOLE

Data structure formats.

- **C** LINKED LIST Memory access
- $\bullet$  XSM FILE Direct-access file
- SEQ BINARY Tracking information mainly
- SEQ\_ASCII Machine independent format
- DIR\_ACCESS XS library file

# **Managing <sup>a</sup> DRAGON Execution 4**

### Variable types.

- INTEGER (signed) Numbers
- REAL $\mathbf{L} \qquad$  (signed) Decimal numbers with  $E$  or .
- <code>DOUBLE (signed)</code> Decimal numbers with  $D$  and .
- STRING 72 character long, enclosed in " "
- $LOGICAL = 5True$  LOGICAL  $=$  5.  $T = 1$

Variable names are case sensitive.



#### Assign or Evaluate variables



Variable in data input deck.

- << . >> access the content of a variable send
- >> . << put a value into a variable  $\qquad \qquad \text{recover}$

#### **Managing <sup>a</sup> DRAGON Execution 6** *,*<br>☞ ÉCOLE<br><mark>rTECHNIQUE</mark>

Operations on variables.

- **C** Reverse Polish Notation (value) (operator) (value) ⇔ (value) (value) (operator)
- Arithmetic operations  $+ -$  \* / \*\* Ex : delta := b 2 \*\* 4. a c \* \* -
- **C** Unary operations COS SQRT ABS NOT LN Ex : delta := delta SQRT
- Relational operations  $\langle \rangle \langle \rangle \langle = + -$ Ex : condition := <sup>a</sup> b <=
- Operations on STRING variables <sup>+</sup> -
- NO mixed mode operations



#### IF/THEN/ELSE statement



#### WHILE and REPEAT statement







#### Working with procedures



#### **Managing <sup>a</sup> DRAGON Execution 9** ÉCOLE **POLYTECHNIQUE**

### Flow chart in DRAGON input decks



#### **Managing <sup>a</sup> DRAGON Execution 10** ÉCOLE

Input cross sections

- **C** Macroscopic library
	- MAC: data input
	- EDI: DRAGON calculations
	- MODULE: other transport codes, WIMS TAPE16 ...  $\bullet$
	- LIB: microscopic library
	- $\Rightarrow$  object input in ASM: or FLU: and EDI:

#### **Managing <sup>a</sup> DRAGON Execution 11** 。<br>€COLE

Tracking validation

- Set relatively low line densities and angle number
	- Verify volume and surface integration errors
- EXCELT: Useful print levels EDIT *iprint* :
	- *iprint* = 0 no printing
	- *iprint* <sup>=</sup> 1 (default) geometric information and echo of data input
	- *iprint* = 2 tracking error on volumes and surfaces
	- *iprint* <sup>=</sup> 5 surface and region numbering and description, cell by cell and then global by plane in 3D

#### **Managing <sup>a</sup> DRAGON Execution 12 O**<br>POLYTECHNIQUE<br>MONTRÉAL MONTRÉAL



# **Managing <sup>a</sup> DRAGON Execution 13**

```
EVALUATE DevLocation := "IN" ;
Volumes Tracks := DevGeo :: <<DevType>> <<DevLocation>> ;
PIJ := ASM: Macrolib Volumes Tracks ;
Fluxes := FLU: PIJ Macrolib Volumes :: TYPE B B1 PNL ;
Edition := EDI: Fluxes Macrolib Volumes ::
 COND 0.626 MERGE COMP SAVE
PIJ Volumes Tracks Fluxes := DELETE: PIJ Volumes Tracks Fluxes ;
EVALUATE DevLocation := "OUT" ;
Volumes Tracks := DevGeo :: <<DevType>> <<DevLocation>> ;
PIJ := ASM: Macrolib Volumes Tracks ;
Fluxes := FLU: PIJ Macrolib Volumes :: TYPE B B1 PNL ;
Edition := EDI: Edition Fluxes Macrolib Volumes :: SAVE ;
PIJ Volumes Tracks Fluxes := DELETE: PIJ Volumes Tracks Fluxes ;
Results := Edition
Edition Macrolib := DELETE: Edition Macrolib ;
END: ;
QUIT "LIST" .
```
#### **Managing <sup>a</sup> DRAGON Execution 14 O**<br>
ECOLE<br> **POLYTECHNIQUE**



#### **Managing <sup>a</sup> DRAGON Execution 15 O**<br>
ECOLE<br> **POLYTECHNIQUE**



#### **Managing <sup>a</sup> DRAGON Execution 16 O**<br>POLYTECHNIQUE<br>MONTRÉAL



#### **Managing <sup>a</sup> DRAGON Execution 17** ÉCOLE

```
GEOMETRY := GEO: : CAR3D 2 1 2
 CELL FC1B MD1B FC1T AD1TX- REFL X+ SYME Y- REFL Y+ SYME Z- REFL Z+ SYME::: FC1B := GEO: CARCELZ 5341
  ... ;
  ::: MD1B := GEO: CAR3D 2 4 1
  ... ;
  ::: FC1T := GEO: CARCELZ 5342
  ... ;
  ::: AD1T := GEO: CARCELY 6242
  ... ;
;
Volumes Tracks := EXCELT: GEOMETRY ::MAXR <<NbReq>> TRAK TISO <<NbAngles>> <<TrkDens>> ;
GEOMETRY := DELETE: GEOMETRY ;
QUIT "LIST" .
```


Some comments and warning on the CP method

- **C** The sources are assumed constant inside each region
	- Select an adequate spatial discretization This may lead to <sup>a</sup> large number of region (CP is proportionnal to  $N$ 2 ) Some regions may be very small causing problem with tracking
- Select <sup>a</sup> problem that is not too heterogeneous



### $C$ **onclusions**

- The angular flux on each external surface are assumed constant and isotropic
	- **Try to get rid of external surfaces with re-entrant** angular flux
	- Select a model where the region of interest is far from the external surfaces