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Rates of Vertical Groundwater Movement Estimated from the Earth's Thermal Profile¹

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Abstract. An analytical solution is developed describing vertical steady flow of groundwater and heat through an isotropic, homogeneous, and fully saturated semiconfining layer. A typecurve method for estimating groundwater velocities from temperature data is presented.

Introduction. Geophysicists have long recognized that moving groundwater can affect the flux of heat within the Earth. Van Orstrand [1934, p. 996] indicated that the transfer of heat by migrating water could cause variations of temperature gradients within the Earth. He discussed these variations at some length, using illustrative data from the Wall Creek Sands of Salt Creek Oil Field, Wyoming.

Because the natural heat-flux density from the Earth is usually small, the upward thermal gradient within the Earth may be affected by groundwater movement. *Stallman* [1960] presented the basic equations for the simultaneous transfer of heat and water within the Earth and suggested that temperature measurements might provide a means of measuring rates of groundwater movement.

We are presenting a solution to Stallman's general equation for the case of steady-state vertical flow of both groundwater and heat. This solution is of interest to groundwater hydrologists, because it may afford a means of calculating vertical rates of groundwater movement and in some instances, where head relationships are known, vertical permeabilities.

General differential equation. The general differential equation for simultaneous nonsteady heat and fluid flow through isotropic, homogeneous, and fully saturated porous mediums is [Stallman, 1960]

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$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} - \frac{c_0 \rho_0}{\kappa} \left[\frac{\partial (v_T T)}{\partial x} + \frac{\partial (v_T T)}{\partial y} + \frac{\partial (v_1 T)}{\partial z} \right] = \frac{c \rho}{\kappa} \frac{\partial T}{\partial t}$$

where

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- = Temperature at any point at time t.
- = Specific heat of fluid.
- = Density of fluid.
- = Specific heat of solid-fluid complex.
- = Density of solid-fluid complex.
- = Thermal conductivity of solid-fluid complex.
- v_x, v_y, v_z = Components of fluid velocity in the x, y, and z direction.

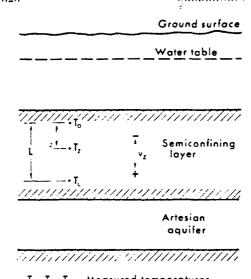
x, y, z = Cartesian coordinates.

t = Time since flow started.

In the problem treated below the flow of heat and fluid is one-dimensional (vertical) and steady. Under these conditions the differential equation reduces to

$$(\partial^2 T | \partial z^2) + (c_0 \rho_0 v_{z-k}) (\partial T | \partial z) = 0$$

Analysis. Consider the vertical steady anisothermal groundwater flow through the semiconfining layer of an aquifer system (Figure 4) in which the temperature along a vertical section of the semiconfining layer has been measured at three or more points. These temperature measurements should satisfy the solution to the following boundary-value problem describing both heat and groundwater flow in the semiconfining layer: 326



 T_0, T_2, T_1 Measured temperatures

v, Leakage rate

Fig. 1. Diagrammatic sketch of typical leaky aquifer.

$(\partial^2 T_z \ \partial z^2) \rightarrow (c_0$	$\rho_0 v_{\pm}(\kappa)$ (d)	$l_{z}(\partial z) = 0$	(1)
$T_{\pm} = T$	o ut	z = ()	(2)
$T_{\pm} = T$	<i>t</i> at	z = I	(1)

where

- T_z = Temperature measurement at any depth z.
- T_{n} = Uppermost temperature measurement.
- T_L = Lowermost temperature measurement.
- L = Length of vertical section over which temperature measurements extend (vertical distance between T_{0} and T_{L}).

and in which the coordinate z is positive downward with origin at T_{00} and the groundwater velocity v_{2} is positive downward.

After solving equation 1 and applying boundary conditions (2) and (3), the solution is

$$(T_z - T_0) (T_L - T_0) = f(\beta, z | L)$$

where

$$(\beta, z, L) = [\exp(\beta z, L) - 1] [\exp(\beta) - 1]$$

and $\beta = c_{\text{montr}_{k}}L_{-k} = a$ dimensionless parameter, that is positive or negative depending on whether v_{k} is downward or upward.

Table 1 gives values of the function $f(\beta, z, L)$ for a practical range of its parameters. Calculations were made using the U. S. Geological Survey's Burroughs 220 computer.

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Application. To estimate the groundwate velocity through the semicontining layer true temperature data it is necessary to solve equation 4 for the velocity e_{ij} which appears in the dimensionless parameter β . This is done graphic ally by a "type-curve" procedure. A set of type curves consisting of arithmetic plots of $f(\beta, \beta, z)L$ against z/L for different values of β , shown or Figure 2, was prepared from Table 1.

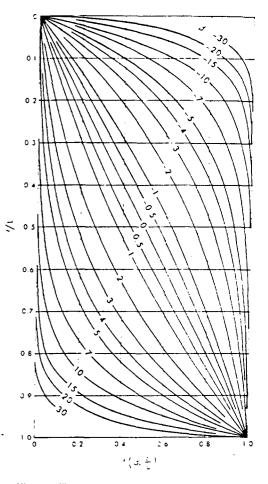


Fig. 2. Type curves of the function f β, z, L .

Ratios $(T_z - T_0)/(T_L - T_0)$ calculated from measured temperature data, are plotted against the depth factor z L at the same scale as the set of type curves. The $T - T_0/(T_L - T_0)$ against z L plot is superimposed on the typecurve set holding the coordinate axes in colucidence. The value of β is interpreted from the type curve that best matches the field data curve.



Vertical Groundwater Morement

TABLE 1.	Values of the Function f (3, z/L	.1
----------	----------------------------------	----

,1,-2 L	0	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9	1.0
0,001	0	0,1000	0,1999	0,3000	0,3999	0, 1999	0.5999	0.6999	0.7999	0,2000	1.000
0.01	0	0.0996	0,1992	0,2990	0.3988	0.4988	0.5988	0.6989	0.7992	0.8995	<u>1.</u> (КЮ)
0.10	Û	0.0956	0,1921	0.2896	0.3880	0.4875	0.5680	0.6894	0.7919	0.8954	1.000
0.25	Ū.	0.0891	0.4805	0.2742	0.3703	0.4688	0.5898	0.6733	0.7795	0.8884	Т.(Н)(Н
0.50	0	0.0790	0.1621	0.2495	0,3413	0.4378	0.5393	0,6460	0.7581	0.8760	1.000
0.75	0	0.0697	0,1449	0.2259	0.3132	0.4073	0.5055	0.6181	0.7360	0.8630	1,000
1.0	0	0.0612	0.1289	0.2036	0.2862	0.3775	0.4785	0,5900	0.7132	0.8495	1.1880
2.0	0	0.0347	0.0770	0.1287	0,1918	0.2689	0.3631	0.4782	0.6187	0.7904	1.1881
3.0	0	0.0183	0.0431	0.0765	0,1216	0.1824	0.2646	0,3755	0.5252	0.7272	1.0км
4.0	Ö	0.0092	0.0229	0.0433	0.0738	0.1192	0.1870	0.2882	0.4391	0.6642	1.000
5,0	0	0.0014	0.0117	0.0236	0.0433	0.7586	0.1295	0.2179	0.3636	0.6039	1.000
6.0	ü	0,0020	0.0058	0.0125	0.0249	0.4743	0.0885	0.1632	0.2995	0.5477	1,000
7.0	Ū.	0,0009	0.0028	0.0065	0.0141	0.2931	0.0600	0.1217	0.2459	0.4961	1.000
5.0	0	0,0004	0.0013	0.0034	0.0079	0,1799	0.0404	0.0204	0.2016	0.4491	1.000
9.0	0	0,0002	0.0006	0.0017	0,0044	0,1099	0.0272	0.0671	0.1652	0.4065	1.000
10	0	0,0001	0.0003	0.0009	0.0024	0.0067	0.0183	0.0497	0.1353	0.3679	1.000
15	0	0.000	0.0000	0.0000	0.0001	0.0006	0.0025	0.0111	0.0498	0.2231	1.(X)
	0	0,000	0.0000	0.0000	0.0000	0.0000	0.0003	0.0025	0.0183	0.1353	1.(88
20 00	ö	0.000	0.0000	0.0000	0.0000	0.0000	0.0000	0.0001	0.0025	0.0498	1.000
	0	0.0000	0.0000	0.0000	0.0000	(1,0000)	0.0000	0.0000	0.0000	0.000	1.000
100	0	0.1000	0,2001	0.3001	0.4001	0.5001	0.6001	0.7001	0.8001	0.9000	1.000
-0.001			0.2001	0.3011	0.4012	0.5012	0.6012	0.7010	0.8008	0.0004	1.008
-0.01	0	0.1005	0.2081	0.3106	0.4120	0.5125	0.6120	0.7104	0.8079	0.9044	1.000
-0.10	0		0.2051 0.2205	0.3267	0.4302	0.5312	0.6297	0.7258	0.8195	0.9109	1.000
-0.25	0	0.1116		0.3540	0.4607	0.5622	0.6587	0,7505	0.8379	0.9210	1.000
-0.50	0	0.1240	0.2419		0.4912	0.5927	0.6868	0.7741	0.8551	0.9303	1.000
-0.75	0	0.1369	0,2640	0.3819 0.4100	0.5215	0.6225	0.7138	0,7964	0.8711	0.9388	1.000
-1.0	0	0.1505	0.2868	0.4100 0.5218	0.6369	0.7311	0.8082	0.8713	0.9230	0.9653	1.000
-2.0	0	0.2096	0.3813	0.5245 0.6245	0.7354	0.8176	0.8784	0.9235	0.9569	0.9817	1.000
-3.0	0	0.2728	0.4748	0.0240 0.7118	0.8130	0.8808	0.9262	0.9567	0.9771	0.9908	1.000
-4.0	()	0.3358	0.5609		0.8705	0.9241	0.9567	0.9764	0.9883	0,9956	1.000
-5.0	0	0.3961	0.6364	0.7821	0.8705 0.9115	0.9526	0.9751	0,9875	0,9942	0.9980	1.000
-0.0	0	0.4523	0.7005	0.8568		0.9707	0.9859	0.9935	0.9972	0.9991	1.000
-7.0	0	0.5039	0.7541	0.8783	0.9400	0.9820	0,9921	0,9966	0.9987	0.9996	1.000
-8.0	0	0.5509	0.7984	0.9096	0.9596		0.9956	0,9983	0,9994	0.9998	1.000
-9.0	0	0.5935	0.8348	0.9329	0.9728	0.9820	0.9976	0,9991	0.9997	0.9999	1,000
-10	0	0.6321	0.8647	0.9503	0.9817	0.9933		1.0000	1.0000	1.0000	1.000
-15	()	0.7769	0.9502	0.9889	0.9975	0.9994	0.9999	1.0000	1,0000	1.(X(K))	1.000
-20	0	0.8647	0.9817	0.9975	0.9997	1.0000	1.0000		1,0000	1.0000	1.400
 30	()	0.9502	0,9975	0.9999	1,19(90)	1,0000	1,000	1,0000			1.000
- 100	0	1.0000	1,0000	1.0000	1,0000	1.0000	1.0000	1.0000	1.0000	1.0000	1.000

The groundwater velocity is calculated from the relation

water from equan the phictype z(L)n on

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 $u = \kappa \beta c_0 \rho_0 L \qquad (5)$

Discussion and conclusions. The type curves (Figure 2) indicate that under isotropic and steady-state conditions with no groundwater flow ($\beta = 0$) the thermal gradient is linear with depth. As groundwater movement occurs, the thermal profile curves and the plot of $(T_z - T_y)^+$ $(T_{|L|} - T_y)$ against z|L| is convex upward or downward, depending upon the direction of water movement. The curvature of the thermal profile increases with increasing groundwater velocity.

The lowest value at which we might expect to detect a curvature in the $(T_z - T_y)/(T_L - T_y)$ against z L plot is $\beta_{\perp} = 0.5$. That is, the lower limit of detectable groundwater velocities is

$v = 0.5\kappa c_0 \rho_0 L$

Typical values for the thermal conductivity of water-saturated clay are approximately $\kappa=2$ \times

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10⁻³ calorie cm sec ^oC [Birch, 1942, p. 259]. The gradient, For a gradient of 1^ee (1 cm m), which quantity $\kappa' c_{0,0,0}$ is therefore approximately

$$\frac{\kappa}{c_0\rho_0} = \frac{2 \times 10^{-3} \text{ cal } \text{cm} - \text{sec} - ^\circ C}{1 \text{ cal } g - ^\circ C \times 1 g \text{ cm}^3}$$
$$= 2 \times 10^{-3} \text{ cm}^2 \text{ sec}$$

With this value of $\kappa | c_{\mu\rho\rho}$ the least magnitude of e_2 that can be determined is

$$|v_z| = \frac{0.5 \times 2 \times 10^{-3} \text{cm}^2 / \text{sec}}{L}$$
$$= \frac{1 \times 10^{-3} \text{cm}^2 / \text{sec}}{L}$$

The magnitude of ε_{i} depends on L; if, for example, L = 40 meters, the least magnitude is

$$|v_{z}| = \frac{1 \times 10^{-3} \text{ cm}^{2} \text{ sec}}{1 \times 10^{3} \text{ cm}}$$

 $= 1 \times 10^{-6}$ cm sec

or approximately 0.1 cm day. From Darcy's law

$$v_z = K dh dz$$

where K = permeability (hydraulic conductivity) of the semiconfining layer, and h = head at any depth z; we can also suggest a lower limit of detectable permeability by assuming a head is easily observed in the field.

$$dh dz = 1 \times 10^{-2} \,\mathrm{cm} \,\mathrm{cm}$$

and $e_t = 1 \times 10^{-6}$ cm sec

$$K = \frac{1 \times 10^{-6}}{1 \times 10^{-2}} = 1 \times 10^{-4} \,\mathrm{cm/sec}$$

Permeabilities of this order of magnitude are in the range commonly associated with semiconfining layers. As Stallman suggests, it appears that temperature measurements afford a means of detecting vertical groundwater movement, a problem that continues to trouble groundwater hydrologists.

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